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# The MPR-300 Reactor System for use in Martian Applications

(NASA-CR-184747) THE MPR-300 REACTOR SYSTEM  
FOR USE IN MARTIAN APPLICATIONS Final Report  
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by

Ted G. Bagwell  
Oscar J. Lessard  
John A. Rennie  
Sandra M. Sloan  
John D. Valentine

*TEXAS A&M UNIV.*  
*COLLEGE STATION*

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## I. Attributes

All project members had an active role in the development of this project. The decision on which project to do, and well as the major decisions of the final project, were a collaborative effort by the entire group. The calculations of the individual systems were divided among the specific group members. Table I.1 basically shows the tasks assigned to each group member.

Table I.1 Member Responsibilities

Ted G. Bagwell	<ul style="list-style-type: none"><li>• group leader</li><li>• coordination of activities</li><li>• reactor thermal-hydraulics</li></ul>
Oscar J. Lessard	<ul style="list-style-type: none"><li>• core neutronics</li><li>• materials considerations</li><li>• background research in gas reactors</li></ul>
John A. Rennie	<ul style="list-style-type: none"><li>• space logistics</li><li>• micrometeoroid shielding</li><li>• solar irradiation analysis</li></ul>
Sandra M. Sloan	<ul style="list-style-type: none"><li>• heat rejection system</li></ul>
John Valentine	<ul style="list-style-type: none"><li>• cycle thermodynamics</li><li>• radiation shielding</li><li>• turbomachinery</li></ul>

## II. Introduction

The motivation behind the design project was to provide power for propulsion for an unmanned GEO to Mars mission in a 2 to 6 month time span and to provide 6 years of terrestrial power on Mars. Research was conducted on the German HTR reactor, current U.S. government journal articles on particle bed space reactors, and the British MAGNOX carbon dioxide reactor of the 1960's. After this background research it was determined that a carbon dioxide gas cooled pebble fueled reactor would be feasible to meet our initial motivation. The main reasons for this choice were the availability of carbon dioxide on the Martian surface and the small reactor size with high power densities achieved with particulate type fuel.

The design objectives consisted of obtaining an approximate electrical power level of 300 kW in order to provide power for MPD thrusters for the GEO to Mars journey and obtaining a high cycle efficiency in order to keep the reactor thermal power output around 1 MW. The entire system mass objective was under 10,000 kg in order to ensure that the system could feasibly fit on one shuttle payload. This led to the choice of heat pipe radiators for waste heat rejection as the total mass of the heat rejection system was of utmost concern. The reliability provided by a redundant system was another design objective in order to meet mission goals.

The report outline consists of seven different sections which are contained in the body of the report. The first section consists of neutronics which calculates flux distributions and fuel requirements. The second section is made up of thermal-hydraulics considerations for calculating reactor core temperature and pressure characteristics. Section three consists of the thermodynamic cycle calculations which defines states and arrives at an overall cycle efficiency. The turbomachinery selections can be seen in section four. The heat pipe radiators for waste heat rejection are explained in section five. The sixth section contains the shadow shield configuration necessary to protect electronic components from radioactivity. Section seven is comprised of propulsion and space logistics in order to successfully complete the GEO to Mars mission.

### III. Body

## 1. NEUTRONICS

### Introduction:

A pebble bed nuclear space reactor is needed to supply power for propulsion and instrumentation from GEO to Mars and to supply terrestrial power on the Martian surface. The entire lifetime of the reactor is limited to seven years. The power required is one megawatt thermal.

The concept of the pebble bed reactor is based on a combination of the German HTR and recent research on the fixed particle bed space reactor. The advantages of this type of reactor are its small size and its high power density due to the large surface area of the fuel particles <sup>1</sup>. The gas coolant will be carbon dioxide because of the availability of this gas on the Martian surface. The Martian atmosphere is over 80% CO<sub>2</sub>. The use of CO<sub>2</sub> in a reactor is based on the British MAGNOX reactor of the 1960's <sup>2</sup>. The only serious problems encountered with this design is the carbon dioxide and graphite corrosion activation at temperatures exceeding 810K.

The conceptual design of the pebble bed reactor can be seen in figure 1.1. The turbine-compressor drive shaft goes through the center of the reactor and BeO/B<sub>4</sub>C control drums are used for reactivity control. The fuel is a BISO type consisting of UN fuel and a PyC cladding. A very thin coating of stainless steel makes up the outer surface of the fuel pebble in case temperatures exceed the 810K discussed above. It can be observed that this is an axial flow pebble bed reactor.

### General Assumptions:

Certain assumptions were made before neutronics calculations could be accomplished. Since the reactor core consisted of an annular cylindrical shape, it was difficult to choose a corresponding computer code for neutronics calculations. It was finally decided that a two group approximation using a finite difference approach would be used. This allowed one to determine the radial flux profile by breaking the center turbine-compressor, the reactor vessel walls, the reactor core, and the control drums into separate regions for neutron flux approximations. The axial

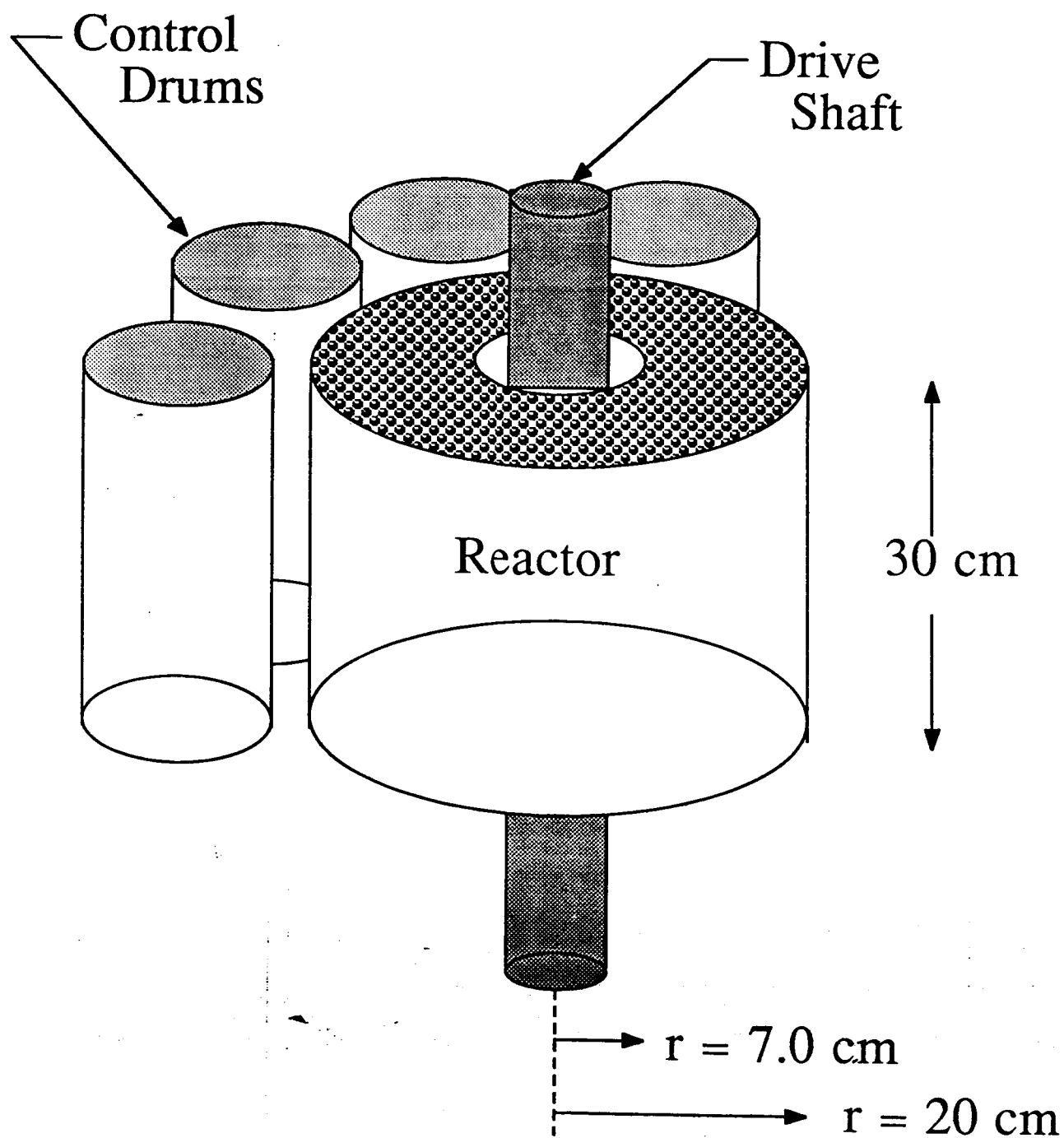


Figure 1.1 Conceptual Reactor Design

neutron flux distribution was then assumed to be a cosine shape which is the normal case for cylindrical geometry.

The gas coolant voiding in the reactor core was assumed to be 37%. This is true as long as the annular radius is larger than five fuel pebble diameters<sup>3</sup>. The other 63% of the reactor core volume was assumed to consist of equal volumes of fuel and cladding. This was considered a reasonable estimate as the fuel radius was 79% of the entire fuel pebble radius. This should allow sufficient fission product gathering in the BISO fuel pebble. The fuel pebble concept can be seen in figure 1.2.

The center turbine-compressor shaft was used as it was our intention to have a reactor diameter equalling that of the turbine and compressor. This would allow for a jet engine type arrangement where the turbine and compressor are welded directly to the reactor vessel. This would also allow for reduced piping and reduced overall area required for the reactor assembly.

The choice of the center shaft made it necessary for one to have an intermediate reactor instead of a fast reactor. This reactor type was chosen to keep the fast neutron flux as low as possible on the surface of the shaft. Neutron embrittlement could become a serious problem if the center shaft is exposed to high neutron fluxes.

These are a few of the general assumptions needed before neutronics modelling could be considered. Many other minor assumptions will be discussed in the text as deemed necessary.

### Neutron Flux Modelling:

The DIF2DK two group, one dimensional diffusion theory code supplied in the NUEN 429, Spring 1987 course was used for neutronics modelling. The two group equations as solved by the computer code are written as<sup>4</sup>:

$$\nabla \cdot D_1 \nabla \phi_1 - \Sigma_{a1} \phi_1 - \Sigma_{1 \rightarrow 2} \phi_1 + \frac{1}{K_{eff}} (v \Sigma_{f1} \phi_1 + v \Sigma_{f2} \phi_2) = 0 \quad (1.1)$$

and

$$\nabla \cdot D_2 \nabla \phi_2 - \Sigma_{a2} \phi_2 + \Sigma_{1 \rightarrow 2} \phi_1 = 0 \quad (1.2)$$



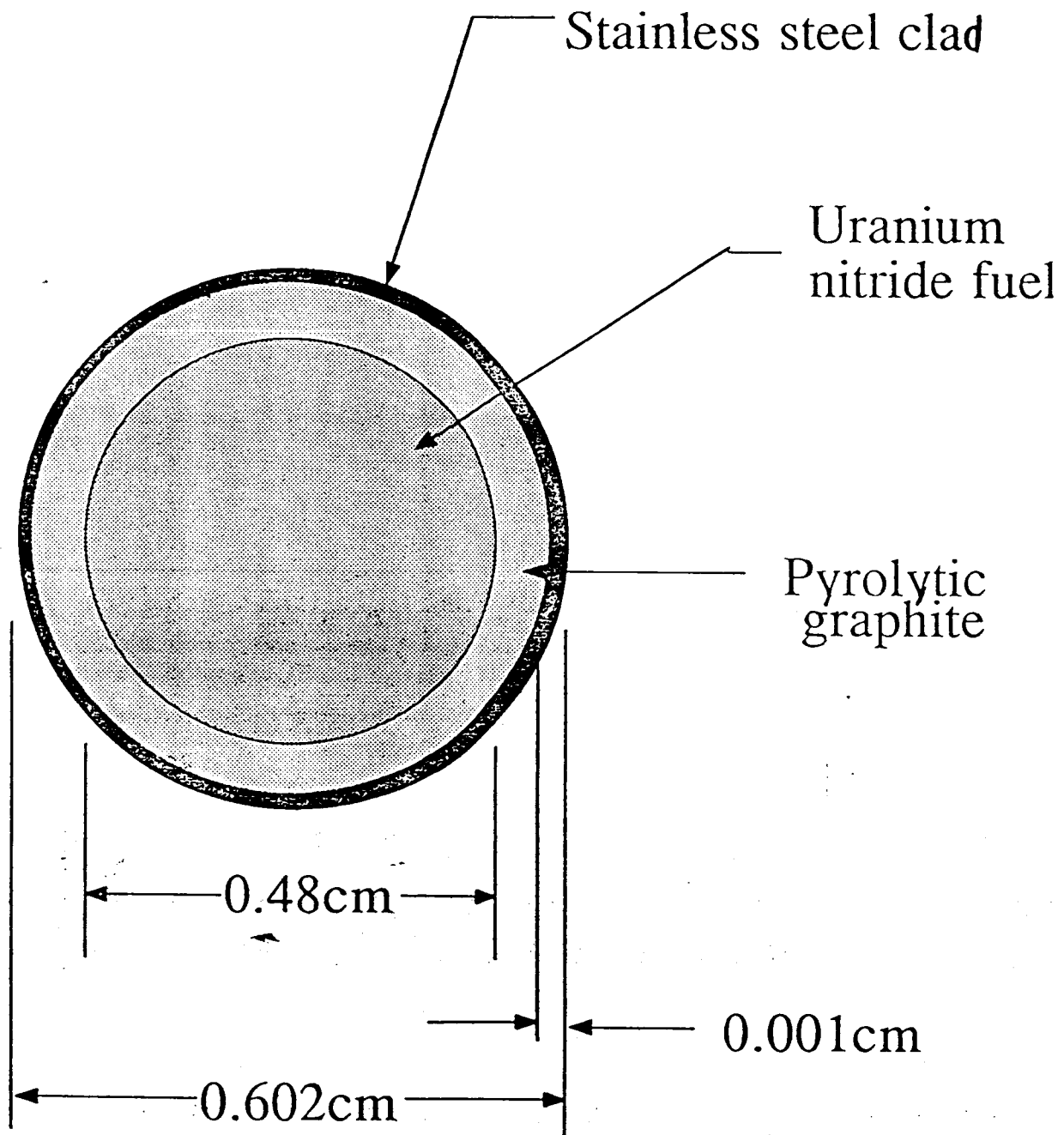


Figure 1.2 Pebble Fuel Design

where :

- $D_1$  = fast group diffusion coefficient
- $D_2$  = thermal group diffusion coefficient
- $K_{eff}$  = multiplication factor
- $\Sigma_{a1}$  = macroscopic fast group absorption cross section
- $\Sigma_{a2}$  = macroscopic thermal absorption cross section
- $\Sigma_{1 \rightarrow 2}$  = macroscopic slowing down cross section
- $\Sigma_{f1}$  = macroscopic fission cross section (fast)
- $\Sigma_{f2}$  = macroscopic fission cross section (thermal)
- $\nu$  = neutrons released per fission
- $\phi_1$  = fast neutron flux
- $\phi_2$  = thermal neutron flux

The code was slightly modified to accept the input thermal power and energy per fission to calculate the actual neutron flux values in the radial direction. The equations used were:

$$A = \frac{P}{K \sum_{i=1}^N \left( \nu \Sigma_{f_i} \frac{\Delta x_i}{2} + \nu \Sigma_{f_{i+1}} \frac{\Delta x_{i+1}}{2} \right) \tilde{\phi}_i} \quad (1.3) \quad \text{and} \quad \phi_i = a \tilde{\phi}_i \quad (1.4)$$

where :

- $P$  = power in MeV
- $K$  = energy per fission
- $\tilde{\phi}_i$  = normalized flux
- $\phi_i$  = actual flux
- $\nu$  = neutrons released per fission
- $\Sigma_f$  = macroscopic fission cross section
- $i$  = number of intervals in finite difference approach

A finite difference approach is used to solve the two group problem. This allows one to define different material regions and the input of DIF2DK requires the neutronic properties of each of these regions. These properties are the diffusion coefficient and the macroscopic absorption, slowing down, and nu-fission cross sections.

A computer code called START was written to solve for the properties needed for the DIF2DK input deck. The microscopic absorption, slowing down, nu-fission, and transport cross sections are required in order to calculate the properties needed. In order to come up with the two

group microscopic cross sections, it was necessary to use the computer code ANISN to collapse the 27 group cross sections. For this intermediate reactor the thermal group was defined from 0 to 1 eV and the fast group was defined from 1 eV to infinity. This is an accepted energy group distribution for an intermediate reactor <sup>5</sup>. The code START calculates the required number densities from the given material densities in order to determine the macroscopic cross sections as follows:

$$N = \frac{\rho A_0}{M} \quad (1.5)$$

where :

$N$  = atom number density

$A_0$  = Avogadro's number ( $6.02 \times 10^{23}$ )

$\rho$  = density

$M$  = atomic or molecular weight

The diffusion coefficient is determined by:

$$D = \frac{1}{3 \Sigma_{tr}} \quad (1.6)$$

where :

$D$  = diffusion coefficient

$\Sigma_{tr}$  = macroscopic transport cross section

The detailed input deck of START can be seen in table 1.1 and the output description can be seen in table 1.2.

The START program is capable of accepting different fuel enrichments, fuel and cladding volume ratios, and fraction of BeO seen in the control drum compared to B<sub>4</sub>C in the radial direction. The output of START is then used as the input to DIF2DK. The input deck of DIF2DK is described in table 1.3. The stainless steel reactor vessel walls and the center turbine-compressor

**Table 1.1**  
**Input Deck of START**

The order of microscopic cross section input is as follows:  
absorption, slowing down, nu-fission, and transport

Line 1: Oxygen-16, fast  
Line 2: Oxygen-16, thermal  
Line 3: Uranium-235, fast  
Line 4: Uranium-235, thermal  
Line 5: Uranium-238, fast  
Line 6: Uranium-238, thermal  
Line 7: Carbon-12, fast  
Line 8: Carbon-12, thermal  
Line 9: Iron, fast  
Line 10: Iron, thermal  
Line 11: Beryllium-9, fast  
Line 12: Beryllium-9, thermal  
Line 13: Boron, fast  
Line 14: Boron, thermal  
Line 15: Nitrogen, fast  
Line 16: Nitrogen, thermal

**Table 1.2:**

**Output of START**

Line 1: Macroscopic absorption cross sections in the following order:

U-235 fast, U-235 thermal, U-238 fast, U-238 thermal

Line 2: Macroscopic nu-fission cross sections in the same order

The remaining lines have the following output order:

diffusion coefficient, macroscopic absorption cross section,  
macroscopic slowing down cross section, macroscopic nu-fission  
cross section

Line 3: Reactor core properties, fast

Line 4: Reactor core properties, thermal

Line 5: Reactor vessel walls and center shaft properties, fast

Line 6: Reactor vessel walls and center shaft properties, thermal

Line 7: Control drum properties, fast

Line 8: Control drum properties, thermal

**Table 1.3:**  
**Input Deck of DIF2DK**

Line 1: # of regions, # of materials, fractional convergence criterion, thermal power  
(Mev), energy per fission

Line 2: geometry (2-cylinder), left B.C. (1-symmetry), right B.C. (0-zero flux at extr.  
boundary)

Line 3: region #1, material in region, # intervals in region, region thickness

Line 4: same as above for region #2

Line 5: same as above for region #3

Line 6: same as above for region #4

Line 7: material #1, fast diffusion coefficient, fast macroscopic absorption cross section,  
fast macroscopic slowing down cross section, fast macroscopic nu-fission cross  
section

Line 8: same as line 7 for thermal properties

Line 9: same as above for material #2 fast properties

Line 10: same as above for material #2 thermal properties

Line 11: same as above for material #3 fast properties

Line 12: same as above for material #3 thermal properties

    \*\* regions

        1) center shaft and inner vessel wall

        2) reactor core

        3) outer vessel wall

        4) control drums

    \*\* materials

        1) homogeneous reactor core

        2) iron

        3) control drums ( $\text{BeO}/\text{B}_4\text{C}$ )

shaft were assumed to be iron for neutron cross section purposes. The reactor core was assumed to be a homogeneous mixture of equal volume PyC cladding and UN fuel and a volume fraction of 0.37 of CO<sub>2</sub> gas. The control drums were assumed to be one region with the option of varying the amounts of BeO and B<sub>4</sub>C. The output of the DIF2DK program consists of the normalized fast and thermal neutron fluxes, the actual fluxes and the value of Keff. START and DIF2DK program listings can be seen in appendices 1.1 and 1.2 respectively.

### Neutron Flux Calculations:

The initial inputs of the DIF2DK program consisted of a 7 cm region of center shaft and inner vessel wall (3 cm shaft radius and 4 cm inner vessel wall region), a 10 cm region of reactor core, a 3 cm region of outer vessel wall, and a 20 cm region of control drums. The fuel enrichment and annular core radius were varied numerous times and it was finally determined that a fuel enrichment of 5.6% resulted in a 12.8 cm annular core radius for criticality or Keff equal to one. An infinite reflector (control drum) was approached with a region equal to 10 cm. This was a good result as it produced a reactor assembly of about 65 cm in total diameter which compares favorably well with turbine and compressor diameters. It was also a favorable result when considering the magnitude of the neutron flux at the center drive shaft surface. The thermal and fast flux profiles can be seen in figures 1.3 and 1.4 respectively. The fast flux at the surface of the shaft was less than 4E14 and the thermal flux was less than 2E13. These flux values do not produce any neutron embrittlement problems for the steel shaft.

The ratio of the fast flux to thermal flux throughout the reactor core was on the order of 15 to 1 with the peak fast flux in the core being 7.3E14 and the peak thermal flux being 4.7E13. The peak to average flux ratio for both thermal and fast fluxes was on the order of 1.12. The average flux ratio was calculated by averaging the 20 data points within the core. This peak to average flux ratio and the average fluxes were passed on for thermohydraulic considerations. The equation used to break the power into fast and thermal components is given by:

$$\rho''' = K \left[ (\sum_{f_{2b}} \phi_1 + \sum_{f_{2s}} \phi_1)_{fast} + (\sum_{f_{2b}} \phi_2 + \sum_{f_{2s}} \phi_2)_{slow} \right] \quad (1.7)$$

# THERMAL FLUX PROFILE

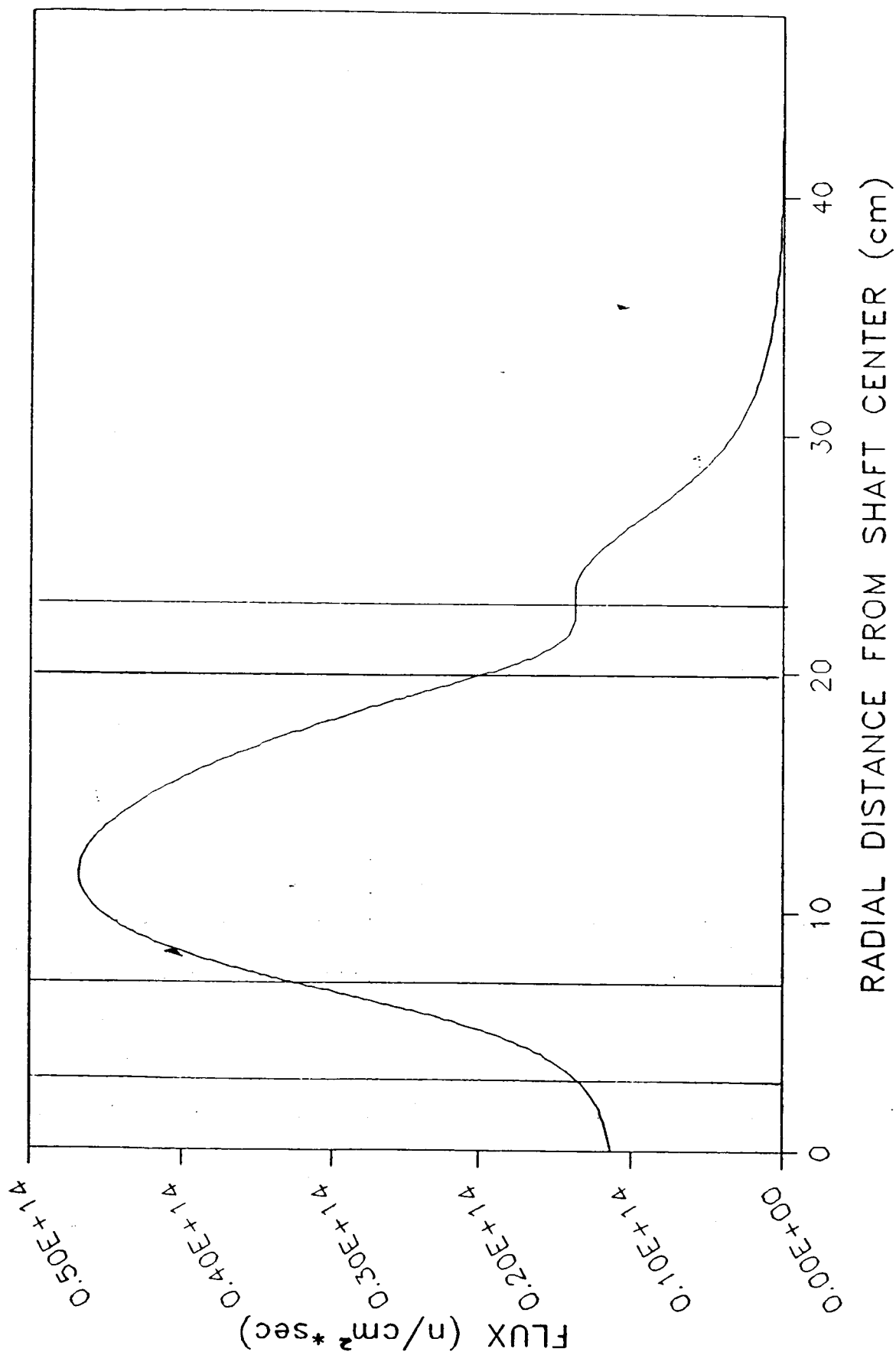


Figure 1.3 Thermal Flux Profile



# FLUX PROFILE

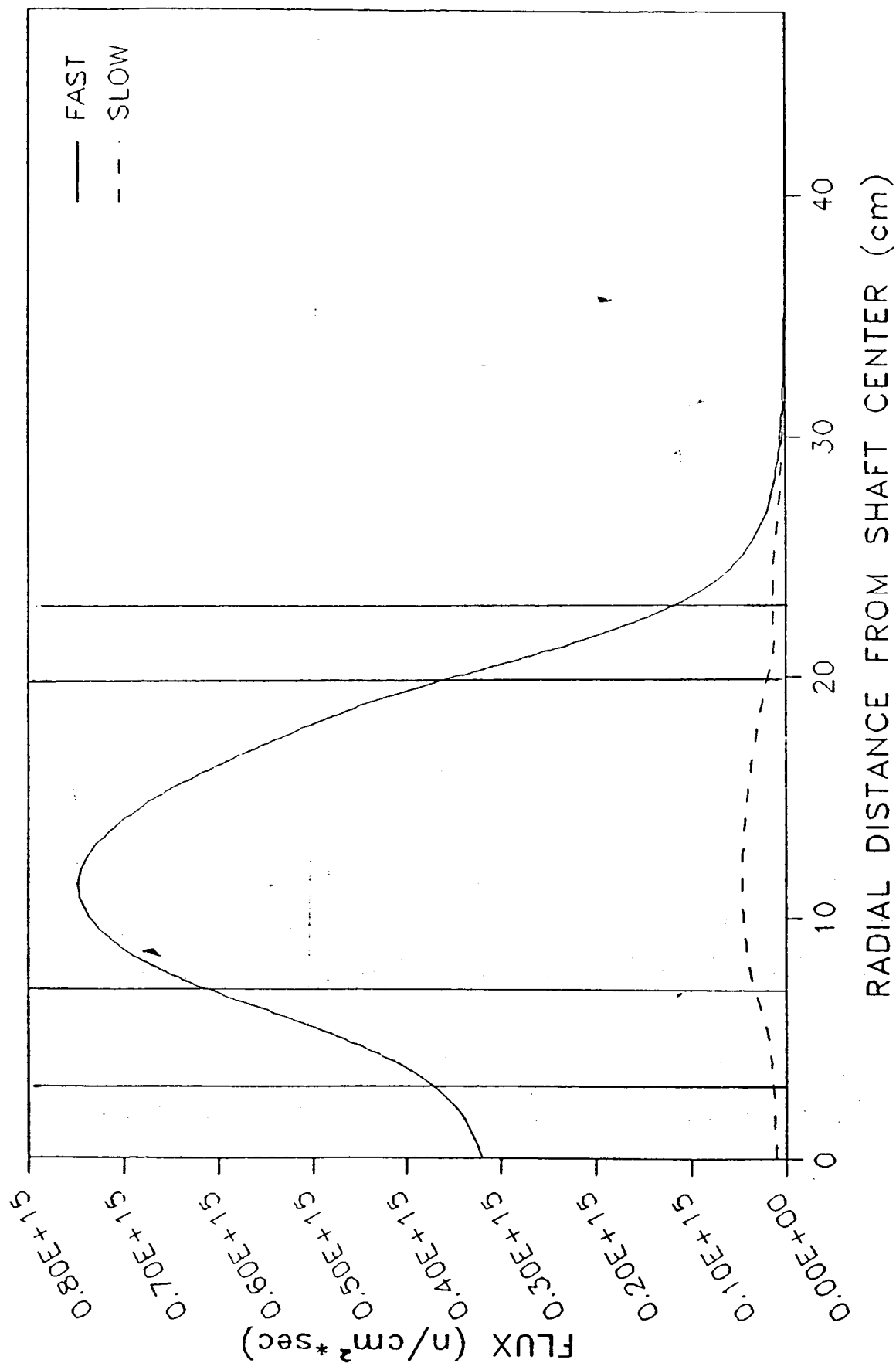


Figure 1.4 Fast and Thermal Flux Profile

where :

$q'''$  = power density

$K$  = energy released per fission

$\Sigma_{f20}$  = macroscopic fission cross section of U-238

$\Sigma_{f25}$  = macroscopic fission cross section of U-235

$\phi_1$  = fast flux

$\phi_2$  = thermal flux

After performing the calculation it was observed that 91% of the power is generated by the thermal flux as is expected even though the fast flux magnitude is much greater.

The method used to account for fuel burnup during the seven year lifetime will be to increase the fuel enrichment. The axial flux distribution will be of the cosine shape as discussed previously for cylindrical geometry. The annular radius of 12.8 cm of the core can be turned into an equivalent cylinder radius by assuming that the volume of the annular core is equal to the volume of a cylinder. The equivalent cylinder radius becomes 18.5 cm. With a reactor height of 30 cm, the height to diameter ratio of the reactor core becomes 30/37 or 0.81. This value compares well to the other HTGR reactors which have height to diameter ratios of 0.75 to 0.90. A complete summary of reactor parameters can be seen in table 1.4.

#### Fuel Burnup Considerations:

The burnup of U-235 in the UN fuel in the reactor core is due to consumption by fission and parasitic absorption of U-235 by radioactive capture. Assuming a thermal reactor, the equation used to calculate the burnup of U-235 is given by <sup>6</sup>:

$$\text{Consumption rate} = 1.05 (1 + \alpha) P \quad \text{g/day} \quad (1.8)$$

**Table 1.4:**  
**Reactor Parameters**

Thermal Output	1 MWt
Electrical Output	300 kWe
Power Density	3.1 W/cm <sup>3</sup>
Active Core Volume	32,330 cm <sup>3</sup>
CO2 fraction	0.37
UN fraction	0.315
PyC fraction	0.315
Peak to Average Flux Ratio	1.12
Power Production	
Thermal Flux	91%
Fast Flux	9%
Control Drum Diameter	10 cm
Reactor core annular radius	12.8 cm
Central Shaft Diameter	6 cm
Reactor Vessel Wall Diameter	
inner	4 cm
outer	3 cm
Peak Fluxes	
thermal	4.7E13
fast	7.3E14
Reactor Height	30 cm
Height to diameter ratio	0.81

where :

$\alpha = 0.169$  which is the radiative capture or parasitic absorption factor

$P =$  power in MeV

The value of alpha accounts for the parasitic absorptions and is 0.169 for U-235. Using the above equation it can be seen that 3.16 kg of U-235 will be consumed during the seven year lifetime. Since 7.61 kg of U-235 is available for the 5.6% enriched reactor fuel, it can be calculated that an 8.0% enriched fuel is needed for a total of 10.77 kg of U-235 at reactor startup. Detailed calculations on fuel burnup can be seen in appendix 1.3. The value of  $K_{eff}$  at startup using DIF2DK and the 8.0% enriched fuel is 1.228.

#### Poisoning Considerations:

Since this is an intermediate reactor, the fission product poisoning of Xenon-135 and Samarium-149 can be neglected. This is true as the buildup of these fission products is negligible in an intermediate space reactor.

#### Decay Heat Considerations:

After a few days of reactor operation, the beta and gamma radiation emitted from decaying fission products amounts to approximately seven percent of the total thermal power output of the reactor. For the one megawatt reactor, the amount of power available at shutdown due to decay heat will be slightly less than 70 watts of thermal power. Recall that 91% of the power produced is from the thermal spectrum. At any time after shutdown, the ratio of power due to decay heat to the original power is given by 7:

$$* \frac{P(t_0, t_s)}{P_0} = \frac{P(t_s)}{P_0} - \frac{P(t_0 + t_s)}{P_0} \quad (1.9)$$

where :

$P =$  power

$P_0 =$  original power before shutdown

$t_0$  = finite time of reactor operation

$t_s$  = time since reactor shutdown

\* Use Figure 8.3 of "Introduction to Nuclear Engineering" by Lamarsh to obtain results

The use of this equation along with the guidance in Lamarsh <sup>7</sup> results in an available power of about 14 watts at one hour after shutdown. This decay heat has been considered and should not be of much concern because of the low power output of the pebble bed space reactor. A method of cooling could be accomplished by circulating the gas coolant through the core after shutdown.

#### **Fuel Considerations:**

The fuel pebble is 0.602 cm in diameter and is composed of equal volumes of 8.0% enriched uranium nitride fuel and pyrolytic graphite cladding thus forming a BISO type fuel pebble. The PyC was chosen because of its outstanding heat transfer and fission product gathering characteristics as revealed by previous space reactor design work <sup>1</sup>. A thin coat of stainless steel is placed on the outer surface of the fuel pebble in case the coolant temperature in the core exceeds 810K <sup>8</sup>. At this temperature the CO<sub>2</sub> coolant and PyC cladding of the fuel will undergo corrosion activation. The equal volumes of fuel and cladding result in the fuel radius being 79% of the entire fuel radius. This provides plenty of fission product gathering when comparing the fuel pebble to other BISO fuels. After simple volume and number density calculations (see appendix 1.3), a list of fuel characteristics was created in table 1.5.

#### **Control Drums and Reactivity Control:**

The BeO and B<sub>4</sub>C control drums serve as a means for reactivity control. There are a total of seventeen 10 cm diameter drums surrounding the reactor core. 120° of the drum is made up of the boron absorber and the other 240° is made up of the beryllium reflector. The drums are capable of

**Table 1.5:**  
**Fuel Characteristics**

UN Fuel Mass	145.7 kg
Fuel Enrichment	8.0%
Mass of U-235	10.8 kg
critical mass at startup	7.6 kg
fuel burnup	3.2 kg
Specific Power	6.86 kW/kg UN
# of Fuel Pebbles	180,000
Pebble Diameter	0.602 cm

adding positive or negative reactivity by letting the reactor core see either the reflector or absorber in the radial direction. Any combination of reflector and absorber is possible with the control drum method.

The reactor also has a mechanism to inject a poison boron dust to cause reactor shutdown as the boron in the control drums may or may not have the capability of shutting down the reactor especially just after startup.

#### **Other Materials Considerations:**

The top and bottom of the reactor core are constructed with a porous grate to allow maximum carbon dioxide flow with enough support to hold the fuel pebbles in place. The thermohydraulics section will discuss these grates in more detail.

The center turbine-compressor drive shaft design is limited by the temperature, the neutron flux at its surface, and the speed of rotation. The maximum temperature will be less than 1000K and the maximum surface flux will be  $4E14$ . These two factors along with normal turbine and compressor drive shaft rotation will allow for the use of almost any ordinary steel. The shaft will be made of a high quality steel of 6 cm in diameter. This should allow for no malfunctions of the shaft due to fatigue or shearing stresses over the seven year lifetime of the system.

The radius of the inner reactor vessel wall is 4 cm thick and the outer vessel wall is 3 cm thick. The extra 1 cm of the inner wall is used to help reduce the neutron flux at the drive shaft surface.

#### **Masses of the Reactor System:**

The masses of the reactor system were calculated using the volumes and densities as follows:

$$m = V \rho \quad (410)$$

Where :

$m$  = mass

$\rho$  = density of reactor component

$V$  = volume of reactor component

The sample calculations can be seen in appendix 1.3 and the masses of the reactor system can be seen in table 1.6. The mass of the control drum drive assembly was assumed to equal the mass of the 17 control drums. This appeared to be a conservative estimate.

### **Conclusion and Recommendations:**

The pebble bed space reactor has been designed after many iterations of optimization and appears to have a good chance of becoming reality. Future work would consist of even more iterations of the entire system and more in depth studies of the effects and reliability of the center drive shaft. A smaller reactor system with higher fuel enrichments may be a possible alternative also.



**Table 1.6:**  
**Reactor System Masses**

UN fuel	145.7 kg
PyC cladding	16.3 kg
CO <sub>2</sub>	0.02 kg
Steel shaft, vessel walls	131.0 kg
Control drums	114.5 kg
Control drum drive assembly	114.5 kg
Total reactor system	522 kg

## References:

- 1) J.R. Powell and T.E. Potts, "FBR and RBR Particle Bed Space Reactors", BNL-33058, 8 pages, Aug 1983.
- 2) G.R. Bainbridge and T.N. Marsham, "Three Generations of Nuclear Power Stations", (United Kingdom Atomic Energy Authority, Risley, England), J. Inst. Fuel, 41, 280-8, Jul 1968.
- 3) S. T. Robinson, " The Pebble Bed Reactor ", Proceedings of the Symposium on Gas Cooled Reactors, Journal of the Franklin Institute, Lancaster, Pa, p. 94, 1960.
- 4) DIF2DK Computer Code, Dr. Parrish, NUEN 429, Spring 1987.
- 5) ANL-5800, Nuclear Reactor Constants, USAEC , p. 581, Jul 1963.
- 6) J.R. Lamarsh, "Introduction to Nuclear Engineering", Chap 3, p. 78, Addison-Wesley Publishing Company, Reading Mass., 1983.
- 7) Lamarsh, Chap 8, p.350-2.
- 8) J.K. Pickard, "Nuclear Power Reactors", Chap 8, p. 310-2, D. Van Nostrand Company Inc, Princeton NJ, 1957.

## Appendix 1.1

DIF2DK Code

4	3	.0001	1.0	200.0
2	1	0		
1	2	7		
2	1	20	12.8	
3	2	3		
4	3	20	25	
1	.84061	.16336	.01806	.06709
1	.38590	.10622	0.0	1.09040
2	.55914	.01214	.02300	0.0
2	.30157	.18987	0.0	0.0
3	.59471	.00038	.09287	0.0
3	.49701	.00067	0.0	0.0

THIS PROGRAM WILL SOLVE THE TWO GROUP-ONE DIMENSIONAL FINITE DIFFERENCE APPROXIMATION TO THE DIFFUSION EQUATION. THE NEUTRON SOURCE IS DUE ONLY TO FISSION AND THEREFORE THE SOLUTION CONSISTS OF DETERMINING KEFF(LARGEST EIGENVALUE) AND THE NORMALIZED FLUX(FUNDAMENTAL MODE) SHAPE. SEVERAL MATERIALS AND BOTH VACUUM AND SYMMETRY BOUNDARY CONDITIONS ARE ALLOWED. UP TO 20 REGIONS CAN BE USED WITH SLAB, CYLINDER, AND SPHERE GEOMETRIES.

# LIST OF PARAMETERS

M = NUMBER OF MATERIALS BEING CONSIDERED (MAX=20)  
 TH = THICKNESS OF MATERIAL SECTIONS  
 INT = NUMBER OF MESH INTERVALS PER MATERIAL  
 DM = DIFFUSION COEFFICIENT  
 SIGA = ABSORPTION CROSS SECTION  
 FIS = NU TIMES FISSION CROSS SECTION  
 D = EXTRAPOLATION LENGTH  
 XKN = NEW K-EFFECTIVE APPROX.  
 XKO = OLD K-EFFECTIVE APPROX.  
 XLAMN = NEW EIGENVALUE APPROX.  
 XLAMO = OLD EIGENVALUE APPROX.  
 FLUXN = NEW FLUX APPROX.  
 FLUXO = OLD FLUX APPROX.  
 IBNL = INDICATOR FOR LEFT BOUNDARY CONDITION  
 IBNR = INDICATOR FOR RIGHT BOUNDARY CONDITION

SIZE IS LIMITED TO 100 MESH INTERVALS

READ IN INITIAL VALUES FOR PROBLEM DESCRIPTION

# ASSIGN VALUES TO THE MESH

```

NPTS=0
DO 101 L=1,NREG
  NPTS=NPTS+INT(L)
  NPTS=NPTS+1
  NPTS1=NPTS-1
DO 102 L=1,NREG
  DT(L)=TH(L)/INT(L)
  X(1)=0.0
  L2=1
DO 103 L=1,NREG
  LXX=INT(L)
DO 104 L1=1,LXX
  J=MR(L)
  L2=L2+1
  X(L2)=X(L2-1)+DT(L)
  DIFC(L2)=DM(1,J)
  SABS(L2)=SIGA(1,J)
  STR(L2)=SIG12(1,J)
  XNF(L2)=FIS(1,J)
  DIFT(L2)=DM(2,J)
  SABT(L2)=SIGA(2,J)
  XNT(L2)=FIS(2,J)
  DX(L2)=X(L2)-X(L2-1)
104 CONTINUE
103 WRITE(6,7171)
  WRITE(6,466)
  DIFC(1)=0.0
  SABS(1)=0.0
  STR(1)=0.0
  XNF(1)=0.0
DO 488 I=1,NPTS
  WRITE(6,467)I,X(I),DIFC(I),SABS(I),STR(I),XNF(I)
  467 FORMAT(I4,F14.4,2X,D12.5,2X,D12.5,2X,D12.5)
  466 FORMAT(//,PT',6X,'POSITION',5X,'DIFF COEF',5X,'ABS XSECT'
    1,7X,'SIG1-2',2X,'NU SIGF XSCT'//)
  WRITE(6,7172)
  WRITE(6,4668)
  DIFT(1)=0.0
  SABT(1)=0.0
  XNT(1)=0.0
DO 4883 I=1,NPTS
  4883 WRITE(6,4675)I,X(I),DIFT(I),SABT(I),XNT(I)
  4675 FORMAT(I4,F14.4,2X,D12.5,2X,D12.5,2X,D12.5)
  4668 FORMAT(//,PT',6X,'POSITION',5X,'DIFF COEF',5X,'ABS XSECT'
    1,2X,'NU SIGF XSCT'//)
DO 3333 J=2,NPTS
  3333 SABS(J)=SABS(J)+STR(J)

```

CALCULATE RIGHT AND LEFT VOLUMES AND AREAS FOR EACH

MESH POINT

```

99  YN1=SABS(L)*DLV(L)+SABS(L+1)*DRV(L)+YN1
    YD1=(XNF(L)+XNT(L)*R)*DLV(L)+(XNF(L+1)+XNT(L+1)*R)
    1  *DRV(L)+YD1
    FLUXO(NPTS)=1.0
    FLUXP(NPTS)=1.0
    FLUOT(NPTS)=STR(NPTS)/SABT(NPTS)

    TOTVOL=TOTVOL+DLV(NPTS)
    YN1=YN1+SABS(2)*DRV(1)
    YD1=YD1+(XNF(2)+XNT(2))*FLUOT(2)*DRV(1)
    YN1=YN1+SABS(NPTS)*DLV(NPTS)
    YD1=YD1+(XNF(NPTS)+XNT(NPTS))*FLUOT(NPTS)/FLUXO(NPTS))*DLV(NPTS)
    XLAMI=YN1/YD1
    XLAMO=XLAMI
    XKO=1./XLAMO

599  ICNT=0
    CONTINUE
    ICNT=ICNT+1
    CALL ALQADF(A, IGEO, IBNL, IBNR, NPTS, DX, DLV, DRV, DLA, DRA, DIFC,
    1  SABS, XNF, XNT, XLAMO, X, B1, FLUXO, FLUOT)
    CALL TRIDG(A, NPTS, 101)
    DO 206 I=1, NPTS
206  FLUXN(I)=A(I,4)

    USE OF POWER METHOD TO FIND NEW LAMDA VALUE

    XNUM=0.0
    XDEN=0.0
    DO 150 I=1, NPTS
    XNUM=B1(I)*FLUXN(I)*B1(I)*FLUXO(I)+XNUM
150  XDEN=(B1(I)**2)*FLUXN(I)**2+XDEN
    XLAMN = XLAMO*XNUM/XDEN
    XKN=1./XLAMN

    CHECK FOR CONTINUED ITERATIONS

    ERROR = DABS((XKN-XKO)/XKN)
    XKO=1./XLAMN
    XLAMO = XLAMN
    FLXAV=0.0
    DO 210 I = 1, NPTS
    FLXAV=FLUXN(I)*(DLV(I)+DRV(I))+FLXAV
210  FLXAV=FLXAV/TOTVOL
    DO 211 I=1, NPTS
    FLUXO(I)=FLUXN(I)/FLXAV
211  FERM=0.0
    DO 400 I=1, NPTS
    FERR=DABS((FLUXP(I)-FLUXO(I))/FLUXO(I))
    IF(FERR.GT.FERRM) FERRM=FERR
400  CONTINUE
    CALL ALQADT(A, IGEO, IBNL, IBNR, NPTS, DX, DLV, DRV, DLA, DRA, DIFT,
    1  SABT, STR, X, FLUXO)

```

```

WRITE(6,500)
DO 745 I=1,NPTS
  WRITE(6,600)I,X(I),FLUXF(I),FLUXT(I)
745 CONTINUE
STOP
END
SUBROUTINE TRIDG(A,N,NDIM)
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION A(NDIM,4)
  DO 10 I=2,N
    A(I,1) = A(I,1)/A(I-1,2)
    A(I,2) = A(I,2)-A(I,1)*A(I-1,3)
    A(I,4) = A(I,4)-A(I,1)*A(I-1,4)
10 CONTINUE
NM1=N-1
A(N,4)=A(N,4)/A(N,2)
DO 20 J=1,NM1
  I=NM1+1-J
  A(I,4)=(A(I,4)-A(I,3)*A(I+1,4))/A(I,2)
20 CONTINUE
RETURN
END

```

```

SUBROUTINE ALOADF(A,IGEO,IBNL,IBNR,NPTS,DX,DLV,DRV,DLA,DRA,
1 DIFC,SABS,XNF,XNT,XLAMO,X,B1,FLUXO,FLUOT)
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION A(101,4),DX(101),DLV(101),DRV(101),DLA(101),DRA(101)
1 DIFC(101),SABS(101),XNF(101),XNT(101),X(101),B1(101),
2 FLUXO(101),FLUOT(101)

```

LOAD THE REQD MATRICES

```

PI=3.1415927
NPTS1=NPTS-1
DA1=1.0
DANPTS=1.0
IF(IGEO.EQ.2)DA1=0.0
IF(IGEO.EQ.3)DA1=0.0
IF(IGEO.EQ.2)DANPTS=2.0*PI**(NPTS)
IF(IGEO.EQ.3)DANPTS=4.*PI*X(NPTS)**2
A(1,1)=0.0
A(NPTS,3)=0.0
DO 105 L=2,NPTS1
  R=FLUOT(L)/FLUXO(L)
  A(L,2)=DIFC(L)/DX(L)*DLA(L)+DIFC(L+1)/DX(L+1)*DRA(L)+SABS(L)*
1 DLV(L)+SABS(L+1)*DRV(L)
  A(L,1)=-DIFC(L)/DX(L)*DLA(L)
  A(L,3)=-DIFC(L+1)/DX(L+1)*DRA(L)
  A(L,4)=XLAMO*((XNF(L)+XNT(L)*R)*DLV(L)+(XNF(L+1)+XNT(L+1)
1 *R)*DRV(L))
  B1(L)=A(L,4)/XLAMO
  A(L,4)=A(L,4)*FLUXO(L)
105 IF(IBNL.EQ.1) A(1,2)=DIFC(2)/DX(2)*DRA(1)+SABS(2)*DRV(1)
  IF(IBNL.EQ.1) A(1,3)=-DIFC(2)/DX(2)*DLA(2)
  IF(IBNR.EQ.0) A(NPTS,2)=1./(.71*3.)*DANPTS+DIFC(NPTS)/DX(NPTS)

```

IF (IBNR.EQ.0)A(NPTS,4)=STR(NPTS)\*DLV(NPTS)\*FLUXO(NPTS)  
IF (IBNR.EQ.1)A(NPTS,4)=STR(NPTS)\*DLV(NPTS)\*FLUXO(NPTS)  
RETURN  
END

NO OF REG	NO OF MAT	REL ERR	POWER(MEV)	MEV/FISSION
4	3	0.00010	1.00000	200.00000

\*\*\* CYLINDRICAL GEOMETRY \*\*\*

LEFT BC EQUALS ZERO CURRENT (SYMMETRY)  
RIGHT BC EQUALS ZERO FLUX AT EXTR. BNDRY (VACUUM)

REG NO	MAT NO	NO INT	THICK
1	2	7	0.70000D+01
2	1	20	0.12800D+02
3	2	3	0.30000D+01
4	3	20	0.25000D+02

CROSS SECTIONS FOR GROUP 1

MAT NO	DIF COEF	SIG ABS	SIG TRFR	NU SGFIS
1	0.84061D+00	0.16336D+00	0.18060D-01	0.67090D-01
2	0.55914D+00	0.12140D-01	0.23000D-01	0.00000D+00
3	0.59471D+00	0.38000D-03	0.92870D-01	0.00000D+00

CROSS SECTIONS FOR GROUP 2

MAT NO	DIF COEF	SIG ABS	SIG TRFR	NU SGFIS
1	0.38590D+00	0.10622D+00	0.00000D+00	0.10904D+01
2	0.30157D+00	0.18987D+00	0.00000D+00	0.00000D+00
3	0.49701D+00	0.67000D-03	0.00000D+00	0.00000D+00



1	0.00	0.213162E+15	0.110393E+14
2	1.00	0.216512E+15	0.112207E+14
3	0.00	0.226699E+15	0.117759E+14
4	3.00	0.244210E+15	0.127442E+14
5	4.00	0.269872E+15	0.142008E+14
6	5.00	0.304908E+15	0.162761E+14
7	6.00	0.350994E+15	0.191953E+14
8	7.00	0.410352E+15	0.233600E+14
9	7.64	0.437340E+15	0.256529E+14
10	8.28	0.460268E+15	0.274565E+14
11	8.92	0.478993E+15	0.288538E+14
12	9.56	0.493512E+15	0.299007E+14
13	10.20	0.503897E+15	0.306369E+14
14	10.84	0.510264E+15	0.310914E+14
15	11.48	0.512758E+15	0.312868E+14
16	12.12	0.511537E+15	0.312420E+14
17	12.76	0.506769E+15	0.309733E+14
18	13.40	0.498629E+15	0.304958E+14
19	14.04	0.487295E+15	0.298241E+14
20	14.68	0.472945E+15	0.289731E+14
21	15.32	0.455757E+15	0.279579E+14
22	15.96	0.435901E+15	0.267953E+14
23	16.60	0.413536E+15	0.255034E+14
24	17.24	0.38804E+15	0.241036E+14
25	17.88	0.361809E+15	0.226212E+14
26	18.52	0.332605E+15	0.210875E+14
27	19.16	0.301163E+15	0.195432E+14
28	19.80	0.267329E+15	0.180428E+14
29	20.80	0.194950E+15	0.169884E+14
30	21.80	0.137934E+15	0.208692E+14
31	22.80	0.919491E+14	0.334831E+14
32	24.05	0.548836E+14	0.442703E+14
33	25.30	0.328016E+14	0.486219E+14
34	26.55	0.196270E+14	0.492851E+14
35	27.80	0.117563E+14	0.478756E+14
36	29.05	0.704866E+13	0.453414E+14
37	30.30	0.422988E+13	0.422360E+14
38	31.55	0.254041E+13	0.388800E+14
39	32.80	0.152688E+13	0.354562E+14
40	34.05	0.918350E+12	0.320662E+14
41	35.30	0.552692E+12	0.287640E+14
42	36.55	0.332809E+12	0.255755E+14
43	37.80	0.200487E+12	0.225106E+14
44	39.05	0.120789E+12	0.195699E+14
45	40.30	0.727292E+11	0.167493E+14
46	41.55	0.436838E+11	0.140415E+14
47	42.80	0.260431E+11	0.114382E+14
48	44.05	0.151989E+11	0.893082E+13
49	45.30	0.832974E+10	0.651045E+13
50	46.55	0.366054E+10	0.416855E+13
51	47.80	0.242961E-13	0.185698E+13

## Appendix 1.2

### START Code

0.0.0.39,0.0.3.27  
0002.0.0.0.0.3.641  
11.15.132.75.45.49.82  
95.0.0.1264.2.621  
15.5.1786.2031.24.08  
2.44.0.0.0.0.11.44  
0.0.0.554.0.0.3.682  
003.0.0.0.0.4.437  
1433.2714.0.0.7.033  
2.24.0.0.0.13.04  
0053.8857.0.0.4.427  
009.0.0.0.0.5.569  
41.82.5146.0.0.44.83  
669.3.0.0.0.0.672.77  
0.0.0.20.0.0.2.0  
0001.0.0.0.0.0.5

CALCULATES THE REQUIRED TERMS FOR THE TWO GROUP EQUATIONS IN LAMARSH AFTER INPUTING THE COLLAPSED TWO GROUP CROSS SECTIONS SUPPLIED BY THE CODE ANASIN.

**VARIABLES:**

1 - FAST GROUP

D - DIFFUSION COEFFICIENT

**SIGA - MICROSCOPIC ABSORPTION CROSS SECTION**

**SIGS - MICROSCOPIC SLOWING DOWN CROSS SECTION**

SIGF\_\_\_\_\_ - MICROSCOPIC NU\*FISSION CROSS SECTION

SIGT - MICROSCOPIC TRANSPORT CROSS SECTION

**EA, ES, EF - MACROSCOPIC CROSS SECTIONS**

**O - OXYGEN**

**C - CARBON (GRAPHITE)**

## N - NITROGEN

**BE - BERYLLIUM**

**BO - BORON**

**25 - U-235**

**28 - U-238**

FE - IRON

**C - STAINLESS STEEL SHAFT AND VESSEL WALL**

**FUEL - FUEL**

**R - CONTROL DRUMS**

.....

ENSION DC(2),EAC(2),ESC(2),EFC(2),DF(2),EAF(2),ESF(2),

(2), DR(2), EAR(2), ESR(2), EFR(2), ETC(2), ETF(2), ETR(2),

AO(2), SIGA25(2), SIGA28(2), SIGAC(2), SIGAFE(2), SIGABE(2),

ABO(2), SIGSO(2), SIGS25(2), SIGS28(2), SIGSC(2), SIGSFE(2),

SBE(2), SIGSBO(2), SIGFO(2), SIGF25(2), SIGF28(2), SIGFC(2),

FFE(2), SIGFBE(2), SIGFBO(2), SIGTO(2), SIGT25(2), SIGT28(2)

TC(2), SIGTBE(2), SIGTBO(2), SIGAN(2), SIGSN(2),

FN(2), SIGTN(2)

```

N(UNIT=1, FILE='INSTART.DAT', STATUS='OLD')

```

```

N(UNIT#2,FILE='OUTSIART.DAT',STATUS='NEW')

```

```

AWFL=(235.04394*ENR)+(238.05082*(1.0-ENR))+14.0067
AWG=32.0+12.011
AWC=12.011
AWFE=55.85
AWBED=25.02
AWB4C=55.26

```

# CALCULATE THE NUMBER DENSITIES

```

AD25=((.6023*RHOFL)/AWFL)*VOLF*ENR
AD28=((.6023*RHOFL)/AWFL)*VOLF*(1.0-ENR)
ADNF=((.6023*RHOFL)/AWFL)*VOLF
ADOG=((.6023*RHOGL)/AWG)*VOLG*2.0
ADCG=((.6023*RHOGL)/AWG)*VOLG
ADC=((.6023*RHOGL)/AWC)*VOLC
ADFE=((.6023*RHOFE)/AWFE)
ADBE=((.6023*RHOBE)/AWBED)*VOLBE
ADB4=((.6023*RHOBE)/AWB4C)*VOLB4.0
ADBE=((.6023*RHOBE)/AWBED)*VOLBE
ADC84=((.6023*RHOBE)/AWB4C)*VOLB4

```

# CALCULATE AND PRINT MACROSCOPIC ABSORPTION AND NU-FISSION CROSS SECTIONS FOR THE U-235 AND U-238

```

EA251=AD25*SIGA25(1)
EA252=AD25*SIGA25(2)
EA281=AD28*SIGA28(1)
EA282=AD28*SIGA28(2)
EF251=AD25*SIGF25(1)
EF252=AD25*SIGF25(2)
EF281=AD28*SIGF28(1)
EF282=AD28*SIGF28(2)
WRITE(*,*) EA251,EA252,EA281,EA282
WRITE(*,*) EF251,EF252,EF281,EF282

```

# CALCULATE AND PRINT PARAMETERS NEEDED FOR THE LAMARSH TWO GROUP EQUATIONS

```

DO 100 I=1,2
EAF(I)=(AD25*SIGA25(I))+(AD28*SIGA28(I))+(ADNF*SIGAN(I))+
      (ADOG*SIGA(I))+(ADCG*SIGAC(I))+(ADC*SIGAC(I))
ESF(I)=(AD25*SIGS25(I))+(AD28*SIGS28(I))+(ADNF*SIGSN(I))+
      (ADOG*SIGS(I))+(ADCG*SIGSC(I))+(ADC*SIGSC(I))
ETF(I)=(AD25*SIGT25(I))+(AD28*SIGT28(I))+(ADNF*SIGTN(I))+
      (ADOG*SIGT(I))+(ADCG*SIGTC(I))+(ADC*SIGTC(I))
DF(I)=1.0/(3.0*ETF(I))
EFF(I)=(AD25*SIGF25(I))+(AD28*SIGF28(I))
100 CONTINUE

```

```

DO 200 I=1,2
EAC(I)=ADFE*SIGAFE(I)

```

9.6170716E-03	8.1939183E-02	0.1537438	2.4202248E-02
6.5076955E-02	1.090395	2.0145397E-03	0.0000000E+00
0.8406107	0.1633609	1.8055988E-02	6.7091495E-02
0.3859047	0.1062184	0.0000000E+00	1.090395
0.5591469	1.2146703E-02	2.3004992E-02	
0.3015706	0.1898717	0.0000000E+00	
0.5947114	3.8594622E-04	9.2896529E-02	
0.4970135	6.6994439E-04	0.0000000E+00	

## Appendix 1.3

### Sample Calculations

## FUEL BURNUP

critical radius = 12.8 cm (annular from 7 to 19.8 cm)

$$\pi (7)^2 (30) = 4618 \text{ cm}^3 \quad \pi (19.8)^2 (30) = 36,949 \text{ cm}^3$$

• volume of annular core =  $36,949 - 4618 = 32,331 \text{ cm}^3$

• cross sectional area of core (axial)

$$= \pi (19.8)^2 - \pi (7)^2 = 1231.6 \text{ cm}^2 - 153.9 \text{ cm}^2 = 1077.7 \text{ cm}^2$$

• volume of UN =  $(0.315)(32,331 \text{ cm}^3) = 10,184 \text{ cm}^3$  of UN

• mass of UN =  $10,184 \text{ cm}^3 \times 14.31 \text{ g/cm}^3 = 145.74 \text{ Kg}$  of UN

• mass of U-235 (5.6% enriched) at startup for criticality

5.6% U-235	.056 (235.04) =	13.1625
94.4% U-238	.944 (238.05) =	224.9560
N		<u>14.0067</u>
		252.13

$$\text{w/o U-235} = \frac{13.1625}{252.13} = .0522 = 5.2\%$$

$$(.0522)(145.74 \text{ Kg UN}) = 7.61 \text{ Kg U-235}$$

• mass of U-235 burned up during lifetime

$$1.05 (1 + \alpha) P \text{ g/day} = 1.23 \text{ g/day}$$

$$\alpha = 0.169$$

$$P = 1 \text{ MWt}$$

$$1.23 \text{ g/day} \times 365 \text{ day/yr} \times 7 \text{ yr} = 3.16 \text{ Kg}$$

• mass of U-235 needed for 7 year lifetime

$$7.61 \text{ Kg} + 3.16 \text{ Kg} = 10.77 \text{ Kg of U-235}$$

• new fuel enrichment

$$\frac{10.77 \text{ Kg}}{145.74 \text{ Kg}} = .0739 \%$$

<u>? (1)</u>	(235.0434) =	18.632
<u>? (2)</u>	(238.0506) =	219.160
N		<u>14.007</u>
		251.82

$$? (1) = .0793$$

$$? (2) = .9207$$

$$7.93\% \approx 8.0\% \text{ enriched fuel}$$

### POWER PRODUCTION

$$q''' = K \left[ (\Sigma_{f20} \phi + \Sigma_{f25} \phi)_{\text{fast}} + (\Sigma_{f20} \phi + \Sigma_{f25} \phi)_{\text{slow}} \right]$$

$$K = 200 \text{ MeV/fission}$$

$$\bar{\phi}_f = 6.7 \times 10^{15} \text{ n/cm}^2 \cdot \text{sec}$$

$$\bar{\phi}_s = 4.1 \times 10^{14} \text{ n/cm}^2 \cdot \text{sec}$$

$$\Sigma_{f20, \text{fast}} = 5.37 \times 10^{-5} \text{ cm}^{-1}$$

$$\Sigma_{f20, \text{slow}} = 0$$

$$\Sigma_{f25, \text{fast}} = 1.90 \times 10^{-3} \text{ cm}^{-1}$$

$$\Sigma_{f25, \text{slow}} = 3.32 \times 10^{-1} \text{ cm}^{-1}$$

$$q''' = 200 \left[ (3.596 \times 10^{11} + 1.2708 \times 10^{13}) + (0 + 1.3606 \times 10^{14}) \right]$$

$$= 200 \left[ (1.307 \times 10^{13}) + (1.361 \times 10^{14}) \right]$$

↑

9% thermal

↑

91% fast

$$q''' = 2.982 \times 10^{16} \frac{\text{MeV}}{\text{cm}^3 \cdot \text{sec}} \times \frac{1.6 \times 10^{-13} \text{ J}}{\text{MeV}} \times 32,331 \text{ cm}^3$$

$$= 1003 \text{ kW} = 1.003 \text{ MW}$$

### EQUIVALENT REACTOR RADIUS

volume of annular core = volume of cylindrical core

$$32,331 \text{ cm}^3 = \pi r^2 h$$

$$r^2 = \frac{32,331}{\pi h}$$

$$h = 30 \text{ cm}$$

$$r = \sqrt{\frac{32,331}{\pi h}}$$

$$r = 18.5 \text{ cm}$$



## REACTOR MASSES

use  $\text{mass} = \text{volume} \times \text{density}$

volumes based on  $\pi r^2 h$  for cylinder volumes

<u>material</u>	<u>densities</u>	<u>volumes</u>	<u>mass</u>
UN fuel	14.31 g/cm <sup>3</sup>	10,184 cm <sup>3</sup>	145.7 Kg
P <sub>4</sub> C clad	1.60 g/cm <sup>3</sup>	10,184 cm <sup>3</sup>	16.3 Kg
Stainless steel	7.86 g/cm <sup>3</sup>	16,663 cm <sup>3</sup>	131.0 Kg
CO <sub>2</sub> gas	0.0019 g/cm <sup>3</sup>	11,962 cm <sup>3</sup>	0.02 Kg
BeO (67%)	3.025 g/cm <sup>3</sup>	40,055 cm <sup>3</sup>	114.5 Kg
B <sub>4</sub> C (33%)	2.52 g/cm <sup>3</sup>		
Drive Assembly			114.5 Kg
TOTAL			522.0 Kg

## 2. REACTOR THERMAL-HYDRAULICS

The reactor thermal-hydraulics are incorporated into the code PEB, which was created specifically for this project. This code models the steady state pressure drop and temperature rise for a pebble bed reactor system in one dimension, using a rough finite difference technique. PEB assumes constant mass flow rate as input by the user. The four basic calculations performed by the code are: the pressure drop across the inlet and outlet grid plates, the pressure drop across the core, the fluid temperature rise across the core, and the fuel pebble centerline temperature. The code is written in standard fortran 77 and uses SI units. A code listing is given in Appendix 2.1.

### PROPERTY ASSUMPTIONS

The CO<sub>2</sub> coolant was assumed to be an ideal gas. This conclusion is based on the principle of corresponding states<sup>1</sup>. The principle of corresponding states predicts that a gas will behave ideally if its reduced temperature is approximately 2.0 and its reduced pressure is approximately 1.0. Since our reduced states are approximately correct, the ideal gas law will correctly model our system.

Four gas properties are needed for the PEB code; the density, the thermal conductivity, the constant pressure specific heat, and the viscosity. The density of the gas is found using the ideal gas relation,

$$\rho = \frac{P}{RT} \quad (2.1)$$

where P is the pressure, R is the gas constant, and T is the temperature. The thermal conductivity, constant pressure specific heat, and the viscosity were approximated as a linear function of

temperature from experimental data<sup>2</sup>. The three correlations used are shown in equations 2.2, 2.3, and 2.4.

$$K = 7.781 \times 10^{-5} T(^{\circ}\text{K}) - 6.59 \times 10^{-3} \text{ (W/m } ^{\circ}\text{K)} \quad (2.2)$$

$$C_p = 0.634 T(^{\circ}\text{K}) + 683.5 \text{ (J / kg } ^{\circ}\text{K)} \quad (2.3)$$

$$\mu = 3.783 \times 10^{-8} T(^{\circ}\text{K}) + 3.894 \times 10^{-6} \text{ (N s / m}^2\text{)} \quad (2.4)$$

The material property necessary for the code is the thermal conductivity. For our system, there are three materials: stainless steel 316, pyrolytic carbon, and uranium nitride. Equations 2.5, 2.6, and 2.7 show the relations used.

$$K_{\text{UN}} = 1.864 T^{0.361} \quad (2.5)^3$$

$$K_{\text{PyC}} = 4.256 - 0.0027 T \quad (2.6)^2$$

$$K_{\text{SS}} = 9.248 + 0.0157 T \quad (2.7)^2$$

## REACTOR INLET AND OUTLET GRID PRESSURE DROP

The pressure drop across the two grids was approximated by a correlation predicting the pressure drop across a grid with rounded edges<sup>4</sup>. The correlation predicts a head loss  $\Delta H$ ,

$$\Delta H = \frac{2 \zeta}{\rho \omega_1^2} \quad (2.8)$$

where  $\rho$  is the density of the fluid,  $\omega_1$  is the inlet velocity, and  $\zeta$  is given by equation 2.9,

$$\zeta = \frac{(\sqrt{\zeta'(1 - \bar{f})} + 1 - \bar{f})^2}{\bar{f}^2} \quad (2.9)$$

where  $\zeta'$  is a function of the porosity of the grid. Our calculations assumed a conservative constant value of 0.44 for  $\zeta'$ .  $\bar{f}$  is the flow area fraction, defined to be the ratio of the grate flow area divided by the inlet flow area, or for the grate in Figure 2.1,

$$\bar{f}_{\text{inlet}} = \frac{\pi}{4} \left( \frac{D}{P} \right)^2 \quad (2.10)$$

$$\bar{f}_{\text{outlet}} = \frac{\pi}{4\epsilon} \left( \frac{D}{P} \right)^2 \quad (2.11)$$

where  $\epsilon$  is the porosity of the particle bed. Using these correlations, the pressure drop across the two plates were calculated in the PEB code.

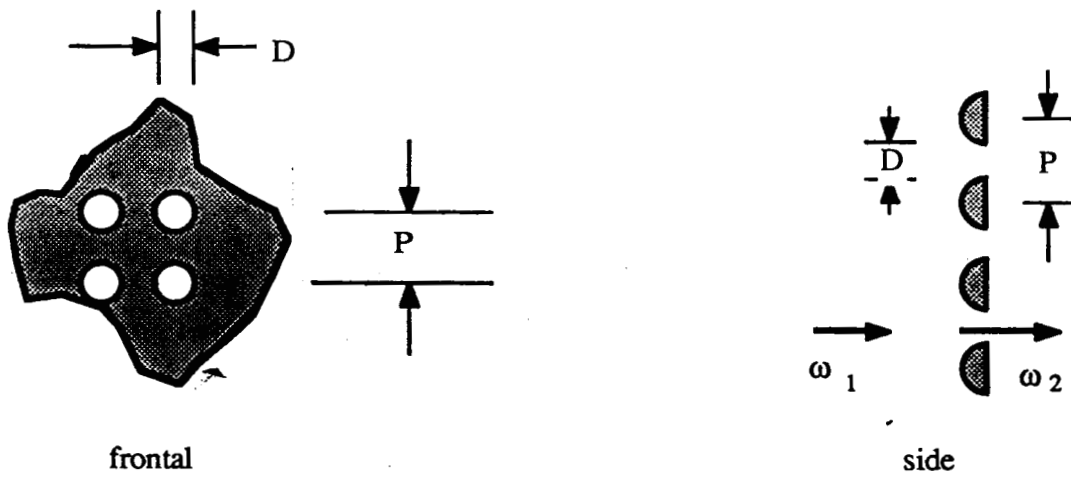


Figure 2.1 Sectional View of Core Grid Plates

## CORE PRESSURE DROP

The pressure drop across a cell,  $\Delta P$ , was approximated by the correlation<sup>5</sup>,

$$\Delta P = \psi \frac{H}{d_h} \frac{\rho}{2} u_p^2 \quad (2.12)$$

where  $H$  is the height of the bed,  $\rho$  is the density of the fluid, and  $u_p$  is the mean velocity of the gas in the gaps at that cell. The hydraulic diameter  $d_h$  is given by,

$$d_h = d \frac{\epsilon}{1 - \epsilon} \quad (2.13)$$

where  $d$  is the diameter of the individual pebble and  $\epsilon$  is the porosity of the bed. The function,  $\psi$ , is a function of the Reynolds number,  $Re$  and  $\epsilon$ , as is shown in equation 2.14.

$$\psi = 320 \left( \frac{Re}{1 - \epsilon} \right)^{-1} + 6 \left( \frac{Re}{1 - \epsilon} \right)^{-0.1} \quad (2.14)$$

## FLUID TEMPERATURE RISE

The heat transfer through the core was found by applying conservation of energy in steady state. The form of the energy conservation equation, assuming  $q'_s$  is the linear heat generation rate, is given in equation 2.15.

$$q'_s = \dot{m} c_p \text{Area} (T_{out} - T_{in}) \quad (2.15)$$

The PEB code assumes that  $c_p$  is a constant at  $T_{in}$  and divides the core region into a set number of cells, as specified by the user. The total heat produced in the core is input by the user, then PEB subdivides the total power to all of the cells in a sinusoidal pattern. The linear heat source,  $q'$ , is

assumed to be a sinusoidal source, with zero power produced at the actual edge of the reactor. Although this is not entirely correct, it does well enough for our calculations, since the total amount of energy deposited in the fluid will be the same. Also, this source shape will overestimate the fuel temperature in the center of the reactor, which is a safety plus. With the inlet temperature and the heat generation rate known for the cell, the cell outlet temperature can be found using equation 2.15. The PEB code then uses the outlet temperature of cell  $i$  to be the inlet temperature of cell  $i+1$ , thus it is a rough finite difference code.

## FUEL CENTER TEMPERATURE

To find the temperature in the center of the individual pebbles,  $T_c$ , three basic thermal resistances are calculated, the thermal resistance across the convection surface, the thermal resistance across the layers of cladding, and the thermal resistance through the fuel itself. The center fuel temperature is then found by using the equation 2.16, where  $Q$  is the heat generated in one pebble.

$$T_c = T_{bulk} + Q(R_{conv} + R_{clad} + R_{fuel}) \quad (2.16)$$

The thermal resistance across the convection surface can be found using the Nusselt number correlation for a packed sphere bed<sup>5</sup>,

$$Nu_d = \frac{hD}{k_c} = f_\epsilon Nu_s \quad (2.17)$$

where  $h$  is the heat transfer coefficient,  $D$  is the diameter of an individual pellet,  $k_c$  is the conductivity of the gas,  $f_\epsilon$  is the arrangement factor, and  $Nu_s$  is the Nusselt number for a single ball. The arrangement factor,  $f_\epsilon$ , is given by equation 2.18.

$$f_{\epsilon} = 1 + 1.5 (1 - \epsilon) \quad (2.18)$$

The Nusselt number for a single sphere is given in terms of a laminar and a turbulent Nusselt number,

$$Nu_s = 2 + \sqrt{Nu_l^2 + Nu_t^2} \quad (2.19)$$

The turbulent and laminar Nusselt number are in turn given by the following empirical correlations,

$$Nu_l = 0.66 \left( \frac{Re}{\epsilon} \right)^{0.5} (Pr)^{0.33} \quad (2.20)$$

$$Nu_t = \frac{0.037 \left( \frac{Re}{\epsilon} \right)^{0.8} Pr}{1 + 2.443 \left( \frac{Re}{\epsilon} \right)^{-0.1} (Pr^{0.667} - 1)} \quad (2.21)$$

With the Nusselt number known for the pebble, the heat transfer coefficient,  $h$ , can be determined from the Nusselt number correlation in equation 2.22.

$$h = \frac{Nu k_c}{D} \quad (2.22)$$

Since the thermal resistance through the cladding then will be equal to,

$$R_{conv} = \frac{T_{surface} - T_{bulk}}{Q} \quad (2.23)$$

and

$$h = \frac{q''}{T_{surface} - T_{bulk}} \quad (2.24)$$

then,

$$R_{\text{conv}} = \frac{1}{\text{Nu } k_c \pi D} \quad (2.25)$$

where  $r$  is the radius of the pebble. The resistance through  $N$  layers of cladding will be equal to,

$$R_{\text{clad}} = \frac{\sum_{n=1}^{N-2} \frac{1}{k_{n+1}} \left( \frac{1}{r_n} - \frac{1}{r_{n+1}} \right)}{4\pi} \quad (2.26)$$

where the subscripts are shown in figure 2.2.

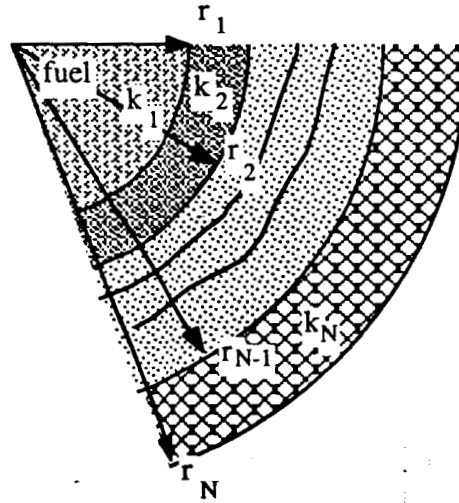


Figure 2.2 Cutaway of fuel pellet showing node numbering

The thermal resistance through the fuel itself can be calculated using the 1-D conduction equation with a source term. The solution to this is given in equation 2.27.

$$R_{\text{fuel}} = \frac{1}{8 k_1 r_1 \pi} \quad (2.27)$$



Using equation 2.16, 2.25, 2.26, and 2.27, the final solution for the centerline fuel temperature will be,

$$T_{\text{center}} = T_{\text{bulk}} + Q \left( \frac{1}{\text{Nu } k_c \pi D} + \frac{\sum_{n=1}^{N-2} \frac{1}{k_{n+1}} \left( \frac{1}{r_n} - \frac{1}{r_{n+1}} \right)}{4\pi} + \frac{1}{8 k_1 r_1 \pi} \right) \quad (2.28)$$

In axial flow pebble beds, there is a significant amount of radial conduction and radiation that will occur. This phenomenon was not accounted for in our calculations, due to the complexity of the problem. However, the radial heat transfer regime will only serve to "spread out" the heat source, thus our calculations are still conservative in that our temperature distribution will appear to be more peaked in the radial direction.

#### PEB INPUT AND OUTPUT DESCRIPTION

The input to the PEB code, shown in Table 2.1, is quite simple. Most of the input is self explanatory, except the material number. This number corresponds to an identification number contained in PEB. 50 is pyrolytic carbon, 61 is SS-316, and 82 is uranium nitride. Also, the aperture of the core plate refers to the ratio of D to P in Figure 2.1. One of the actual input decks to PEB is contained in Appendix 2.2.

Table 2.1 PEB Input description			
line #			
1	# of cells (NAX)	# of layers in a pellet (NLAY)	total reactor length
2	inlet temperature	inlet pressure	mass flow rate
3	total reactor power		

4	inner radius of core region	outer radius of core region	porosity of bed	aperture of core plate
5	radius of fuel	fuel material number		
6	radius of first clad	first clad material number		
7	radius of second clad	second clad material number		
.	.	.		
.	.	.		
.	.	.		
NLAY+4	radius of last clad	last clad material number		

An example PEB output is given in Appendix 2.3. The output is clearly labelled, therefore no explanation will be given.

## RESULTS

The thermal-hydraulic calculations in PEB were carried out to determine the size of the reactor, the size of the pebbles, the aperture in the core grid plate, and the mass flow rate through the core. All calculations were performed with the limiting parameters being that the fluid outlet temperature could not exceed 900 °K and the fuel centerline temperature could not exceed 1800 °K. With these limiting margins, no serious interactions or material degradation will occur. In the determination of the entire core size, a corroborative effort with the neutronics personnel had to be maintained, since the overall core size affects both fields. In the end, however, there was a large margin in which to work with and thus no severe problems were encountered. The size of the individual pebbles was strictly a thermal-hydraulics problem. If the pebbles are too large, the center temperature will exceed the design criterion. Also, if the pebbles are too small, there will be the possibility of significant manufacturing problems and too large of a pressure drop. Since we did not know the extent of the manufacturing problems, the centerline fuel temperature was raised to a maximum. The massflow rate was adjusted to obtain an outlet temperature of approximately 850 °K. Previous experiments<sup>6</sup> have shown that the porosity of the particle bed will be 0.37, if the ratio

of the core radius to the pebble radius is greater than 5. Since our reactor radius is on the order of 10 cm and the pebbles are less than a centimeter in size, we could safely assume that the porosity is a constant 0.37. The inlet temperature was assumed to be 500 °K and the inlet pressure was assumed to be 6.9 MPa.

After several iterations, the pebble size was found to be 0.6 cm in diameter. The mass flow rate was 2.43 kg/s. The core size was found by the neutronics personnel. The final actual input conditions are the ones listed in Appendix 2.2. With these input conditions, the output in Appendix 2.3 was generated. These results are displayed graphically in Figures 2.3, 2.4, 2.5, and 2.6.

Figure 2.3 shows the total pressure drop across the reactor. Notice that most of the pressure losses are incurred at the inlet and outlet grid plate. Since this does not make intuitive sense, there is a significant possibility that the correlation in equation 2.12 is either incorrect or is used inappropriately. First, this is probably not the exact plate geometry that will be used, however this should not be a large factor. The one factor that was overassumed was the  $\zeta'$  factor. However, in the end result, the pressure drop was assumed to be that of Figure 2.3, since this will be conservative.

Figure 2.4 shows the fluid temperature rise through the core. The outlet temperature was found to be about 868 °K, which is within our design criterion. Figure 2.5 shows the fuel center temperature as a function of axial location. Two curves are shown, the uppermost being the hot channel temperature distribution and the lower one being the average channel temperature distribution. Both curves have a maximum of less than 1400 °K, which is well below the design criterion. Figure 2.6 shows two curves, the average center fuel temperature and the fluid temperature as a function of core axial distance. The calculations came out quite well, since the two curves had nearly the same value at the outlet.

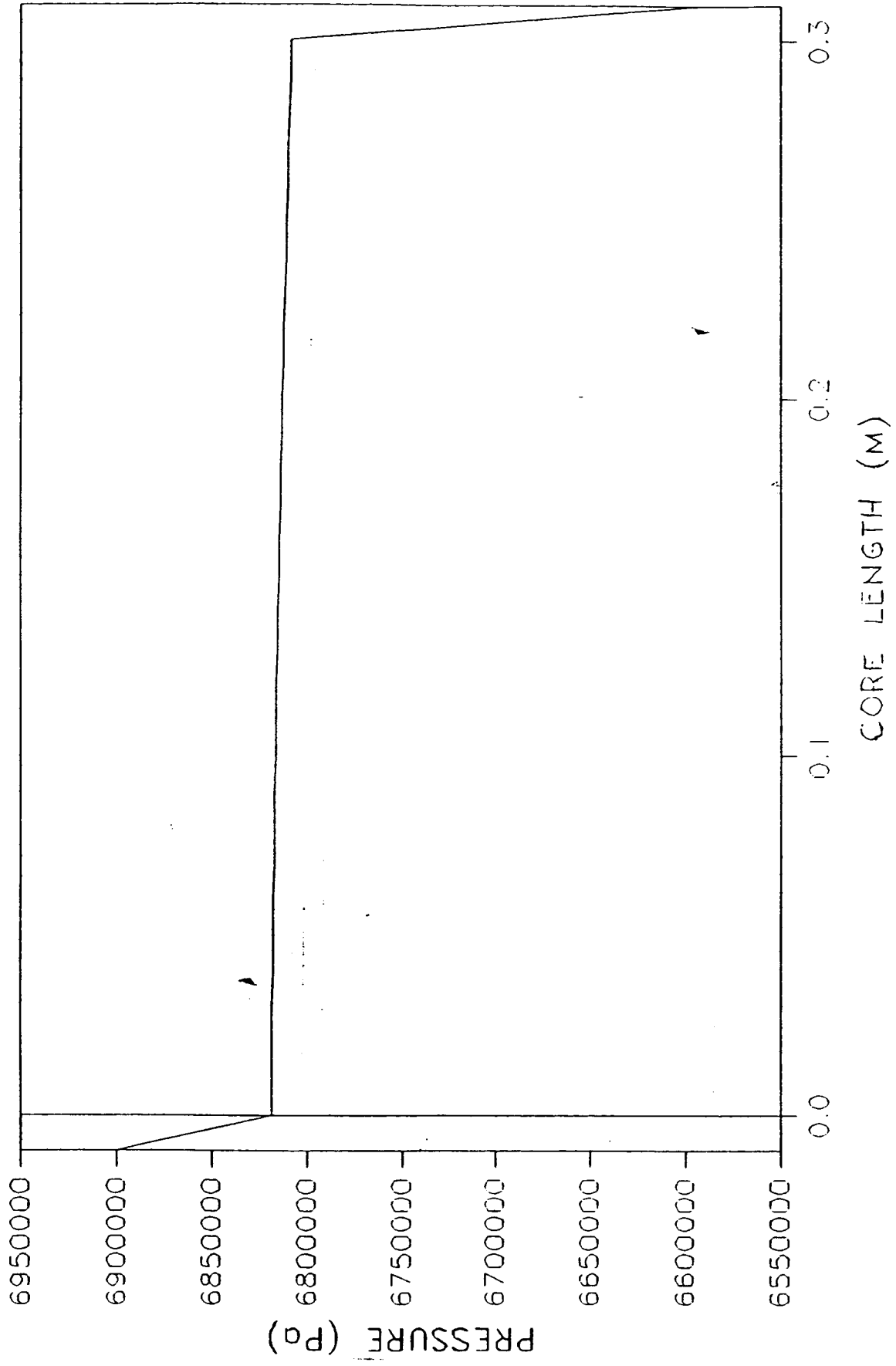


Figure 2.3 Pressure Drop Across the Core

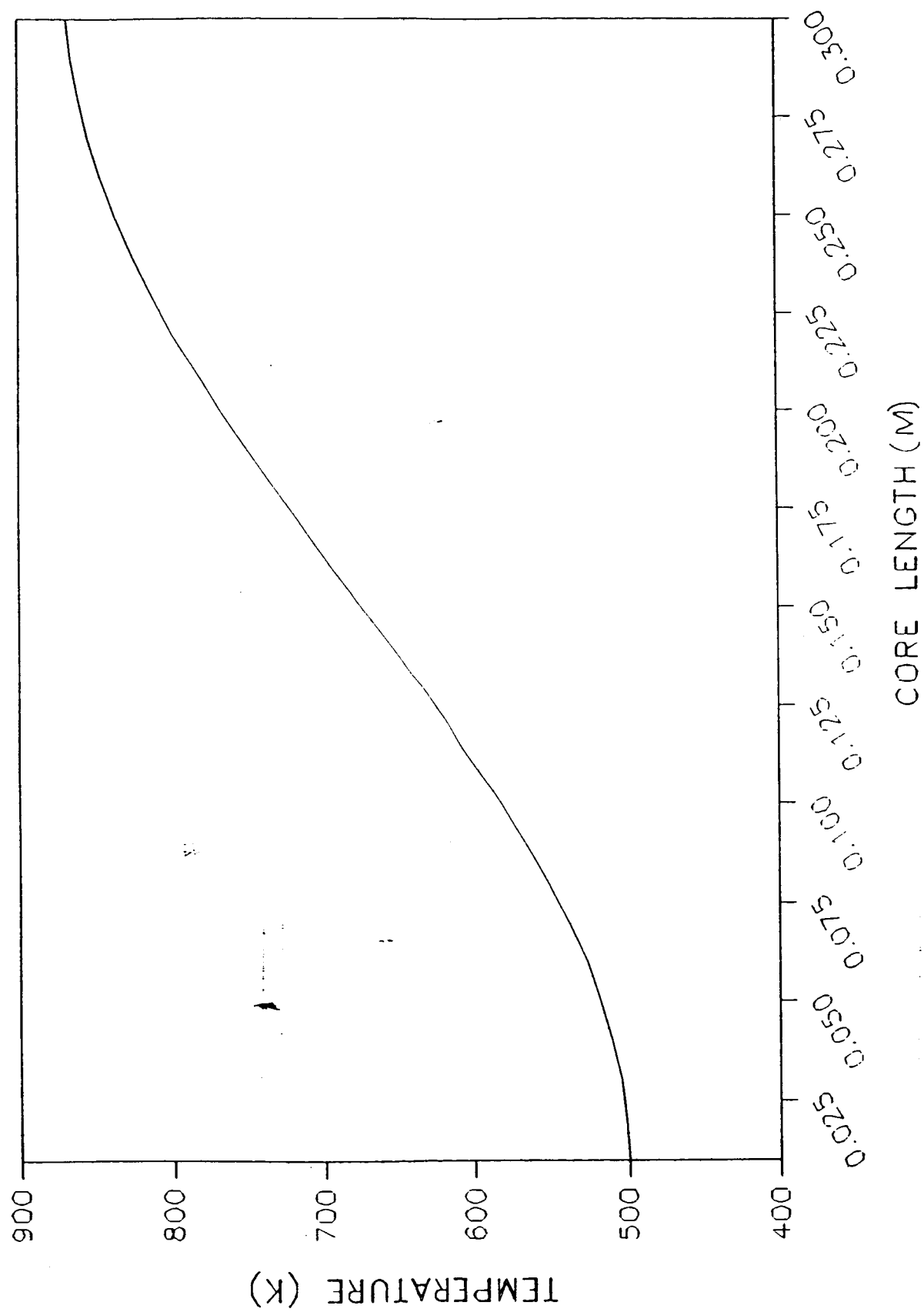


Figure 2.4 Fluid Temperature Rise Across the Core

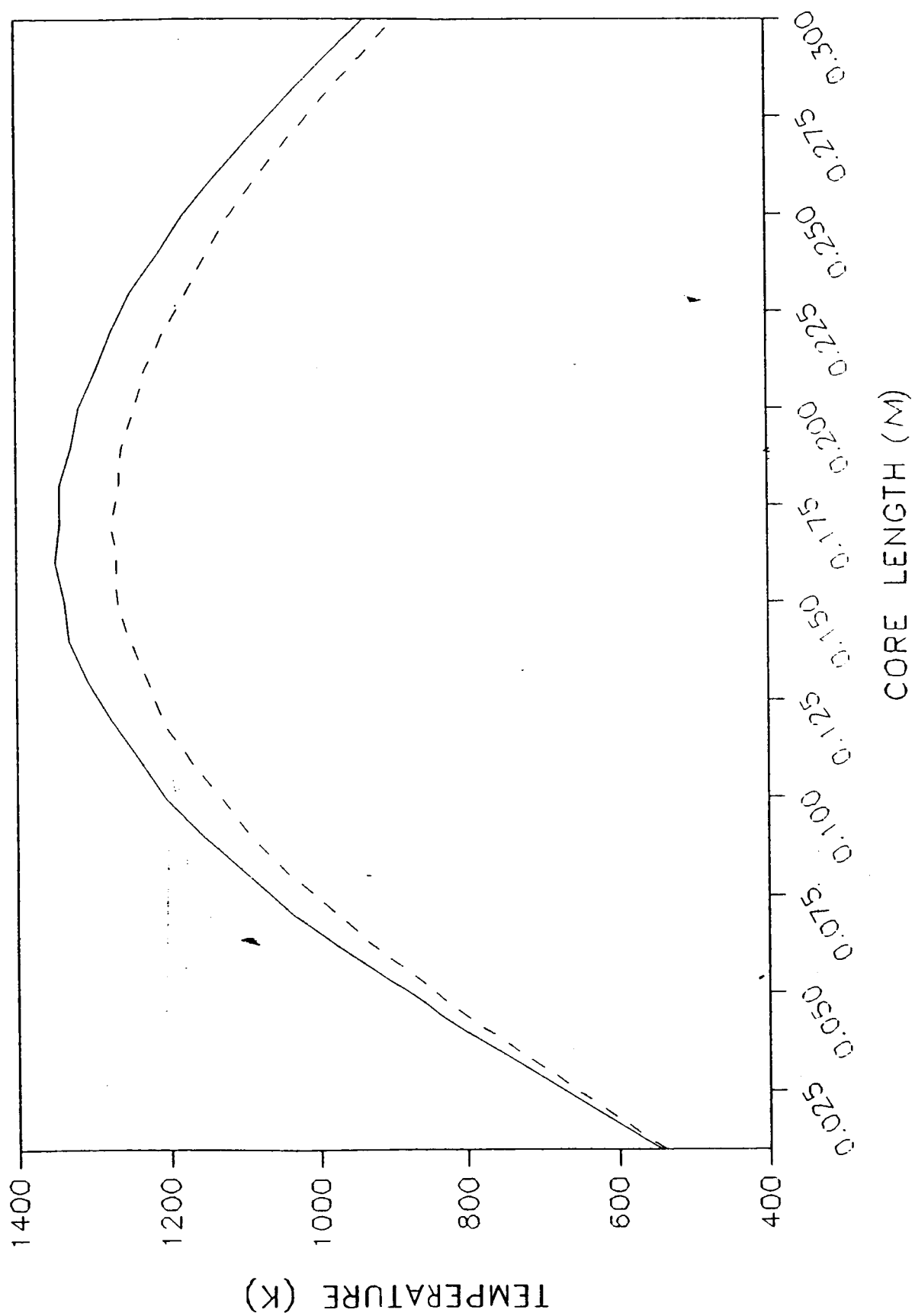


Figure 2.5 Maximum and Average Fuel Center Temperature

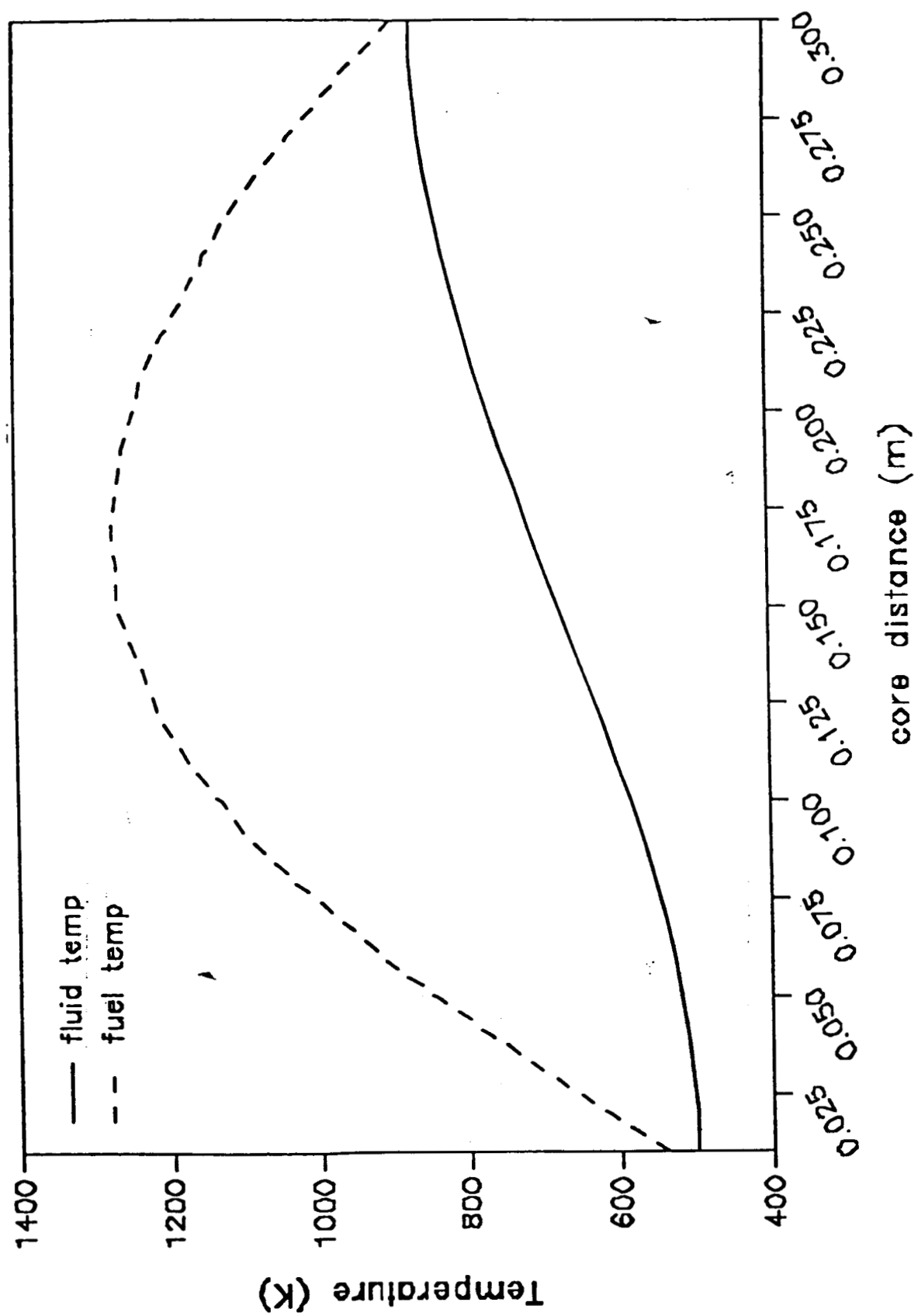


Figure 2.6 Fluid and Fuel Center Temperature Across the Core

## REFERENCES

1. Wark, K., "Thermodynamics," McGraw-Hill Book Company, New York, p 137, 1983.
2. Incropera, F. P. and DeWitt, D. P., "Fundamentals of Heat and Mass Transfer," John Wiley & Sons, New York, 1985, Appendix A.
3. Thomas, J. K. and Hayes, S. L., "Material Property and Irradiation Performance Correlations for Nitride Fuels," Fifth Symposium for Space Nuclear Power Systems, Albuquerque, NM, 1988.
4. Chik, I. E. Idel, "Handbook of Hydraulic Resistance, Coefficients of Local Resistance, and of Friction," Gosudarstvennoe Energeticheskoe Izdatel'stvo, Moskva - Leningrad, 1960.
5. French, H. "Heat Transfer and Fluid Flow in Nuclear Systems," Pergamon Press, New York, 1987, Chapter 5, part 2.
6. S. T. Robinson, "The Pebble Bed Reactor", Proceedings of the Symposium on Gas Cooled Reactors, Journal of the Franklin Institute, Lancaster, Pa, p. 94, 1960.



## **Appendix 2.1 PEB Code Listing**



C 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789

C SUBROUTINE PLATEPD(P1,P2,T1,K)

C     PARAMETER MAXCEL=100  
C     PARAMETER PI=3.14

C     COMMON/THERMAL/ TIN,PIN,VELIN,TEMP(MAXCEL,2),PRESS(MAXCEL),  
C     &VEL(MAXCEL),QSDUR(MAXCEL),RMASFLO  
C     COMMON/GEOM/ DIA,NAX,NLAY,PORIS,CLEN,XAREA,RADIUS(MAXCEL),  
C     &MID(MAXCEL),APTUR

C     IF(K.EQ.1)FACT=PORIS  
C     IF(K.EQ.2)FACT=1.0

C     ZETAP = 0.44  
C     AREAR = PI\*APTUR\*\*2.0/(4.0\*FACT)  
C     QUO = (RMASFLO/(FACT\*XAREA))\*\*2.0\*DENS(T1,P1)  
C     DP = 0.5\*QUO\*((ZETAP\*(1.0-AREAR))\*0.5+(1.0-AREAR))/AREAR\*\*  
C     &     2.0

C     P2=PI-DP

C     RETURN  
C     END

C 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789

C SUBROUTINE PRESR(N)

C     PARAMETER R=188.9  
C     PARAMETER PI=3.14  
C     PARAMETER MAXCEL=100

C DATA

C     COMMON/THERMAL/ TIN,PIN,VELIN,TEMP(MAXCEL,2),PRESS(MAXCEL),  
C     &VEL(MAXCEL),QSDUR(MAXCEL),RMASFLO  
C     COMMON/GEOM/ DIA,NAX,NLAY,PORIS,CLEN,XAREA,RADIUS(MAXCEL),  
C     &MID(MAXCEL),APTUR

C     FACT1=RYNOLD(N-1)/(1-PORIS)  
C     PSI=320.0\*FACT1\*\*-1.0+6.0\*FACT1\*-0.1  
C     DP = (PSI\*CLEN\*(1.0-PORIS)\*RMASFLO\*\*2.0)/(2.0\*  
C     &     RADIUS(NLAY)\*PORIS\*\*3.0\*DENS(TEMP(N,2),PRESS(N-1))\*XAREA\*\*2.0)  
C     PRESS(N)=PRESS(N-1)-DP  
C     RETURN  
C     END

C 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789 123456789

C SUBROUTINE TFLUID(N)

C     PARAMETER R=188.9  
C     PARAMETER PI=3.14  
C     PARAMETER MAXCEL=100

C DATA

C     COMMON/THERMAL/ TIN,PIN,VELIN,TEMP(MAXCEL,2),PRESS(MAXCEL),  
C     &VEL(MAXCEL),QSDUR(MAXCEL),RMASFLO

**NUSLAM=0.66\*FACT1\*\*0.5\*FACT2\*\*0.333**

$$\text{NUSS}' = 2.0 + \text{SQRT}(\text{NUSLAM} \cdot 2.0 + \text{NUSTURB} \cdot 2.0)$$
$$\text{NUSSLT} = (1.0 + 1.5 * (1 - \text{PORIS})) * \text{NUSS}$$

## RETURN

C23456789 123456789 123456789 123456789 123456789

## FUNCTION PRANDL(N)

**PARAMETER R=188.9**

**PARAMETER PI=3.14**

PARAMETER MAXCEL = 100

## DATA

COMMON/THERMAL/ TIN,PIN,VELIN,TEMP(MAXCEL,2),PRESS(MAXCEL),

 $\&VEL(MAXCEL),QSOUR(MAXCEL),RMA SFLO$ 

COMMON/GEOM/ DIA,NAX,NLAY,PORIS,CLEN,CLEN,XAREA,RADIUS(MAXCEL),

**&MID(MAXCEL), APTUR**

$$\text{PRANDL} = \text{SPHEAT}(\text{TEMP}(\text{N}, 2)) * \text{VISC}(\text{TEMP}(\text{N}, 2)) / \text{TCOND}(\text{TEMP}(\text{N}, 2))$$

## RETURN

C23456789 123456789 123456789 123456789 123456789 123456789

# FUNCTION RYNOLD(N)

PARAMETER R=188.9

**PARAMETER PI=3.14**

**PARAMETER MAXCEL=100**

**DATA:**

COMMON/THERMAL/ TIN,PIN,VELIN,TEMP(MAXCEL,2),PRESS(MAXCEL),

$$\text{RVEI}(\text{MAXCEI}) \quad \text{OSOUR}(\text{MAXCEI}) \quad \text{RMASEI} \quad 0$$

COMMON/GEOM/ DIA.NAX.NLAY.PORIS.CLEN.CLEN.XAREA.RADIUS(MAXCEL).

**BMID(MAXCEL) - APTUR**

$$RYNOLD=RMAFLO*RADIUS(NLAY)/(XAREA*PORIS*VISC(TEMP(N,2)))$$

## RETURN

223456789 123456789 123456789 123456789 123456789 123456789 123456789

## FUNCTION VISC(T)

FROM INCRYPTERA & DEWITT

$$VISC = 3.783E-8 * T + 3.894E-6$$

## RETURN

## **Appendix 2.2 PEB Code Input**

30 3 0.3  
500.0 6.9E6 2.43  
1.0E6  
0.07 0.196 0.375 0.6  
0.00312 82  
0.0039 50  
0.004 61

Form: UNGB01 10/21/05/85

## **Appendix 2.1 PEB Code Output**

NAX NLAY RLEN

30 3 0.3000000

TIN PIN MASFLO

500.0000 6900000. 2.430000

QTOT

1000000.

RAD1 RAD2 PORIS APTUR

7.000000E-02 0.1960000 0.3750000

CROSS SECTIONAL AREA 0.1052402

FUEL PELLET DATA

RADIUS MID

3.120000E-03 82 3.8999999E-03

PRESSURE OUT OF INLET = 6819191.

RMASFLO

2.4300000

POS TFUEL TFLUID PRESSURE DENS

1.000000E-02	538.8765	500.0000	
2.000000E-02	615.8722	501.1426	6819191.
3.000000E-02	692.4456	504.5212	6818941.
4.000000E-02	770.9007	510.0877	6818689.
5.000000E-02	842.2253	517.7615	6818434.
6.000000E-02	914.4991	527.4313	6818175.
7.000000E-02	974.9374	538.9583	6817911.
8.000000E-02	1036.378	552.1797	6817641.
9.000000E-02	1091.851	566.9126	6817364.
0.1000000	1131.411	582.9578	6817078.
0.1100000	1173.516	600.1045	6816784.
0.1200000	1209.238	618.1339	6816481.
0.1300000	1228.035	636.8229	6816167.
0.1400000	1251.129	655.9481	6815843.
0.1500000	1268.404	675.2880	6815509.
0.1600000	1269.245	694.6263	6815163.
0.1700000	1275.444	713.7536	6814807.
0.1800000	1266.192	732.4696	6814439.
0.1900000	1262.288	750.5845	6814061.
0.2000000	1244.129	767.9199	6813673.
0.2100000	1230.816	784.3102	6813275.
0.2200000	1204.700	799.6031	6812868.
0.2300000	1175.313	813.6599	6812451.
0.2400001	1149.172	826.3564	6812027.
0.2500001	1117.865	837.5826	6811595.
0.2600001	1078.105	847.2433	6811156.
0.2700000	1035.755	855.2584	6810712.
0.2800000	991.0170	861.5621	6810263.
0.2900000	944.0694	866.1039	6809811.
0.3000000	895.6300	868.8483	6809355.
PRESSURE OUT OF CORE =	6582291.		6808898.

50 4.000000E-03

61

6819191.	72.19894
6818941.	72.03169
6818689.	71.54668
6818434.	70.76326
6818175.	69.71183
6817911.	68.43110
6817641.	66.96487
6817364.	65.35880
6817078.	63.65759
6816784.	61.90282
6816481.	60.13141
6816167.	58.37484
6815843.	56.65900
6815509.	55.00433
6815163.	53.42632
6814807.	51.93622
6814439.	50.54170
6814061.	49.24753
6813673.	48.05624
6813275.	46.96865
6812868.	45.98436
6812451.	45.10212
6812027.	44.32017
6811595.	43.63645
6811156.	43.04882
6810712.	42.55518
6810263.	42.15359
6809811.	41.84240
6809355.	41.62019
6808898.	41.48594



### 3. BRAYTON CYCLE

#### INTRODUCTION

The Brayton cycle utilizes super-heated vapor throughout the cycle, the fluid does not boil and does not operate in the liquid-vapor dome. This requires that turbomachinery have high component efficiency to compensate for the work of compression that is introduced. An actual Brayton cycle temperature-entropy diagram is shown in Figure 3.1, for the system shown in Figure 3.2. The direct-closed cycle uses the primary working fluid for the entire cycle and circulates the same gas repeatedly. Energy is added to the gas in the reactor, the gas is then expanded through the turbine, the waste heat is rejected by the heat pipe radiator, and finally compressed by the compressor. The reactor, turbine, and heat pipe radiator all have significant pressure losses in the Brayton cycle due to the working fluid being gaseous. Also the pressure losses within the ducting may be significant for the same reason.

Carbon dioxide was chosen as a working fluid for this particular design for several reasons. First, the mission is proposed to spend most of its time on the Martian surface and the primary constituent of the Martian atmosphere is carbon dioxide. Therefore, in case of a minor loss of coolant there would be an abundant supply to refurbish the system. Also, in case of a major loss of coolant, the reactor would be flooded with atmospheric gas, in this case the adverse affects would be minimized by the fact that carbon dioxide is both the coolant and atmospheric fluid. Carbon dioxide is also an extremely suitable fluid for material, thermal, and nuclear designs because of inertness, good heat transfer characteristics, and low absorption cross section. Finally, the departure of carbon dioxide from perfect gas laws results in less work of compression, this increases the cycle efficiency (1).

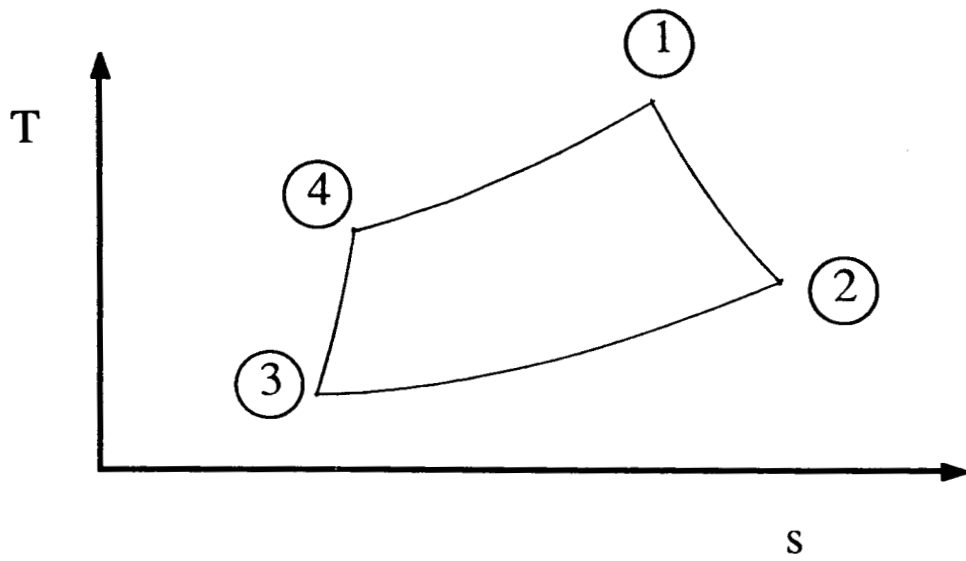


Figure 3.1 Brayton Cycle Temperature-Entropy Diagram

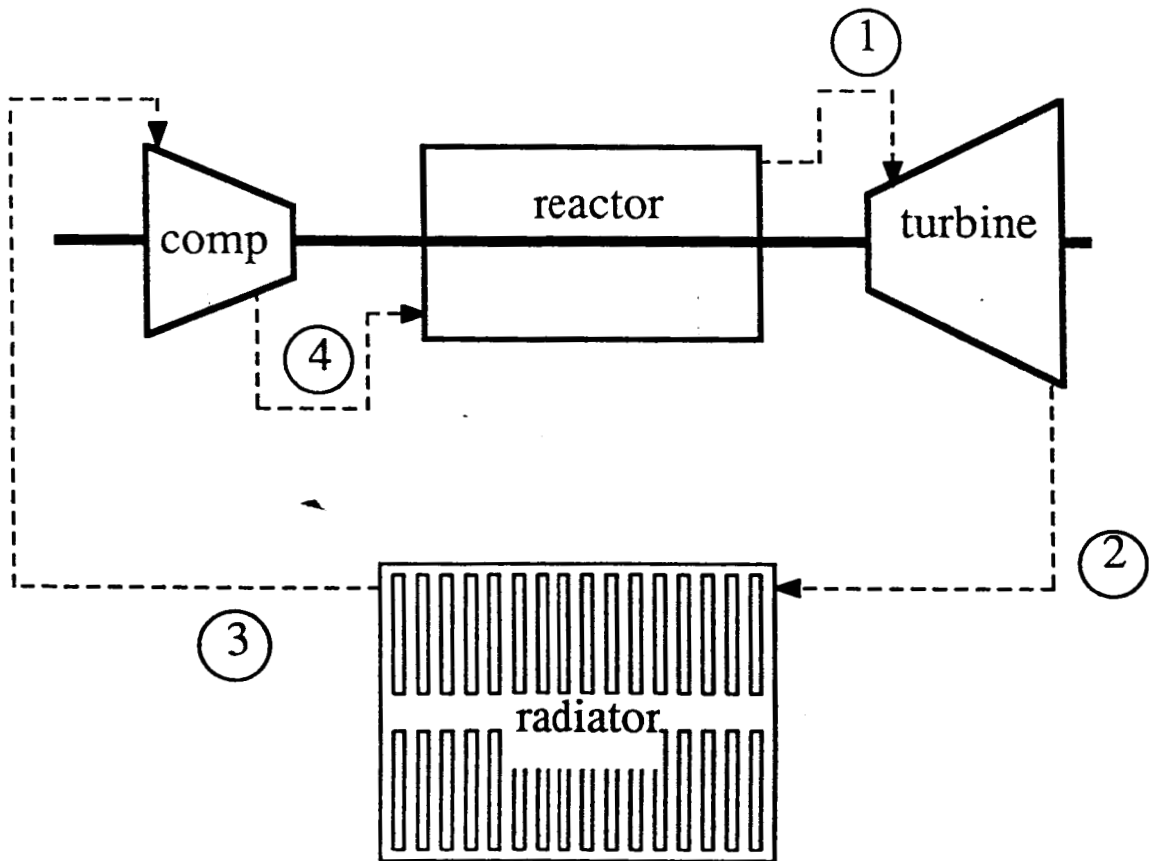


Figure 3.2 System Brayton Cycle

## THEORY

In order to evaluate the performance of a Brayton cycle the standard overall cycle efficiency,  $\eta$ , must be determined from Equation 3.1,

$$\eta = \frac{W_{\text{net}}}{Q_{\text{in}}} \quad (3.1)$$

where  $Q_{\text{in}}$  is the heat added to the system by the reactor and  $W_{\text{net}}$  is the net work done by the system and is calculated by Equation 3.2:

$$W_{\text{net}} = W_T - W_C \quad (3.2)$$

where  $W_T$  is the work produced by the turbine and  $W_C$  is the work necessary to run the compressor. The turbomachinery is not ideal and therefore each component's performance is determined by the efficiency of that component. The turbine efficiency ( $\eta_T$ ) is represented by Equation 3.3:

$$\eta_T = \frac{\text{Actual Work}}{\text{Ideal Work}} \quad (3.3)$$

whereas, the compressor efficiency ( $\eta_C$ ) is the inverse as in Equation 3.4:

$$\eta_C = \frac{\text{Ideal Work}}{\text{Actual Work}} \quad (3.4)$$

These efficiencies for the turbomachinery are given for the given working fluid and thermodynamic conditions and thus by calculating the ideal work, the actual work can be determined. Equation 3.5 is used to calculate the actual work necessary to drive the compressor ( $W_C$ ):

$$W_C = \frac{c_p T_3}{\gamma_c} (r_p^{(\gamma-1)/\gamma} - 1) \quad (3.5)$$

where  $c_p$  is the specific heat at constant pressure,  $T_3$  is the compressor inlet temperature,  $\gamma$  is the ratio of specific heat at constant pressure to that at constant volume, and  $r_p$  is the compressor pressure ratio. The work that the turbine and is defined by Equation 3.6:

$$W_T = \gamma c_p T_1 \left[ 1 - \frac{\beta}{r_p^{(\gamma-1)/\gamma}} \right] \quad (3.6)$$

where  $T_1$  is the reactor outlet temperature, and  $\beta$  is the pressure losses between the compressor and turbine and is defined by Equation 3.7:

$$\beta = \left( \frac{P_4}{P_1} \cdot \frac{P_2}{P_3} \right)^{(\gamma-1)/\gamma} \quad (3.7)$$

In Equation 3.7 the pressures are at state 1, the reactor outlet, state 2, the turbine outlet, and states 3 and 4, the inlet and outlet of the compressor.

Once the net work is determined by subtracting Equation 3.5 from 3.6, the heat input by the reactor ( $Q_{in}$ ) is the last necessary calculation for determining the cycle efficiency. Equation 8 is used to calculate  $Q_{in}$  :

$$Q_{in} = c_p T_1 \left[ \left( \frac{T_1}{T_3} - 1 \right) - \frac{1}{\gamma_c} (r_p^{(\gamma-1)/\gamma} - 1) \right] \quad (3.8)$$

where the states are described above.

The work and heat input and output can also be calculated by  $c_p \Delta T$  across each component if all of the temperatures around the cycle are known and the gas displays ideal characteristics throughout:

$$Q_{in} = c_p (T_1 - T_4) \quad (3.9)$$

$$Q_{out} = c_p (T_2 - T_3) \quad (3.10)$$

$$W_T = c_p (T_1 - T_2) \quad (3.11)$$

$$W_C = c_p (T_4 - T_3) \quad (3.12)$$

where  $Q_{out}$  is the heat removed from the system by the heat pipe radiator (2).

## RESULTS

In order to evaluate the closed-direct Brayton cycle several assumptions were made. It was assumed that the pressure losses in the ducting would be insignificant compared to the losses across the components, that the working fluid performs ideally throughout the system, that the specific heat at constant pressure remained constant around the cycle, and that carbon dioxide is closely approximated by argon in the turbomachinery. The last assumption allows for the use of  $\eta_C = \eta_T = 0.86$  (3), while the third assumption results in the use of  $r_p = 1.22$  kJ/kg-K and  $\gamma = 1.290$  for the entire cycle analysis. The last assumption also allows for the use of  $r_p = 1.90$ , and a turbine pressure ratio of 1.75 (4).

After defining the reactor inlet and outlet states and considering the above assumptions the pressures at the turbine outlet and compressor inlet were found using the turbine and compressor ratios. Once all of the pressures are known  $\beta$  can be calculated and was found to be  $\beta = 1.0188$ . Using Equation 3.6 the turbine work was calculated to be  $W_T = 513$  kW, then rearranging Equation 3.11 the turbine outlet temperature was determined to be  $T_2 = 677$  K. Considering that our net work must be 300 kWe, for propulsion purposes, the work to run the compressor was found to be 213 kW. Rearranging Equation 3.12 the compressor inlet temperature was calculated

to be 428 K. Table 3.1 list the temperatures and pressures around the system and Table 3.2 list the energy balance and cycle efficiency (  $\eta = 28.9\%$  ).

The assumptions made produced an overall cycle efficiency somewhat higher than expected without a recuperator in the system. However, carbon dioxide is an excellent working fluid and thus the system seems to be a feasible and beneficial design for the specific application it is intended for.

Table 3.1 System Thermodynamic States

State	Temperature (K)	Pressure (MPa)
1	850	6.58
2	677	3.76
3	428	3.63
4	500	6.90

Table 3.2 System Energy Balance

$$\begin{aligned}
 Q_{\text{in}} &= 1038 \text{ kW} \\
 Q_{\text{out}} &= 738 \text{ kW} \\
 W_{\text{T}} &= 513 \text{ kW} \\
 W_{\text{C}} &= 213 \text{ kW} \\
 \eta &= 28.9 \%
 \end{aligned}$$

## REFERENCES

1. English, R. E. : Power Generation from Nuclear Reactors in Aerospace Applications. NASA TM 83342, Nov. 1982.
2. Rust, J. H. : Nuclear Power Plant Engineering. Georgia Institute of Technology, 1979.
3. Blumenberg, J. ; Ruppe, H. O. : Comparison of Nuclear and Solar Power Plants with Turboelectric Generators for Application in Space. Technical University of Munich D-8000, Apr. 1983.
4. English, R. E. : Power Generation from Nuclear Reactors in Aerospace Applications. NASA TM 83342, Nov. 1982.



#### 4. TURBOMACHINERY

Optimizing the cycle efficiency was a primary concern in choosing the turbine and compressor for this system. However, a consideration of equal importance was choosing turbomachinery that utilized the system's unique design. The use of carbon dioxide as the primary working fluid for the system constrained the availability of information and the applicability of previously tested systems.

Cycle efficiency optimization for turbine and compressor considerations means achieving the maximum efficiency for both components with minimum mass. In order to accomplish this the turbine needs to have as high an inlet temperature as possible and the compressor pressure ratio ( $r_p$ ) optimized. According to calculations by Blumenberg and Ruppe (1) component efficiencies of  $\eta_T=0.86$  and  $\eta_C=0.86$  can be expected for the turbine and compressor, respectively, for the system characteristics under consideration. For tests typical of the system under consideration compressor pressure ratios of  $r_p=1.90$  were found to be optimum by English (2). The combined mass of these two components and the rotating shaft connecting them is estimated to be 215 kilograms by English.

The system's unique design (Figure 4.1), most significantly having the heat source located between the turbine and compressor on the shaft is perfect for the use of axial flow turbomachinery. Outlet flow from the reactor can more easily be directed to an axial flow turbine (Figure 4.2) for this design. Likewise, both inlet and outlet flow of the compressor is more easily directed for this design with axial flow (Figure 4.3). Axial flow turbomachinery has the advantage of higher efficiency over radial flow, but have the disadvantage of a smaller pressure rise per unit mass in the compressor. These effects balance themselves out and the design concern becomes the factor that makes the axial flow turbomachinery the best alternative for this specific design.

Not much research has been conducted in the area of carbon dioxide turbomachinery and therefore the parameters taken in references 1 and 2 may have some error. Since no available data on turbomachinery could be found for carbon dioxide the efficiencies for both components were taken from data for an argon working fluid at similar conditions. Argon and carbon dioxide have

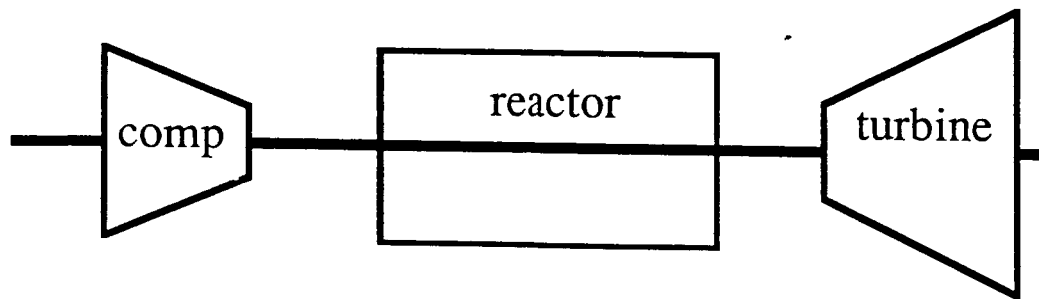


Figure 4.1 Reactor and Turbomachinery Schematic



CS-40265

Figure 4.2 Axial Flow Turbine

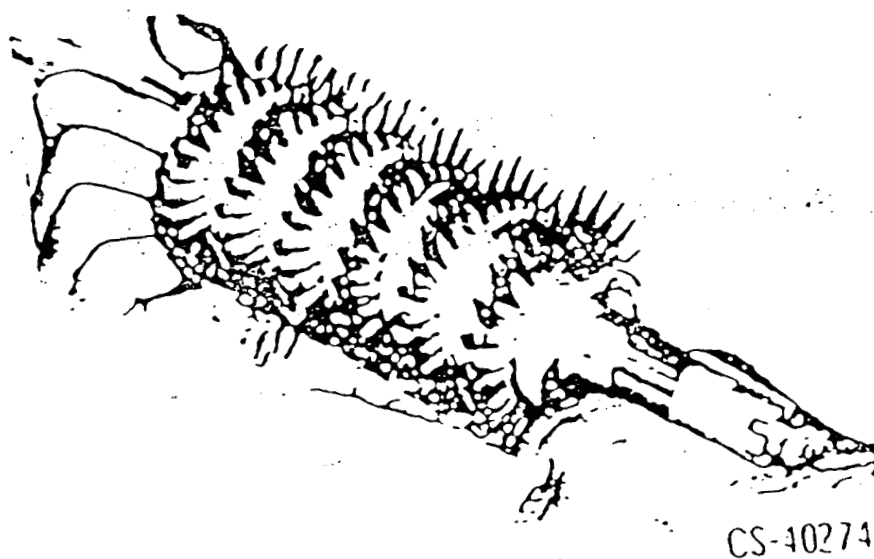


Figure 4.3 Axial Flow Compressor

similar atomic masses, 40 and 44 amu respectively, and therefore the approximation is good as long as the carbon dioxide does not dissociate. At the temperatures under consideration dissociation has a very low probability and thus the approximation should be good as far as turbomachinery is concerned.

## References

1. Blumenberg, J. ; and Ruppe, H. O. : Comparison of Nuclear and Solar Power Plants with Turboelectric Generators for Application in Space. Technical University of Munich D-8000, Apr. 1983.
2. English, R. E. : Power Generation from Nuclear Reactors in Aerospace Applications. NASA TM 83342, Nov. 1982.

## 5. HEAT REJECTION SYSTEM

### 5.1 Introduction

Heat rejection techniques for space reactor systems are distinctly different from those used in terrestrial reactor systems. Terrestrial reactors reject heat by convection to the atmosphere through such mechanisms as cooling towers and to bodies of water such as cooling ponds or rivers. Space power systems can employ neither of these techniques for dissipating waste heat. The only viable mode of heat rejection for space power systems is that of radiation to the surrounding space. Space power systems cannot employ convective heat rejection systems because they operate in a virtual vacuum.

There are two operable types of radiator systems available for space power systems. The first system employs a continuous fluid loop which serpentinely winds through the radiator, such as the pumped fluid radiator used on the Shuttle (Pearson and Dabrowski, p. 806). The second type employs a heat pipe radiator system which transfers heat from the reactor coolant to heat pipes. A heat pipe radiator system was chosen for the space power system under discussion.

The decision to use heat pipes was based on safe operating criteria for a long duration of unmanned operation. The continuous fluid loop design has the disadvantage that if the radiator sustains meteoroid damage, there is a great probability that the entire cooling loop will be lost. To compensate for this, redundancy and isolation valves must be built into the radiator system. This addition increases system mass and reduces reliability, both critical requirements for space operation of long duration. The heat pipe radiator system is a more reliable system because each heat pipe operates independently of the other heat pipes in the radiator. If one heat pipe sustains meteoroid damage, the other heat pipes and cycle fluid are unaffected. The reactor coolant flows through a main manifold which can be shielded for meteoroid protection. By including sufficient heat pipe elements to account for losses from meteoroid damage and shielding the manifold, the heat pipe radiator has a high level of reliability. In the system under discussion, it was also crucial to have high levels of reliability in the radiator because one of the design criteria was that the power

system also be operable on the planet of Mars. The winds on Mars are equivalent to 20 mile per hour winds on earth and dust storms are common. This criteria again pointed toward the use of a heat pipe radiator to insure safe operation of the planet of Mars due to the independent operation of each heat pipe unit.

## 5.2 Principles of Operation

The structure of a heat pipe is shown in Figure 5.1. The main regions are the evaporator section and the condenser section. The pipe itself consists of the pipe wall, the wick, and the working fluid. The heat pipe operates on the principle of capillary forces. In order for the heat pipe to operate properly, the maximum capillary pumping head must be greater than the total pressure drop in the pipe, which is made up of the pressure drop required to return the liquid from the condenser to the evaporator and the pressure drop necessary to cause the vapor to flow from the evaporator to the condenser and the gravitational head.

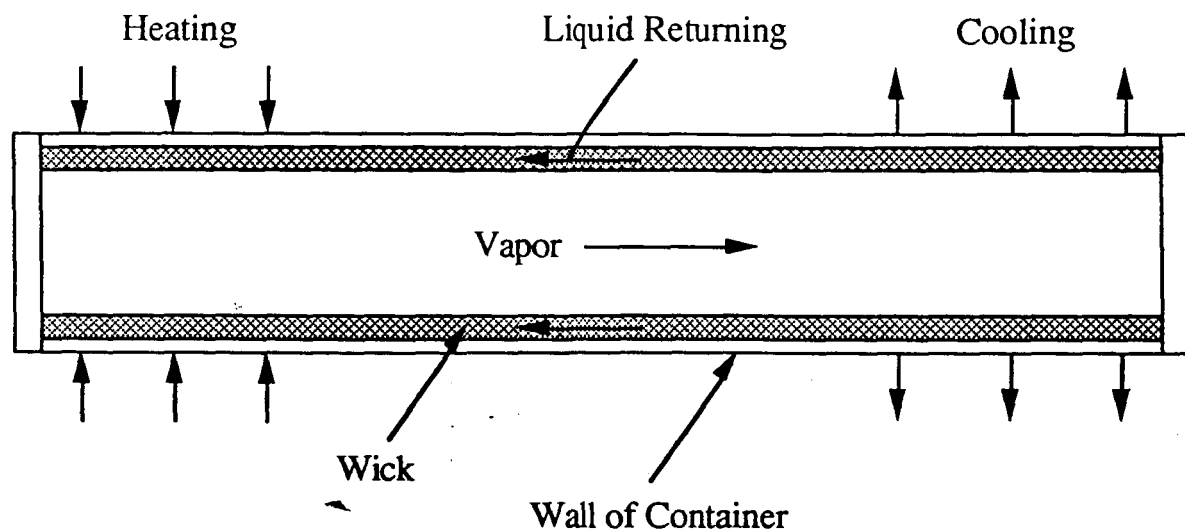


Figure 5.1 - Diagram of Heat Pipe

Energy is transferred to the working fluid of the pipe in the evaporator region. The working fluid begins to evaporate and the vapor moves toward the condenser section of the pipe. When cooled the fluid condenses in the condenser region. Capillary forces in the wick return the



condensate to the evaporator region.

Several characteristics of heat pipes are:

1. can operate in any orientation, evaporator position is not restricted (Dunn and Reay, p. 1)
2. very high effective thermal conductance
3. a near isothermal surface of low thermal impedance, the condenser surface of heat pipe will tend to operate at uniform temperature (Dunn and Reay, p. 3).

In space the latent heat produced when the fluid condenses is dissipated by radiation to the surrounding space. In order to increase the area of radiative heat transfer, fins can be attached to the heat pipe. Since radiative heat transfer is proportional to  $T^4$ , the quantity of energy radiated is larger at higher temperatures. Thus, from the heat transfer standpoint, it is best to reject energy at the highest temperature possible.

### 5.3 Selection of Materials

#### 5.3.1. Working Fluid

The working fluid of the heat pipe was chosen to be sodium. Of primary importance was the useful temperature range of the working fluid. The heat pipe radiator design criteria specified that it should operate in the range of 400-800 K. The two working fluids considered whose useful range fell between these two values were potassium and sodium. Two other working fluids, mercury and cesium, had operating ranges which were compatible with the design criteria. Mercury was not considered due to its toxicity and cesium was not considered due to its high cost. The principle criteria for deciding upon a working fluid were latent heat, thermal conductivity, demonstrated radial heat flux capability, and liquid surface tension. The corresponding properties for sodium and potassium are shown in Table 5.1.

Table 5.1 - Properties for Sodium and Potassium

<u>Fluid</u>	<u>Latent Heat</u>	<u>Thermal Conductivity</u>	<u>Radial Heat Flux Capability</u>	<u>Liquid Surface Tension</u>
Na	4400 kJ/kg	70 W/m °C	200-1250 W/cm <sup>2</sup>	1.5 N/m x 10
K	2050 kJ/kg	49 W/m °C	150-250 W/cm <sup>2</sup>	9.0 N/m x 10 <sup>2</sup>

Another means of comparing working fluids is provided by the Merit number, which is given by:

$$M = \rho_1 \sigma_1 L / \mu_1$$

where  $\rho_1$  is the density of the liquid working fluid  
 $\sigma_1$  is the surface tension  
 $L$  is the enthalpy of vaporization or latent heat  
 $\mu_1$  is the viscosity of the liquid working fluid.

Sodium was chosen over potassium because it has superior qualities over the temperature range of interest as shown in Table 5.1 and because it has a higher merit number than potassium. When the Merit numbers at the boiling points of both fluids are compared, the Merit number for sodium is an order of magnitude larger than that of potassium (Dunn and Reay, p. 91).

### 5.3.2. Container

The container material was selected to be stainless steel (SS 304). The criteria for the container were compatibility with the working fluid, strength, and thermal conductivity. The thermal conductivities of several materials are shown in Table 5.2.

Table 5.2 - Properties of Container Materials

<u>Material</u>	<u>Thermal Conductivity (W/m °C)</u>
Aluminum	205
Copper	394
Stainless Steel 304	17.3

Aluminum and copper both had significantly higher thermal conductivity than stainless steel, but stainless steel was the only material which was compatible with sodium. Stainless steel is also a stronger material than aluminum or copper, which contributes to the reliability of the system.

### 5.3.3. Wick

The wick material chosen was a stainless steel mesh, specifically 508 x 3600 mesh s/s twill. The primary criteria for wick material was demonstrated radial heat flux capability with the working

fluid. Various types of wicks have been tested with sodium as the working fluid. The results are shown in Table 5.3 (Dunn and Reay, p. 95).

Table 5.3 - Wick Performance with Sodium

<u>Wick Material</u>	<u>Radial Heat Flux (W/cm<sup>2</sup>)</u>
s/s mesh	230
various	200-400
3 x 65 mesh s/s	214
508 x 3600 mesh s/s twill	1250

The 508 x 3600 mesh s/s twill had the highest demonstrated radial heat flux capacity. This particular mess twill has a small pore radius, resulting in a large maximum capillary head being generated which aids axial flow.

#### 5.3.4. Fin

Aluminum was chosen to be the fin material. The criteria used for selecting a fin material was strength, weight, and thermal conductivity. Materials considered were aluminum and copper. Stainless steel was also considered since the heat pipes themselves were to be fabricated from it. The thermal conductivities of these three materials are shown in Table 5.2. The densities of the materials are shown in Table 5.4 (Chi, p. 230).

Table 5.4 - Density of Fin Materials

<u>Material</u>	<u>Density</u>
Aluminum	2700 kg/m <sup>3</sup>
Copper	9000 kg/m <sup>3</sup>
SS 304	7850 kg/m <sup>3</sup>

The ultimate tensile strength was compared and stainless steel had the highest, followed by aluminum, which had a slightly higher ultimate tensile strength than copper over the temperature range of interest, 500-800 K (Chi, p. 231). Although stainless steel was strongest, it was eliminated because of its very low thermal conductivity and high density. The ratio of thermal conductivity to density was computed for aluminum and copper. The ratio was  $7.6 \times 10^{-2}$  for

aluminum and  $4.4 \times 10^{-2}$  for copper. Thus, aluminum was chosen since it had the greatest strength and the largest ratio of thermal conductivity to density.

#### 5.3.5. Coating of Heat Pipe Radiator

The energy radiated by a surface is proportional to the emissivity of the surface material. Since the area required is inversely proportional to the emissivity, the mass of the system can be reduced if the emissivity of the surface is increased. A method to increase the emissivity of a surface whose material does not have a large enough emissivity is to coat the surface with a material which has a greater emissivity. Coatings considered were  $\text{Al}_2\text{O}_3$ ,  $\text{ZrO}_2$ , and  $\text{MgO}$ .  $\text{Al}_2\text{O}_3$  was chosen because it had the highest value of emissivity over the temperature range of interest. Its emissivity ranged from 0.85-0.95 (Dieckamp, p. 112). Both the heat pipe and radiator are coated with  $\text{Al}_2\text{O}_3$ .

#### 5.4. Calculations

A basic diagram of the proposed heat pipe system is shown in Figure 5.2. Hot gas enters one end of the radiator system, flows through the manifold into which the evaporator section of heat pipes protrude, and exits at a lower temperature. The purpose of the calculations was to determine the temperature of each heat pipe unit and thus the energy dissipated by that unit, on which basis the number of heat pipes necessary to dissipate the required amount of heat could be found.

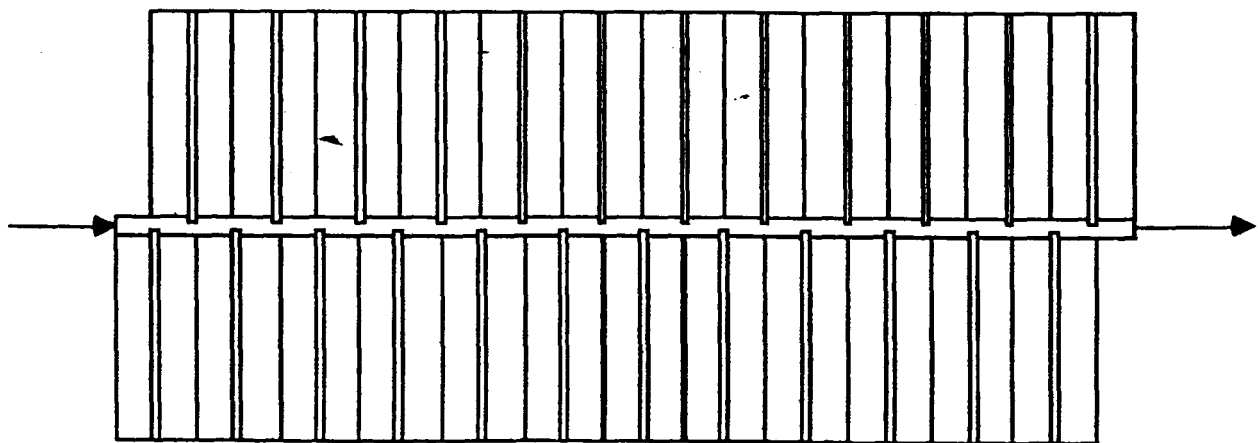


Figure 5.2 - Section of Heat Pipe Radiator System

#### 5.4.1. Basic Energy Equation

The equation used to calculate the energy dissipated by each heat pipe was:

$$Q(z) = \epsilon \sigma A_{hp} (T_{hp}(z)^4 - T^4) + \eta \epsilon \sigma A_f (T_{hp}(z)^4 - T^4)$$

where  $Q(z)$  = energy dissipated by a heat pipe unit

$\epsilon$  = emissivity of the coating

$\sigma$  = Stefan-Boltzmann constant,  $5.670 \times 10^{-8} \text{ W/m}^2 \text{ K}^4$

$\eta$  = fin efficiency

$A_{hp}$  = surface area of the heat pipe

$A_f$  = surface area of the fin

$T_{hp}(z)$  = operating temperature of the heat pipe (position dependent)

$T$  = ambient temperature.

A computer code was written which uses this equation to find the number of heat pipes necessary to dissipate a given amount of heat when the inlet and exit temperatures are specified. The code calculates the operating temperature of each heat pipe unit, the amount of heat dissipated by each unit, and finally, the number of heat pipe units needed and the total radiator mass.

#### 5.4.2. Assumptions

Several simplifying assumptions were made before the equation above was used in the program. These assumptions were:

1.  $\eta = 0.5$ , a conservative estimate since fin efficiency is likely to be above this value because the fins are relatively small in width
2.  $\epsilon = 0.85$ , a conservative estimate since it is the lower end of the range of 0.85-0.95
3. each heat pipe is isothermal (see 5.2. Principles of Operation)
4. the fins operate at the same temperature as the heat pipe; a reasonable assumption considering that the fins are not very wide and aluminum has a high thermal conductivity

5. the temperature of each heat pipe is the temperature of the gas with which it is in contact in the manifold; a reasonable assumption considering that the evaporator region will be completely submersed in the gas and will have numerous small fins protruding into the gas flow to improve convective heat transfer
6. the ambient temperature in space is 0 K.

#### 5.4.3. Preliminary Calculations

Preliminary calculations were necessary to determine the area of the fins and the surface of the heat pipe. Given that the radius of the heat pipes was to be 1.5 cm and the length was to be 4 m, the area was calculated as follows:

$$A_{hp} = 2 \pi r h = 2 \pi (.015 \text{ m}) (4 \text{ m}) = 0.377 \text{ m}^2.$$

The fin area was calculated given that the length was the same as that of the heat pipe, 4 m, and that the width was 0.1 m. The product of length times width was multiplied by two to account for the fin radiating from both sides.

$$A_f = 2 l w = 2 (4 \text{ m}) (0.1 \text{ m}) = 0.8 \text{ m}^2$$

Based on the assumptions and the preliminary calculations, the basic equation can be simplified to the following form:

$$Q = \epsilon \sigma (T_{hp}(z))^4 (A_{hp} + \eta A_f)$$

$$Q = 3.74 \times 10^{-11} \text{ kW/K}^4 (T_{hp}(z))^4.$$

An additional preliminary calculation was undertaken to insure that the heat pipe capability was not going to be exceeded given the operating conditions. The first heat pipe unit in the system dictated the limitations because it would be operating at the highest temperature and rejecting the largest amount of heat. The limiting radial heat flux was estimated to be  $800 \text{ kW/cm}^2$  (see Table 5.1). This value is reasonably conservative because it falls in the middle of the demonstrated radial heat flux range, but is below that actually demonstrated for the chosen mesh. The heat dissipated by the first heat pipe at the maximum operating temperature was determined:

$$Q_1 = 5.34 \times 10^{-11} \text{ kW/K}^4 (800 \text{ K})^4 = 21.87 \text{ kW}.$$

The maximum length of evaporator section which could protrude into the gas flow and still maintain the radial heat flux below its limiting value was calculated as follows, assuming that  $Q_1$  was 25 kW to insure conservatism:

$$25.0 \text{ kW} / 2 \pi r h_e = 0.8 \text{ kW/cm}^2$$

$$h_e = 25.0 \text{ kW} / (0.8 \text{ kW/cm}^2 \times 2 \times \pi \times 1.5 \text{ cm}) = 3.32 \text{ cm}.$$

Thus, the evaporator section can protrude 3.32 cm into the gas flow in the manifold without exceeding the limiting radial heat flux value.

#### 5.4.4. Mass Calculation

The mass of the unit was calculated in the computer code by multiplying the necessary number of heat pipe units by the weight of an individual heat pipe unit. The weight of the unit has three main components: the container, the fluid, and the fin. The mass of the wick can be neglected compared to the masses of these three components. The mass of a unit can then be calculated as follows:

$$\text{mass} = \pi h (r_o^2 - r_i^2) \rho_{ss} + 0.5 \pi h (r_i^2 - r_w^2) \rho_{Na} + h w t \rho_{Al}$$

where  $h$  = length of the unit, 4 m

$r_o$  = outer radius of the heat pipe container, .015 m

$r_i$  = inner radius of the heat pipe container, .0135m

$r_w$  = inner radius of the wick, .0125 m

$w$  = fin width, 0.1 m

$t$  = fin thickness, 3 mm

$$\rho_{ss} = 7850 \text{ kg/m}^3$$

$$\rho_{Na} = 900 \text{ kg/m}^3$$

$$\rho_{Al} = 2700 \text{ kg/m}^3$$

and the 0.5 factor in the second term accounts for the fact that the wick is half full of the working fluid.

The result is that the mass of a single heat pipe unit, consisting of the heat pipe and fin is 7.607 kg.

#### 5.4.5. The Program

The program written to do the heat pipe calculations allows the user to input the inlet temperature, the outlet temperature, and the amount of heat to be dissipated. It uses the basic energy equation and the unit mass formulation to determine the size of the system necessary. A listing of the program is provided in Appendix 5.1.

##### 5.4.5.1. Mass Flow Rate

Using the specified inlet and outlet temperatures and the amount of heat to be rejected, the mass flow rate necessary is calculated from:

$$Q = m c (T_i - T_o)$$

where  $m$  = mass flow rate

$c$  = specific heat of the gas.

Although the specific heat actually varies with temperature, the relationship is fairly linear (Incropera and DeWitt, p. 768). Thus, to calculate the mass flow rate the inlet and outlet temperatures are averaged and the specific heat for the average temperature obtained.

##### 5.4.5.2. Algorithm

The basic algorithm of the program begins with the calculation of the heat loss by the first pipe using the basic energy equation. The temperature of the gas after it transfers heat to the first pipe is then calculated using:

$$T_2 = T_1 - Q / (m \times c)$$

This temperature is used to calculate the heat rejection by the second pipe. Since the heat pipe units are arranged in a series configuration, the temperature of the fluid immediately after it transfers heat to the first pipe is the temperature of the gas when it reaches the next pipe. This process is repeated for additional heat pipes until the specified amount of heat is rejected or the specified outlet temperature is reached. The number of heat pipes necessary is then multiplied by the unit mass to



determine the system mass.

#### 5.4.5.3. Environment

The code has the capability of calculating radiator size for three environments, in space, on the surface of Mars, and buried underground on Mars. Since the only mechanism for energy loss in space is radiation, the basic energy equation can be used unaltered. On Mars energy is lost through convection and radiation. The basic energy equation must be modified with a term to account for convection. The value used for ambient temperature on Mars is 333 K. An appropriate convection coefficient is calculated for each heat pipe unit based on operating temperature. The convection coefficient is calculated using the properties of carbine dioxide, since the atmosphere of Mars is predominantly composed of CO<sub>2</sub>. When the radiator is located underground heat is rejected by conduction. This scenario was treated as conduction from a flat plate in a semi-infinite solid. The heat rejection was calculated from (Rohsenow and Hartnett, p.3-120):

$$Q = k a (T_{hp} - T_w) 2\pi / (\ln (2\pi x/b))$$

where  $k$  = thermal conductivity

$a$  = length of the unit, 4 m

$T_w$  = atmospheric temperature, 333 K

$x$  = depth at which the radiator is buried

$b$  = width of the unit, 0.1 m.

The depth of burial is assumed to be 2 m. Since the thermal conductivity for Martian soil was unavailable the thermal conductivities of sand (0.27 W/mK) and soil (0.52 W/mK) were averaged to obtain an estimate (0.37 W/mK) to be used in the calculations. The resulting expression for energy dissipation was:

$$Q = .00192 \text{ W/K } (T_{hp} - 333)\text{K}.$$

#### 5.4.5.4. Numerical Correlations

Numerical correlations were incorporated as functions in the code. Functions were included to calculate the temperature dependent parameters of specific heat, thermal conductivity, and Nusselt

number. The numerical expressions for specific heat and thermal conductivity were obtained by fitting a curve to tabular data (Incropera and DeWitt, p. 768). Expressions for Nusselt number and the Prandtl and Rayleigh numbers necessary to calculate it were found in a heat transfer textbook (Incropera and DeWitt, p. 767-768).

## **5.5. Results**

### **5.5.1. Preliminary Results**

The specified inlet and outlet temperatures of the radiator were 677 K and 428 K, respectively. The amount of heat to be dissipated was 738 kW. The program was then run under the condition that the environment was space. The resulting mass flow rate was calculated to be 2.43 kg/sec. The number of heat pipes necessary to meet these design criteria was 246. The resulting mass of the radiator was 1870 kg. With the design specifications remaining the same, the program was run under the condition that the environment was on Mars. The number of heat pipes necessary for energy rejection in this case was less than that required in space. The program was then run for the case of the radiator buried underground. The number of heat pipes necessary under this condition was significantly larger than under the other two conditions.

The option of operating the radiator underground on Mars was disregarded because the number of heat pipes necessary for this mode was significantly larger than for either of the other two modes. The limiting case then was operation in space. Since the number of heat pipes required for energy rejection in space was more than that on Mars, the final heat pipe unit numbers and radiator mass are based on this case. The program output is included in Appendix 5.2.

### **5.5.2. Modified Results**

The probability of non-failure was calculated to be 0.864 for the heat pipe design (see VII. Space Logistics). The number of heat pipe units necessary to compensate for failure due to micrometeoroid penetration was obtained by dividing the number of heat pipes units necessary to dissipate the required amount of heat by 0.864. This resulted in a value of 280 heat pipe units. This made the final mass of the radiator 2130 kg. Due to the addition of heat pipe units to compensate for failure, a valve was inserted between the initially calculated number of heat pipe

units, the primary radiator, and the additional heat pipe units, secondary unit. The valve is initially closed but can be opened when failure occurs in the primary radiator unit which warrants additional heat pipe units.

A diagram of the heat pipe radiator unit, including the dimensions of the components, is shown in Figure 5.3. A diagram of the primary radiator system is shown in Figure 5.4. The primary radiator system is 12.3 m long and 8 m wide. The total area of the primary radiator system is 98.4 m<sup>2</sup>. The secondary radiator system is 0.4 m long and 8 m wide and is 3.2 m<sup>2</sup>.

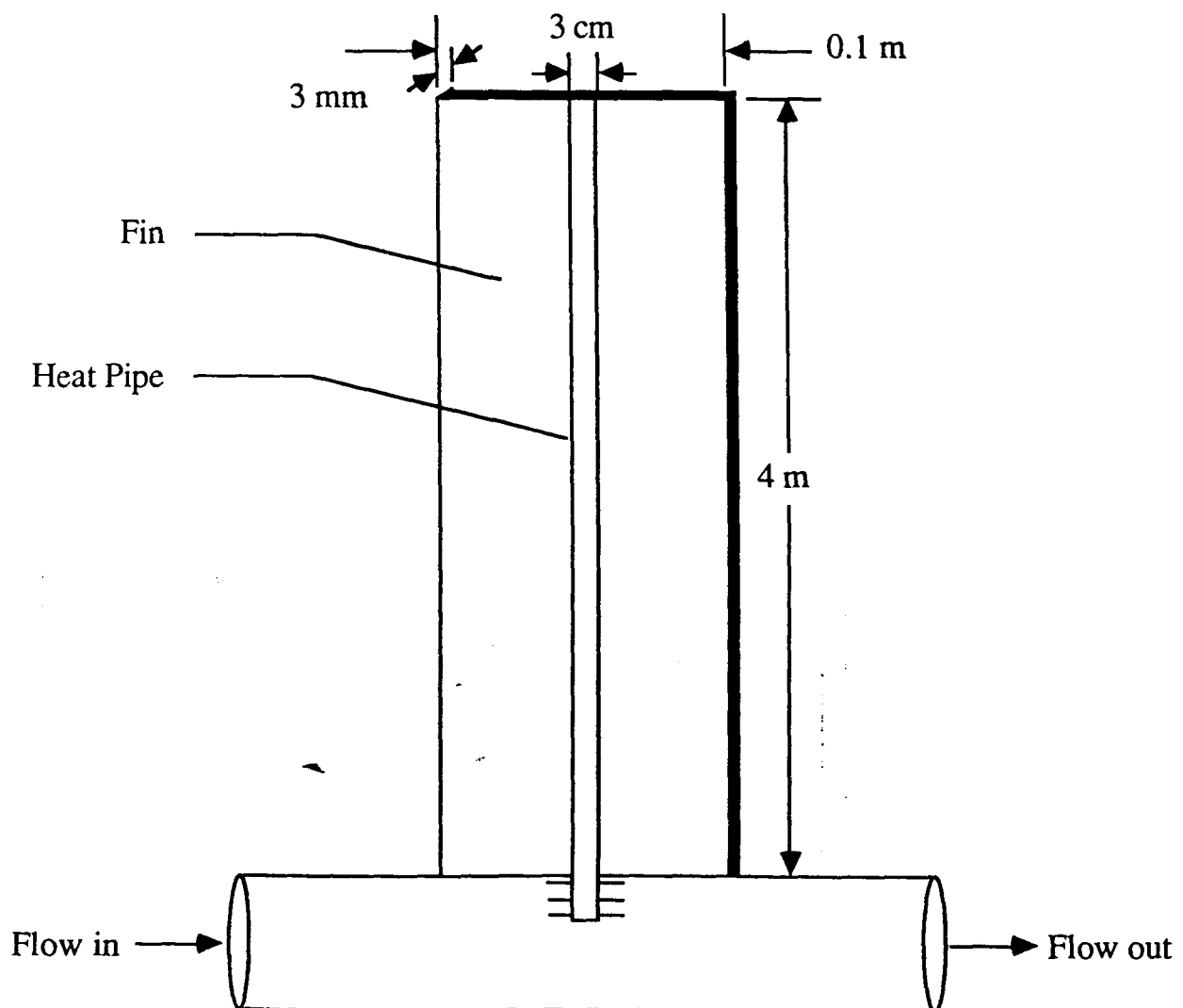
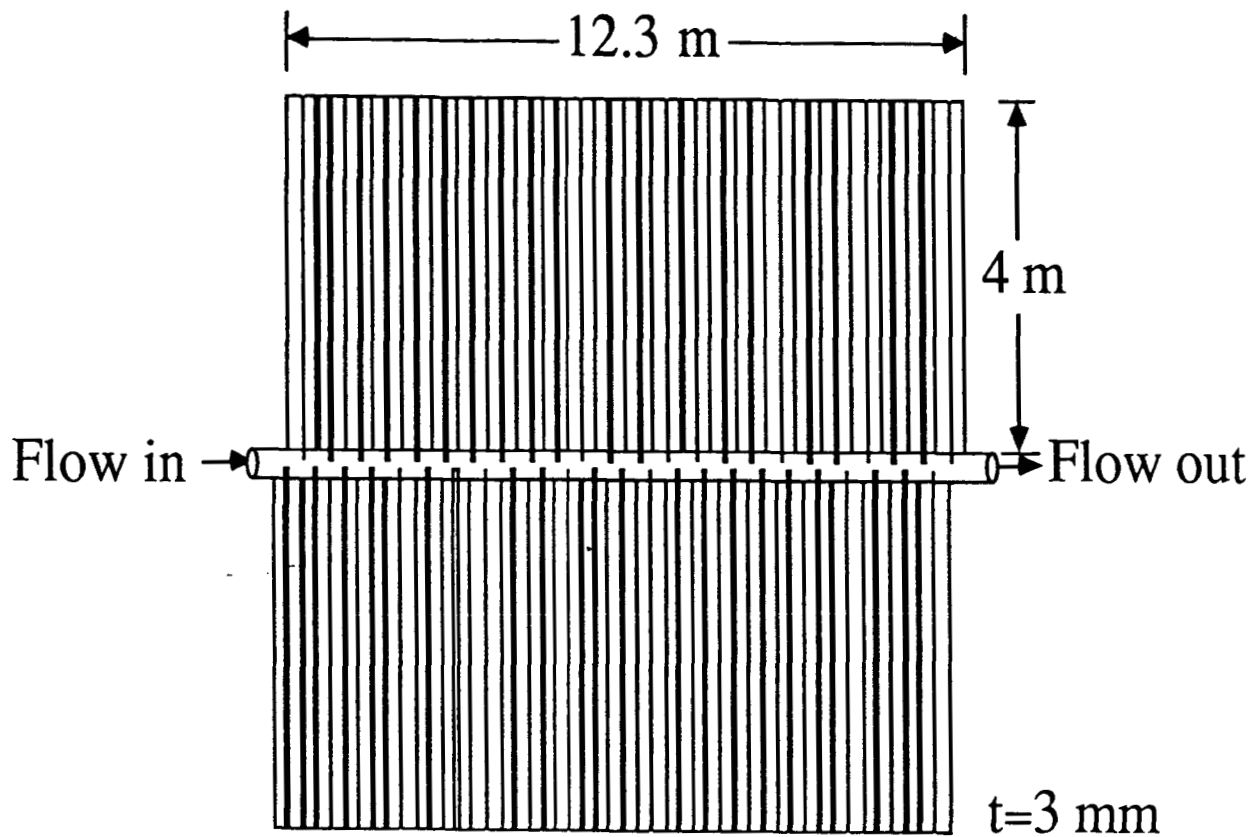


Figure 5.3 - Heat Pipe Unit



- Total Area  $98 \text{ m}^2$
- Mass  $1870 \text{ kg}$
- Heat Pipes  $246$
- Container Material  $\text{SS 304}$
- Coating  $\text{Al}_2\text{O}_3$
- Working Fluid  $\text{Sodium}$
- Fin Material  $\text{Aluminum}$

Figure 5.4 - Heat Rejection System

## 5.6. References

Chi, Heat Pipe Theory and Practice: A Sourcebook, Hemisphere Publishing, Washington, 1976.

Dieckamp, H. M., Nuclear Space Power Systems, Atomics International, Canoga Park, 1967.

Dunn, P. D. and D. A. Reay, Heat Pipes, First Edition, Pergamon Press, New York, 1976.

Incropera, F. P. and D. P. DeWitt, Fundamentals of Heat and Mass Transfer, Second Edition, John Wiley & Sons, New York, 1985.

Pearson, R. and D. Dabrowski, "Optimization of Heat Rejection Subsystem for Solar Dynamic Brayton Cycle Power System," Document Number 860999, Grumman Aerospace Corp., Bethpage, NY.

Rohsenow, W. M. and J. P. Hartnett, Editors, Handbook of Heat Transfer, McGraw-Hill Book Company, New York, 1973.

## Appendix 5.1 - Program Listing

# HEAT PIPE RADIATOR CALCULATIONS

Sandra M. Sloan  
Nuclear Engineering 410  
Design Project

The purpose of this program is to calculate the necessary set of parameters to design a heat pipe radiator. The program calculates the operating temperature of and heat dissipated by each heat pipe which is necessary to meet the design conditions. Based on these calculations, the number of heat pipes necessary to dissipate the required amount of heat is known.

Numerical correlations are used for heat transfer constants such as Nu, c, etc. They are incorporated into the program as functions.

## THE PARAMETERS OF THE SYSTEM

TI = heat pipe radiator inlet temperature  
TO = heat pipe radiator outlet temperature  
QRAD = heat to be dissipated in the radiator  
Q = heat dissipated in the heat pipe  
QTOT = total amount of heat dissipated in the radiator  
MFLOW = mass flow rate  
MASS = mass of unit (heat pipe and fin)  
SYSMAS = mass of radiator system  
NUSSLT = Nusselt number  
C = specific heat  
RK = k, thermal conductivity  
H = h, convective heat transfer coefficient  
ENVIR = environment in which the radiator operates

## DECLARE VARIABLES AND OPEN DATA FILES

REAL MASS, MFLOW  
INTEGER ENVIR

OPEN(UNIT=10, FILE='HEATPIPE1.OUT', STATUS='NEW')  
OPEN(UNIT=12, FILE='HEATPIPE2.OUT', STATUS='NEW')  
OPEN(UNIT=15, FILE='HEATPIPE3.OUT', STATUS='NEW')

PRINT\*, 'INLET TEMPERATURE(K)?'  
READ\*, TI  
PRINT\*, 'OUTLET TEMPERATURE(K)?'  
READ\*, TO  
PRINT\*, 'HEAT TO BE DISSIPATED(KW)?'  
READ\*, QRAD

```

*****
C If the environment is the MARTIAN ATMOSPHERE, then the radiator
C rejects heat by radiation and convection.
C
C
C DO 80 I=1,300
C
C   H=NUSSLT(TOLD)*RK(TOLD)/O.20
C
C   Q=3.74E-11*(TOLD**4-1.23E10)+H*O.80*(TOLD-333)
C
C   PRINT*,I,Q,TOLD
C   WRITE(12,*)I,Q,TOLD
C
C   TNEW=(-1.0*Q/MFLOW/C(TOLD))+TOLD
C   QTOT=QTOT+Q
C
C   ERRORT=TOLD-T0
C   ERRORQ=QTOT-QRAD
C
C   IF(ERROR.LT.O.OR.ERRORQ.GT.O)THEN
C     PRINT*,QTOT
C     WRITE(12,*)QTOT
C     GO TO 100
C   ENDIF
C
C   TOLD=TNEW
C
C   CONTINUE
C
C GO TO 100
C
*****
C If the heat pipe system is UNDERGROUND, it rejects heat by conduc-
C tion to the surrounding soil.
C
C
C DO 140 I=1,350
C
C   Q=.00192*(TOLD-333)
C   PRINT*,I,Q,TOLD
C   WRITE(15,*)I,Q,TOLD
C
C   TNEW=(-1.0*Q/MFLOW/C(TOLD))+TOLD
C   QTOT=QTOT+Q
C
C   ERRORT=TOLD-T0
C   ERRORQ=QTOT-QRAD
C
C   IF(ERROR.LT.O.OR.ERRORQ.GT.O)THEN
C     PRINT*,QTOT
C     WRITE(15,*)QTOT
C     GO TO 100
C   ENDIF
C
C   TOLD=TNEW
C
C   CONTINUE
C
140

```



RETURN  
END

[illegible]

## Appendix 5.2 - Program Output

# ENVIRONMENT IS SPACE

No.	Q(z)	T(z)
1	7.856449	677.0000
2	7.742754	674.5373
3	7.631761	672.1068
4	7.523382	669.7078
5	7.417531	667.3397
6	7.314126	665.0016
7	7.213094	662.6932
8	7.114356	660.4136
9	7.017840	658.1622
10	6.923476	655.9385
11	6.831200	653.7419
12	6.740946	651.5718
13	6.652653	649.4277
14	6.566266	647.3091
15	6.481723	645.2154
16	6.398973	643.1461
17	6.317961	641.1008
18	6.238640	639.0790
19	6.160957	637.0803
20	6.084868	635.1041
21	6.010328	633.1500
22	5.937291	631.2177
23	5.865718	629.3068
24	5.795566	627.4167
25	5.726798	625.5472
26	5.659375	623.6978
27	5.593261	621.8682
28	5.528422	620.0581
29	5.464821	618.2670
30	5.402427	616.4947
31	5.341210	614.7408
32	5.281138	613.0049
33	5.222180	611.2869
34	5.164309	609.5862
35	5.107497	607.9028
36	5.051715	606.2361
37	4.996940	604.5861
38	4.943145	602.9523
39	4.890307	601.3345
40	4.838403	599.7325
41	4.787410	598.1461
42	4.737303	596.5748
43	4.688063	595.0185
44	4.639669	593.4769
45	4.592101	591.9499
46	4.545338	590.4371
47	4.499364	588.9384
48	4.454157	587.4535
49	4.409702	585.9822
50	4.365982	584.5243
51	4.322977	583.0796
52	4.280675	581.6479
53	4.239056	580.2289
54	4.198108	578.8226
55	4.157814	577.4287
56	4.118161	576.0470

118	2.500220	508.4832
119	2.483212	507.6163
120	2.466396	506.7547
121	2.449770	505.8986
122	2.433331	505.0477
123	2.417075	504.2021
124	2.401000	503.3617
125	2.385103	502.5264
126	2.369381	501.6962
127	2.353832	500.8711
128	2.338453	500.0509
129	2.323241	499.2357
130	2.308193	498.4254
131	2.293309	497.6199
132	2.278584	496.8192
133	2.264017	496.0232
134	2.249604	495.2319
135	2.235345	494.4453
136	2.221236	493.6632
137	2.207276	492.8857
138	2.193462	492.1128
139	2.179792	491.3442
140	2.166264	490.5801
141	2.152876	489.8204
142	2.139626	489.0650
143	2.126511	488.3138
144	2.113531	487.5670
145	2.100683	486.8243
146	2.087965	486.0858
147	2.075376	485.3514
148	2.062912	484.6211
149	2.050574	483.8948
150	2.038358	483.1725
151	2.026265	482.4543
152	2.014290	481.7399
153	2.002434	481.0294
154	1.990695	480.3229
155	1.979070	479.6201
156	1.967558	478.9211
157	1.956159	478.2259
158	1.944869	477.5344
159	1.933688	476.8466
160	1.922614	476.1624
161	1.911647	475.4819
162	1.900783	474.8049
163	1.890022	474.1315
164	1.879363	473.4616
165	1.868805	472.7953
166	1.858346	472.1323
167	1.847984	471.4728
168	1.837719	470.8167
169	1.827549	470.1640
170	1.817472	469.5146
171	1.807489	468.8685
172	1.797597	468.2256
173	1.787796	467.5861
174	1.778084	466.9498
175	1.768460	466.3166
176	1.758922	465.6866
177	1.749471	465.0598
178	1.740105	464.4361

240	1.289566	430.9164
241	1.283948	430.4463
242	1.278370	429.9780
243	1.272833	429.5116
244	1.267336	429.0471
245	1.261879	428.5845
246	1.256461	428.1237

QTOT = 739.2075 kW

SYSMAS = 1871.322 kg

# ENVIRONMENT IS MARTIAN ATMOSPHERE

No.	Q(z)	T(z)
1	8.981834	677.0000
2	8.831504	674.1845
3	8.685153	671.4117
4	8.542637	668.6805
5	8.403819	665.9899
6	8.268565	663.3390
7	8.136749	660.7267
8	8.008252	658.1523
9	7.882959	655.6147
10	7.760760	653.1132
11	7.641548	650.6469
12	7.525225	648.2150
13	7.411695	645.8168
14	7.300862	643.4515
15	7.192644	641.1183
16	7.086950	638.8167
17	6.983700	636.5457
18	6.882819	634.3049
19	6.784231	632.0936
20	6.687865	629.9112
21	6.593650	627.7570
22	6.501520	625.6305
23	6.411414	623.5311
24	6.323268	621.4582
25	6.237024	619.4113
26	6.152626	617.3899
27	6.070017	615.3934
28	5.989150	613.4215
29	5.909972	611.4736
30	5.832431	609.5491
31	5.756485	607.6478
32	5.682088	605.7692
33	5.609193	603.9127
34	5.537763	602.0780
35	5.467754	600.2647
36	5.399128	598.4724
37	5.331849	596.7007
38	5.265877	594.9492
39	5.201180	593.2176
40	5.137723	591.5055
41	5.075475	589.8126
42	5.014402	588.1384
43	4.954474	586.4828
44	4.895660	584.8453
45	4.837935	583.2256
46	4.781269	581.6235
47	4.725634	580.0386
48	4.671005	578.4707
49	4.617360	576.9194
50	4.564669	575.3845
51	4.512913	573.8657
52	4.462068	572.3628
53	4.412112	570.8754
54	4.363023	569.4034
55	4.314780	567.9464
56	4.267364	566.5043

118	2.386726	498.6860
119	2.367563	497.8533
120	2.348632	497.0268
121	2.329927	496.2064
122	2.311446	495.3922
123	2.293184	494.5840
124	2.275138	493.7818
125	2.257305	492.9855
126	2.239680	492.1950
127	2.222261	491.4103
128	2.205045	490.6314
129	2.188027	489.8581
130	2.171205	489.0903
131	2.154576	488.3281
132	2.138137	487.5714
133	2.121885	486.8201
134	2.105817	486.0741
135	2.089930	485.3334
136	2.074220	484.5980
137	2.058686	483.8677
138	2.043326	483.1426
139	2.028135	482.4226
140	2.013112	481.7076
141	1.998253	480.9975
142	1.983558	480.2924
143	1.969022	479.5921
144	1.954644	478.8967
145	1.940422	478.2060
146	1.926352	477.5201
147	1.912434	476.8388
148	1.898664	476.1621
149	1.885040	475.4901
150	1.871560	474.8225
151	1.858223	474.1595
152	1.845026	473.5009
153	1.831966	472.8467
154	1.819043	472.1968
155	1.806254	471.5513
156	1.793597	470.9100
157	1.781070	470.2730
158	1.768672	469.6402
159	1.756401	469.0115
160	1.744255	468.3869
161	1.732232	467.7664
162	1.720330	467.1499
163	1.708548	466.5374
164	1.696884	465.9289
165	1.685337	465.3242
166	1.673905	464.7235
167	1.662586	464.1266
168	1.651378	463.5335
169	1.640281	462.9442
170	1.629293	462.3586
171	1.618412	461.7767
172	1.607637	461.1985
173	1.596966	460.6239
174	1.586399	460.0529
175	1.575933	459.4855
176	1.565568	458.9216
177	1.555302	458.3612
178	1.545133	457.8043

240 1.063363 428.9617  
241 1.057457 428.5735  
242 1.051597 428.1873  
QTOT = 738.8571 kW  
SYSMAS = 1840.894 kg



# ENVIRONMENT IS UNDERGROUND

No.	Q(z)	T(z)
1	0.6604800	677.0000
2	0.6600825	676.7930
3	0.6596852	676.5861
4	0.6592880	676.3792
5	0.6588911	676.1725
6	0.6584944	675.9658
7	0.6580978	675.7593
8	0.6577015	675.5529
9	0.6573052	675.3465
10	0.6569093	675.1403
11	0.6565135	674.9341
12	0.6561179	674.7281
13	0.6557225	674.5222
14	0.6553274	674.3163
15	0.6549323	674.1106
16	0.6545375	673.9050
17	0.6541429	673.6994
18	0.6537484	673.4940
19	0.6533542	673.2886
20	0.6529601	673.0834
21	0.6525662	672.8782
22	0.6521726	672.6732
23	0.6517791	672.4683
24	0.6513858	672.2634
25	0.6509927	672.0587
26	0.6505998	671.8541
27	0.6502071	671.6495
28	0.6498145	671.4451
29	0.6494222	671.2407
30	0.6490301	671.0365
31	0.6486381	670.8323
32	0.6482463	670.6283
33	0.6478548	670.4244
34	0.6474634	670.2205
35	0.6470722	670.0168
36	0.6466811	669.8131
37	0.6462904	669.6096
38	0.6458998	669.4061
39	0.6455093	669.2028
40	0.6451191	668.9995
41	0.6447290	668.7964
42	0.6443391	668.5933
43	0.6439495	668.3904
44	0.6435601	668.1876
45	0.6431708	667.9848
46	0.6427817	667.7822
47	0.6423928	667.5796
48	0.6420041	667.3771
49	0.6416156	667.1748
50	0.6412272	666.9725
51	0.6408391	666.7704
52	0.6404511	666.5683
53	0.6400633	666.3663
54	0.6396758	666.1645
55	0.6392884	665.9627
56	0.6389012	665.7610

118	0.6152640	653.4500
119	0.6148887	653.2545
120	0.6145135	653.0591
121	0.6141385	652.8638
122	0.6137637	652.6686
123	0.6133893	652.4736
124	0.6130148	652.2786
125	0.6126406	652.0837
126	0.6122666	651.8889
127	0.6118928	651.6942
128	0.6115190	651.4995
129	0.6111456	651.3050
130	0.6107723	651.1106
131	0.6103992	650.9163
132	0.6100263	650.7220
133	0.6096535	650.5279
134	0.6092810	650.3339
135	0.6089087	650.1400
136	0.6085365	649.9461
137	0.6081645	649.7524
138	0.6077927	649.5587
139	0.6074211	649.3652
140	0.6070498	649.1718
141	0.6066785	648.9784
142	0.6063075	648.7852
143	0.6059366	648.5920
144	0.6055659	648.3989
145	0.6051954	648.2059
146	0.6048251	648.0131
147	0.6044550	647.8203
148	0.6040850	647.6276
149	0.6037153	647.4351
150	0.6033457	647.2426
151	0.6029763	647.0502
152	0.6026071	646.8578
153	0.6022381	646.6656
154	0.6018692	646.4736
155	0.6015006	646.2816
156	0.6011322	646.0897
157	0.6007638	645.8978
158	0.6003957	645.7061
159	0.6000279	645.5145
160	0.5996602	645.3230
161	0.5992926	645.1316
162	0.5989252	644.9402
163	0.5985581	644.7490
164	0.5981911	644.5579
165	0.5978243	644.3668
166	0.5974577	644.1759
167	0.5970913	643.9850
168	0.5967250	643.7943
169	0.5963590	643.6036
170	0.5959931	643.4131
171	0.5956274	643.2226
172	0.5952619	643.0322
173	0.5948966	642.8420
174	0.5945314	642.6518
175	0.5941665	642.4617
176	0.5938017	642.2717
177	0.5934371	642.0818
178	0.5930727	641.8920

240	0.5708356	630.3102
241	0.5704827	630.1264
242	0.5701299	629.9427
243	0.5697773	629.7590
244	0.5694249	629.5755
245	0.5690727	629.3920
246	0.5687206	629.2087
247	0.5683687	629.0254
248	0.5680171	628.8422
249	0.5676655	628.6591
250	0.5673141	628.4761
251	0.5669631	628.2933
252	0.5666121	628.1105
253	0.5662614	627.9278
254	0.5659108	627.7452
255	0.5655603	627.5627
256	0.5652100	627.3802
257	0.5648600	627.1979
258	0.5645101	627.0157
259	0.5641604	626.8336
260	0.5638108	626.6515
261	0.5634615	626.4695
262	0.5631123	626.2877
263	0.5627633	626.1059
264	0.5624145	625.9242
265	0.5620658	625.7426
266	0.5617174	625.5612
267	0.5613691	625.3798
268	0.5610211	625.1985
269	0.5606732	625.0173
270	0.5603254	624.8362
271	0.5599779	624.6552
272	0.5596305	624.4742
273	0.5592833	624.2934
274	0.5589363	624.1127
275	0.5585895	623.9320
276	0.5582428	623.7515
277	0.5578963	623.5710
278	0.5575500	623.3906
279	0.5572038	623.2103
280	0.5568579	623.0302
281	0.5565121	622.8500
282	0.5561665	622.6700
283	0.5558210	622.4901
284	0.5554758	622.3103
285	0.5551307	622.1306
286	0.5547858	621.9509
287	0.5544410	621.7714
288	0.5540965	621.5919
289	0.5537522	621.4126
290	0.5534080	621.2333
291	0.5530640	621.0542
292	0.5527202	620.8751
293	0.5523766	620.6962
294	0.5520331	620.5173
295	0.5516899	620.3385
296	0.5513468	620.1598
297	0.5510039	619.9812
298	0.5506611	619.8027
299	0.5503186	619.6243
300	0.5499762	619.4459

## 6. RADIATION SHIELDING

The proposed design is for an unmanned mission to Mars followed by six years of operation on Mars. Consequently, radiation shielding for the mission was determined to be on the order of that required for the SNAP reactor nuclear power systems with doses on the order of the SP-100 mission to the equipment. The specifications for the SP-100 dose rates for an instrument-rated mission are  $10^{13}$  neutrons/cm<sup>2</sup> and  $5 \times 10^5$  rads for gamma rays over a 7.3 year full power lifetime at a 4.5 meter diameter dose plane that is 22.5 meters from the center of the reactor. This results in  $5.1 \times 10^5$  rads over the system lifetime at the dose plane (1).

Utilizing the weight scaling factors developed by Hedgecock and German (2) for instrument-rated shadow shields similar to that of SNAP reactor nuclear power systems the shielding weight was determined. Reactor power, dose rate, and principle overall dimensions make up seven weight scaling factors ( $w_A, w_B, w_C, \dots, w_G$ ) that are multiplied by the base case shielding weight ( $m_0$ )

$$m = m_0 \times (w_A \times w_B \times w_C \times w_D \times w_E \times w_F \times w_G) \quad (6.1)$$

to find the resulting weight ( $m$ ) as in Equation 1. The instrument-rated shadow shield parameters are determined according to Figure 6.1 and the weight scaling factors are determined using the parameters and Figure 6.2. For the instrument-rated shadow shield the base case weight is 844 pounds while the other base case parameters are listed with the actual parametric values and resulting weight scaling factors in Table 6.1. This results in a shield mass of 1122 pounds or 509 kilograms.

The shielding is comprised of tungsten to attenuate the gamma rays and lithium hydride to absorb and moderate neutrons. The total neutron shield weight was taken equal to 1.39 times the weight of the lithium hydride neutron shield material alone to allow for the weight of structural members.

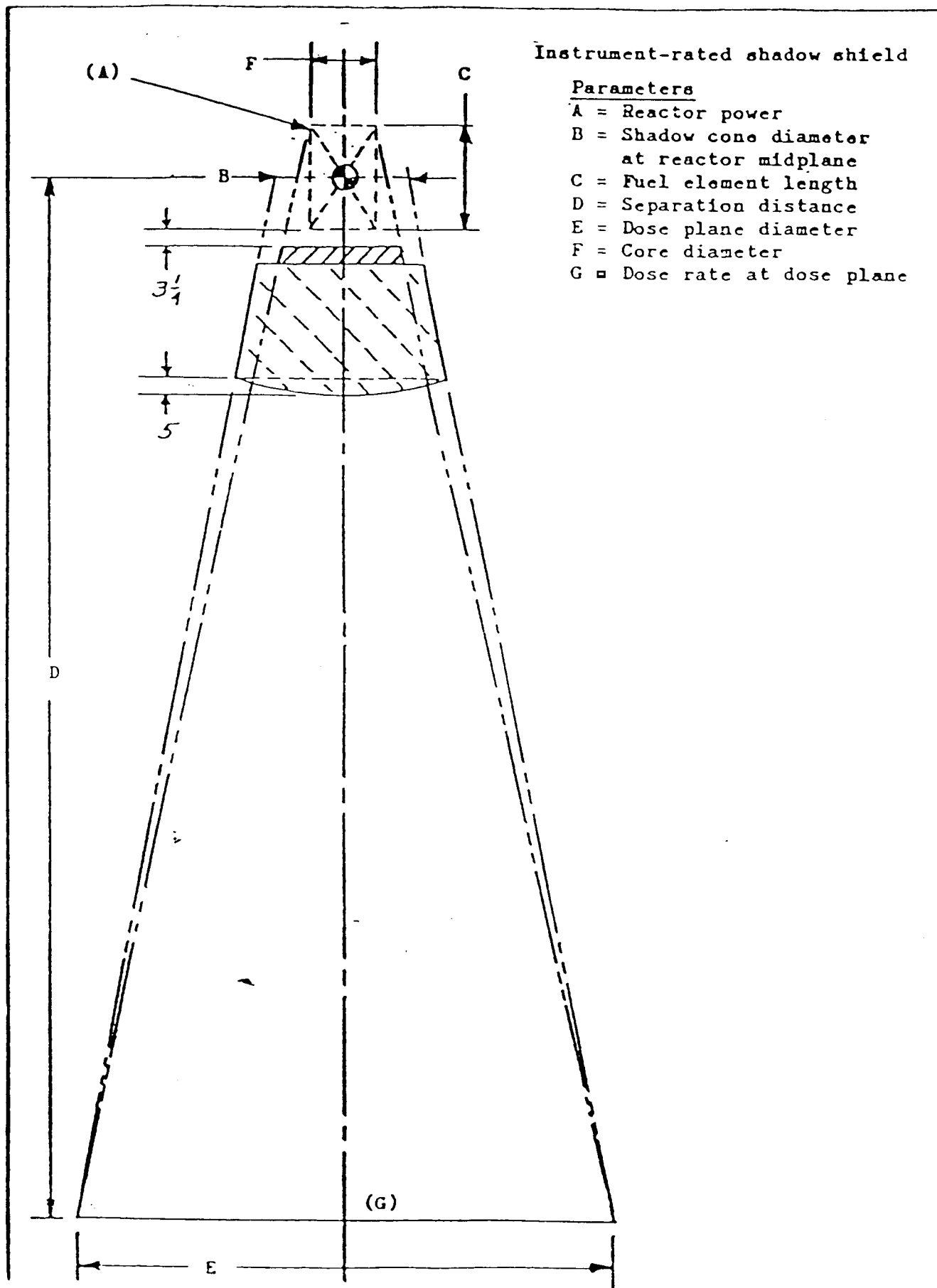


Figure 6.1 Reactor Radiation Shielding Design and Parameters

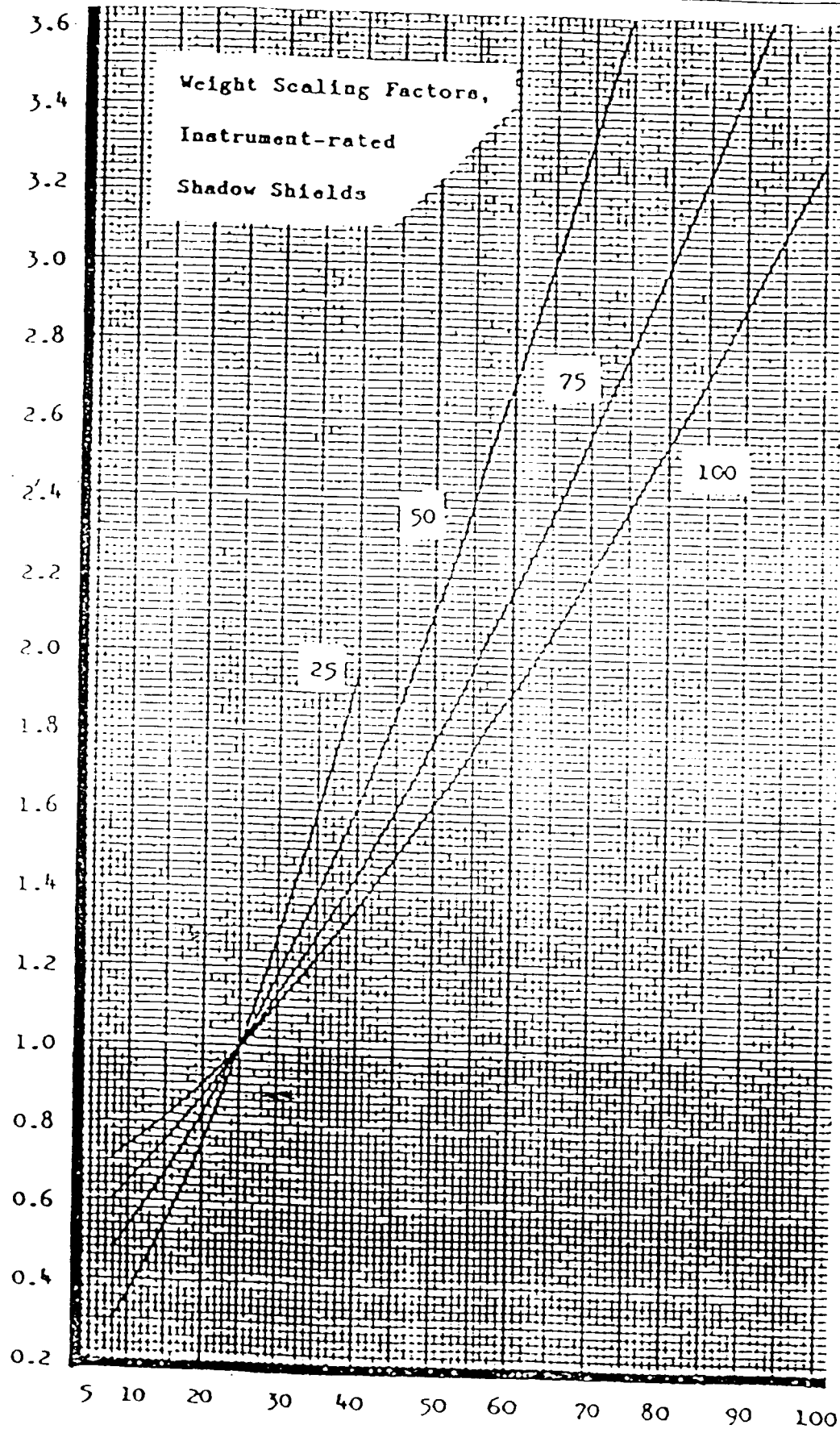


Figure 6.2 Weight Scaling Factors (1 of 3)

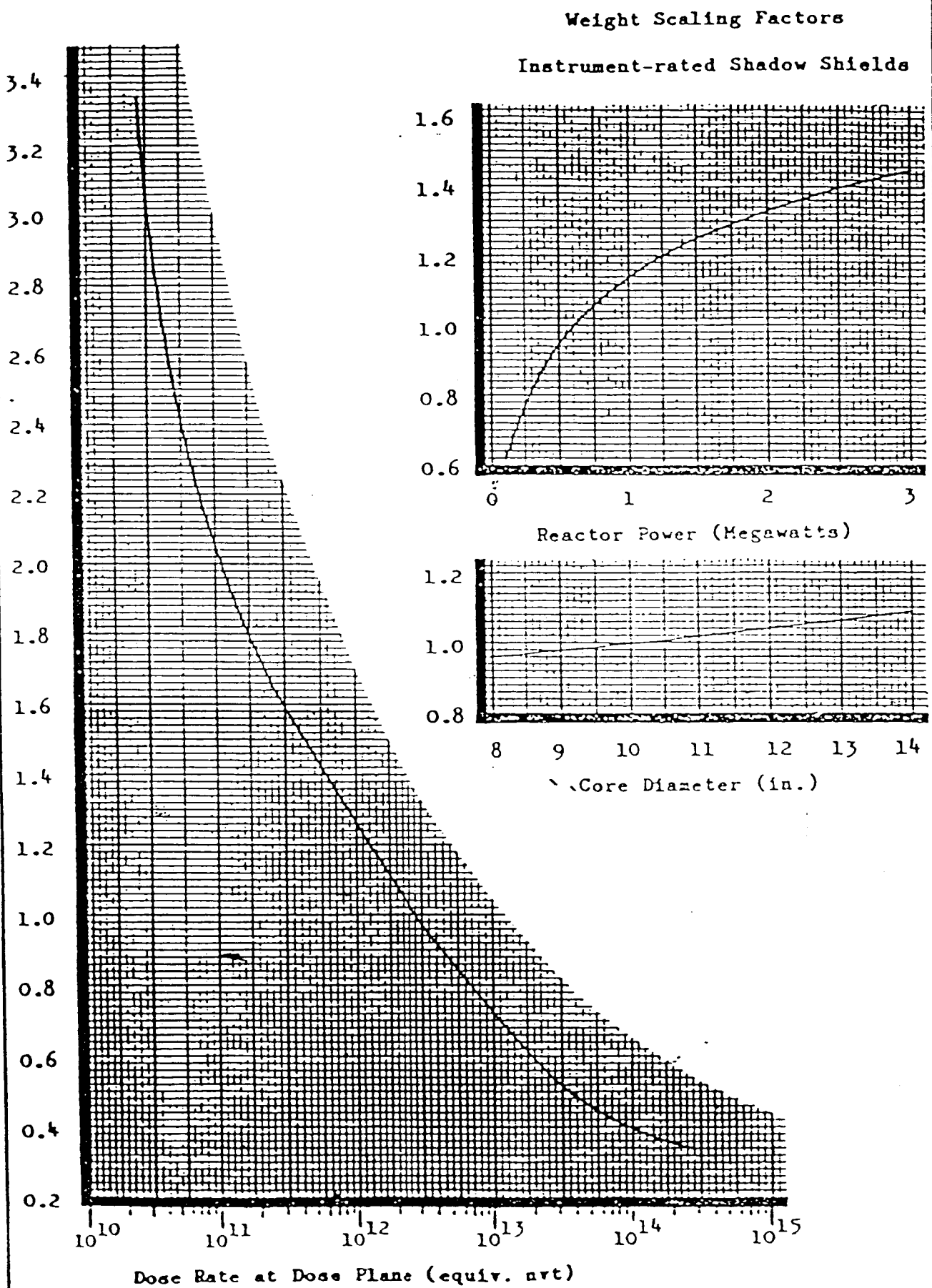
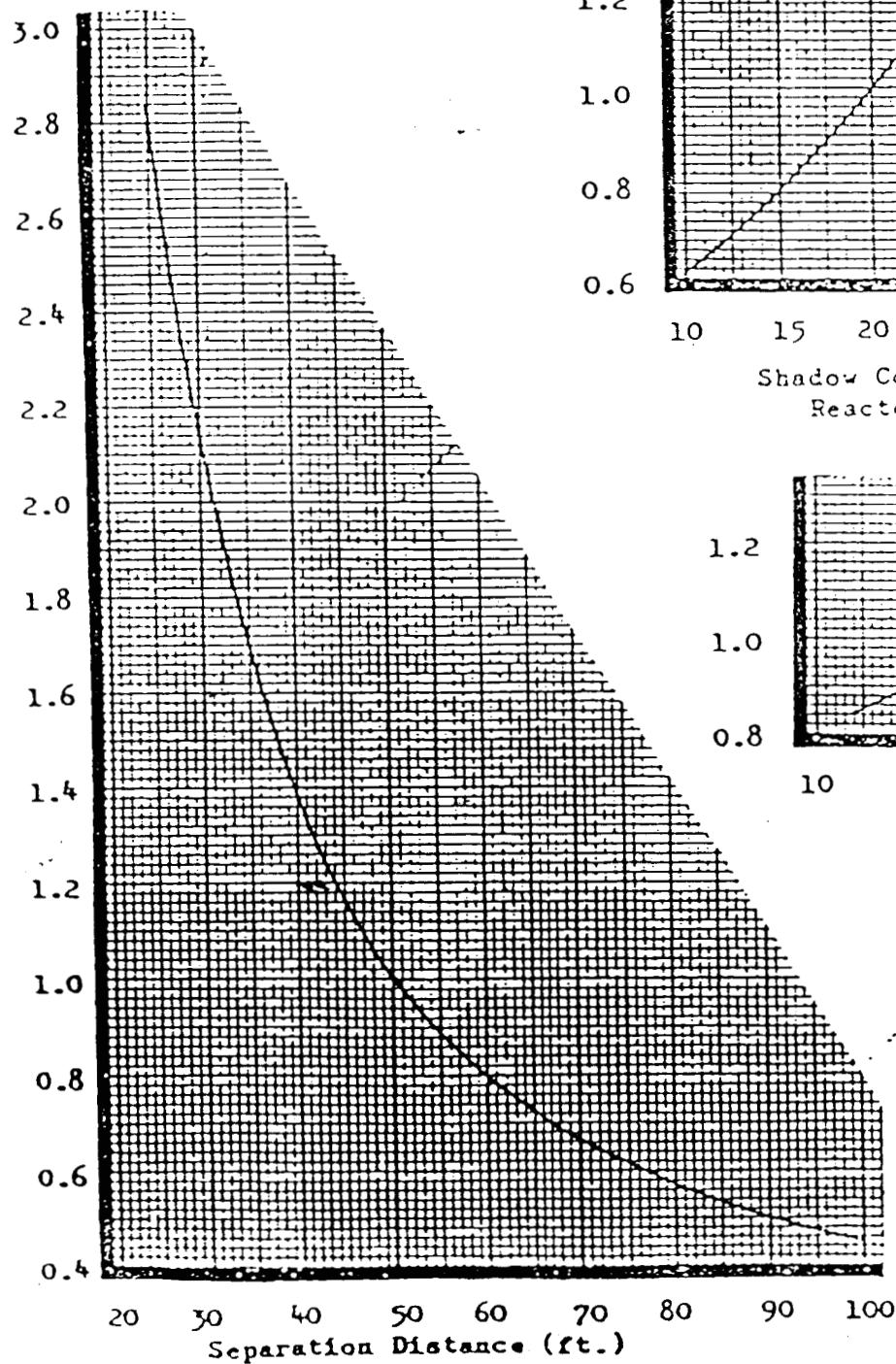


Figure 6.2 Weight Scaling Factors (2 of 3)

# Weight Scaling Factors -

Instrument-rated

Shadow Shields



1.8

1.6

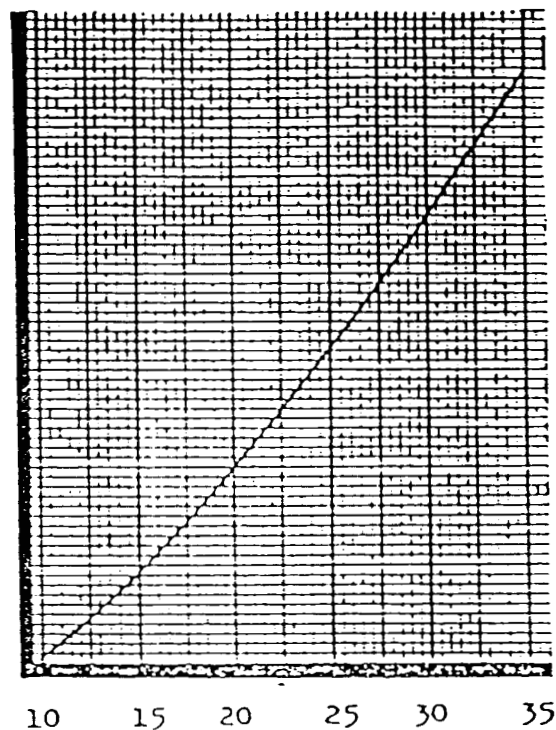
1.4

1.2

1.0

0.8

0.6

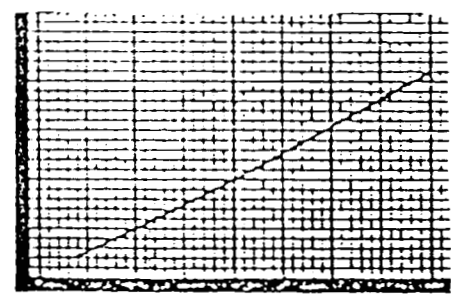


Shadow Cone Diameter at  
Reactor Midplane (in.)

1.2

1.0

0.8



Fuel Length (in.)

Figure 6.3 Weight Scaling Factors (3 of 3)



Table 10.1  
Instrument-Rated Shadow Shield

Parameter	Base Case Value (2)	Actual Value	Weight Scaling Factor
A=Reactor Power	600 kWe	300kWe	0.81
B=Shadow cone diameter at reactor midplane	20 in.	47.2 in.	2.61
C=Fuel element length	20 in.	11.8 in.	0.84
D=Separation Distance	50 ft.	73.8 ft.	0.62
E=Dose plane diameter	25 ft.	14.8 ft.	0.72
F=Core diameter	10 in.	23.6 in.	1.29
G=Dose rate at dose plane	$2.8 \times 10^{12}$ nvt. eq.	$9.15 \times 10^{11}$ nvt. eq.	1.30

## References

1. General Electric Sp-100 Reference Design. GE TM-73219, Jun. 1986.
2. Hedgecock, J. L. ; and German, G. E. : Weight Scaling Factors for SNAP Reactor Shields. Atomic International TDR 11971, Jun. 1966.

## 7. SPACE LOGISTICS

### Mission trajectory

The mission will be initiated with the launch of the MPR-300 aboard one of the space shuttles. Maximum payloads for this launch vehicle are estimated at 35 metric tonnes (MT). Assuming the dimensions of the MPR-300 are compatible with those of the shuttle cargo bay, a single launch would deliver the 8,000-10,000 kg craft into low earth orbit (LEO = 320km). There, the craft would be deployed and cast away from the shuttle. In preparation of obtaining escape conditions from the Earth's gravitation, a chemical booster rocket attachment would be ignited. The transport from one orbit to another for the case of a tangential thrust is described by the following relation:

$$\text{Equation 7.1} \quad t = \frac{WR}{F\sqrt{g_e}} (1/r_0 - 1/r) \quad [1]$$

where  $F/W$  = thrust to weight ratio

$R$  = Earth's radius

$g_0$  =  $9.81 \text{ m/s}^2$

$r_0$  = initial orbit radius

$r$  = orbit radius at time  $t$

With a 100 N chemical thruster, a 10 MT craft has a thrust-to-weight ratio of  $1.02 \times 10^{-3}$ ; to obtain a geosynchronous orbital radius from an initial orbit 320 km in altitude, the craft requires 100 N constant tangential thrust for about 38.8 days. Alternatively, this task could be accomplished using the MPR-300's main propulsion system which was designed for the interplanetary portion of the flight. The main thrusters provide a thrust-to-weight ratio of  $3.67 \times 10^{-5}$ ; although there is some atmospheric drag in LEO, its effect only becomes significant below altitudes of 200 km and thrust-to-weight ratios of  $10^{-6}$  are sufficient to compensate for this effect. Using this in Equation 7.1 yields an orbit transfer time of 1077 days or about 36 months. The reason for choosing a chemical

rocket booster is obvious: mission objectives require a total flight time of approximately 6 months and so 38.8 days is acceptable for LEO to GEO orbit transfer while 1077 days is not.

The next portion of the flight is that from high Earth orbit to Mars orbit. The flight mechanics involved were highly simplified by assuming the craft is far enough from the Earth to neglect its gravitation. This assumption is reasonable considering the previously described orbit transfer. The spacecraft never really approaches a stable geosynchronous orbit (35,870 km) but a highly eccentric one which allows it to eventually escape Earth orbit. In addition, the flight distance was computed from the average radii of revolution about the sun for the Earth and Mars: this implies the necessity of good timing for a rendezvous with Mars on its closest approach to Earth. To further simplify the mission, a "crow flight" path was assumed as well as constant acceleration and craft mass throughout the flight.

Taking all the simplifications into consideration yields a simple equation for rough estimates of flight time and propellant consumption (see Equation 7.2). Although this formulation

Equation 7.2       $t = [2mx/T]^{1/2}$

where  $m$  = total craft mass

$x$  = flight distance

$t$  = flight time

$T$  = total thrust

neglects gravitational effects while near the earth, it also neglects decreasing system mass due to propellant consumption (a large percentage of the total mass) and initial velocity: both of these effects would shorten the flight. With these considerations, Equation 7.2 should yield a conservative estimate of flight time and thus propellant consumption.

Once having reached Mars, the MPR-300 craft will implement aerobraking in the Martian upper atmosphere to slow itself and approach a stable orbit of about 320 km in altitude. This will keep the craft sufficiently high to avoid significant atmospheric drag losses over long periods. The

craft will then initiate entry into the Martian atmosphere with retroactive bursts of the main thrusters. An ablative heat shield on the ship's fore end will protect the craft in the entry stage. Shortly thereafter, parachutes will be deployed to slow the descent to less than 20 fps [2]. As a reference, 30 fps is a typical parachute letdown of a man; the more sensitive reactor system would probably tolerate only less than this. It is therefore necessary to have a damping mechanism to absorb the shock of touchdown.

## Propulsion

The main propulsion system consists of an array of MPD (magnetoplasmadynamic) thrusters, 4 operating simultaneously. These thrusters were selected from a variety of electrically driven ion thrusters on the basis of high efficiency, high specific impulse (see Table 7.1), and compact size (each is less than 12" long and 6" in diameter). Also, the unit's small power requirements allow for modular fitting to the supplied power. Ammonia was chosen over argon and hydrogen as the propellant for the MPD's based on experimental data [3] which showed it to be the most efficient of any studied. For the operation parameters given in Table 7.1, the MPD produces a thrust of 0.9 N. Power supply limitations restrict the number of units, to be operated simultaneously, to 4 for a total thrust of 3.6 N and power consumption of 296 kw<sub>e</sub>. To insure reliability, a total of 13 thrusters were included in the design based on a unit life of 3 months for similar ion thrusters [3].

Table 7.1

MPD Thruster Parameters [3]

current	: 2000 A
voltage	: 37.0 V
thrust	: 0.9 N
propellant consumption rate	: 0.015 g/s

specific impulse	: 6140 s
efficiency	: 36.6%

Using Equation 7.2 and parameters from Table 7.1, a 10 MT craft with 4 MPD's operating would travel to Mars in 242 days and consume 1257 kg of  $\text{NH}_3$ . For reasons of improving reliability, an extra  $\text{NH}_3$  propellant tank has been added to the final design.

The MPD thruster operates in a pseudo steady state mode and requires a pulsed current with a large DC component. The unusually high currents necessary to operate them are readily supplied by homopolar generators. These devices, although used only experimentally at present, are capable of producing large currents in the pulse mode for extended periods [4]. For their inclusion in this design it was assumed that current work on modularization of these generators has been successful and that homopolar generators are a scalable power source. Advantages of using these devices, as opposed to ordinary generators, is the omission of transformers and wave rectification equipment.

### Propellant tank design considerations

Proper analysis of pressurized tanks in a space environment requires detailed consideration of radiative heat transfer. While free space behaves as a blackbody radiating at 3 K, objects in space do not necessarily approach this condition since they are continually being irradiated by a variety of celestial bodies such as the sun and planets. By evaluating the impact of each radiation source on pressurized tanks of cryogenic fluid it is possible to determine the minimal design specifications necessary for their survivability in a stellar environ.

First, consider a spherical tank one astronomical unit from the sun and orbiting about the Earth. To determine the tank wall thickness required to contain the pressurized  $\text{NH}_3$  it is first necessary to calculate the equilibrium temperature. Performing an energy balance on the tank

$$\text{Equation 7.3} \quad \alpha GA_s - \epsilon \sigma T_s^4 A_s = M c_p dT_s / dt$$

where  $\alpha$  = total hemispherical absorptivity

$G$  = total irradiation ( $\text{W/m}^2$ )

$A_s$  = surface area

$\xi$  = total hemispherical emissivity

$\sigma$  = Stephan-Boltzmann constant

$T_s$  = surface temperature

$M$  = mass of tank

$c_p$  = specific heat of tank material

For the given orbit, the total solar irradiation  $G_{\text{solar}}$  is  $1353 \text{ W/m}^2$  and its spectral dependence may be approximated as if it had been emitted by a blackbody radiating at 5800 K. For the Earth, its contribution at low orbits is a total irradiation  $G_{\text{earth}}$  of  $340 \text{ W/m}^2$ ; it may be modeled as a blackbody emitting at 280 K. The spectral dependence of blackbody emitters is well known (see Figure 7.1), having first been determined by Planck, and is tabulated along with other blackbody radiation functions [5].

The tank model consists of a sphere continually and diffusely irradiated over one entire hemisphere while the other side remains in complete darkness. Because the irradiation consists of essentially parallel rays, the actual irradiation varies sinusoidally over the spheres surface (see Figure 7.2a). For conservatism, the tanks will assume the maximum irradiation over the whole hemisphere as shown in Figure 7.2b. Another important consideration of this model is the thermal properties of the tank material and content; if the material used has a sufficiently high thermal conductivity and tank size is small, the tank's dark side will radiate at approximately the same temperature as the irradiated side.

The initial scoping calculation was based on a plain stainless steel tank with a lightly oxidized surface. The spectral absorption for this surface was approximated using the following equation and the data in Figure 7.3.

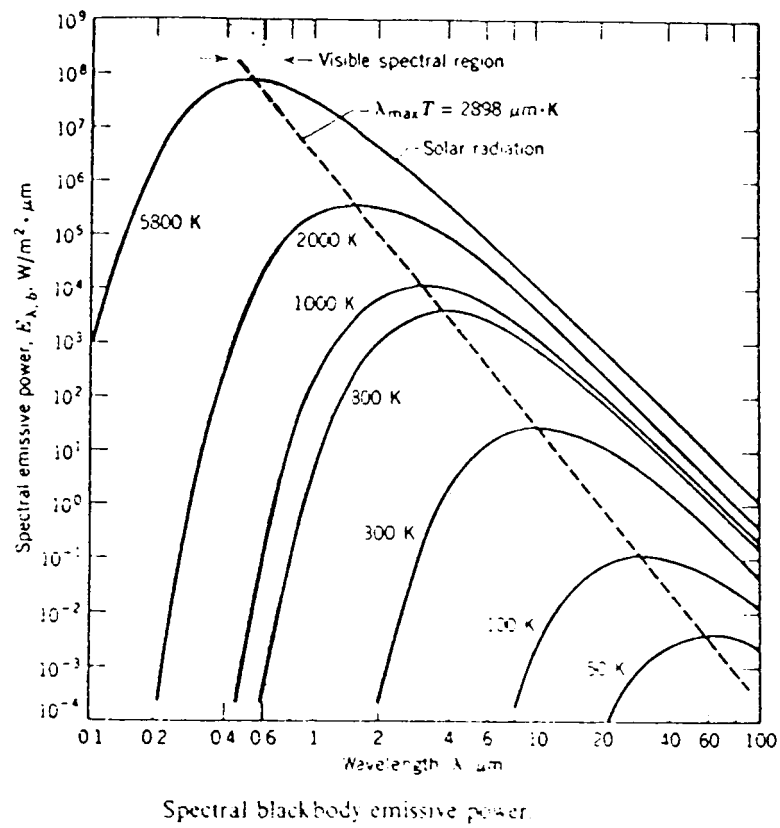


Figure 7.1



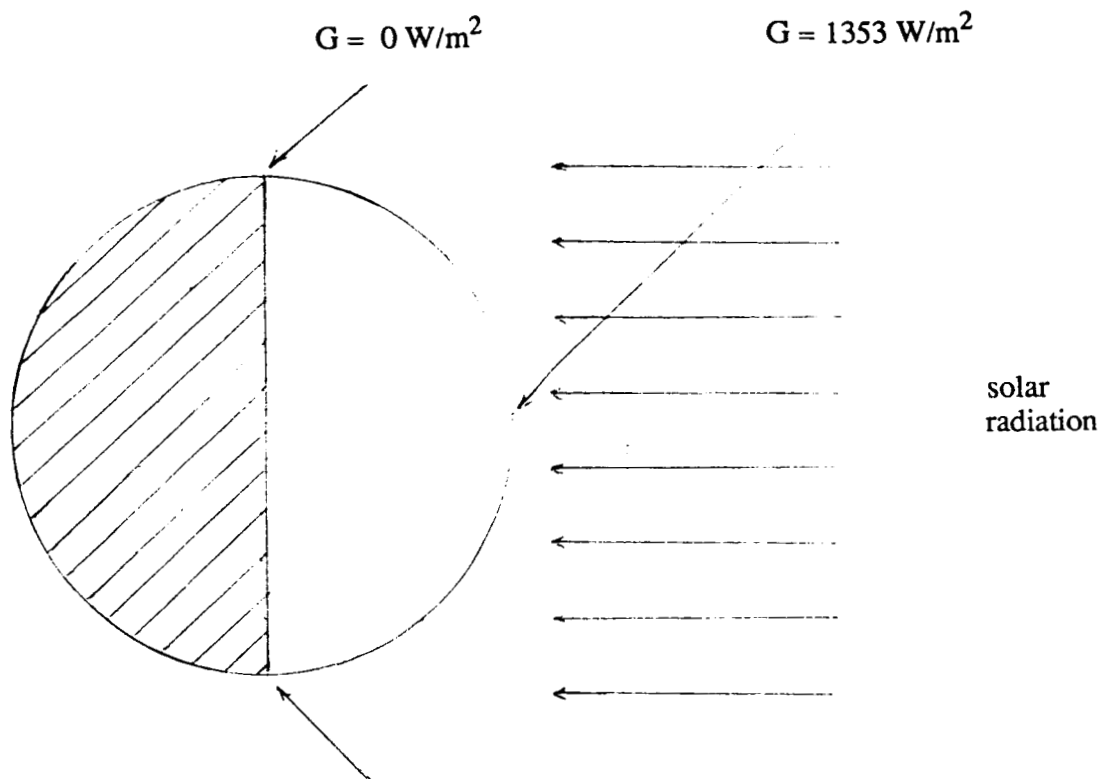


Figure 7.2a Monodirectional irradiation.

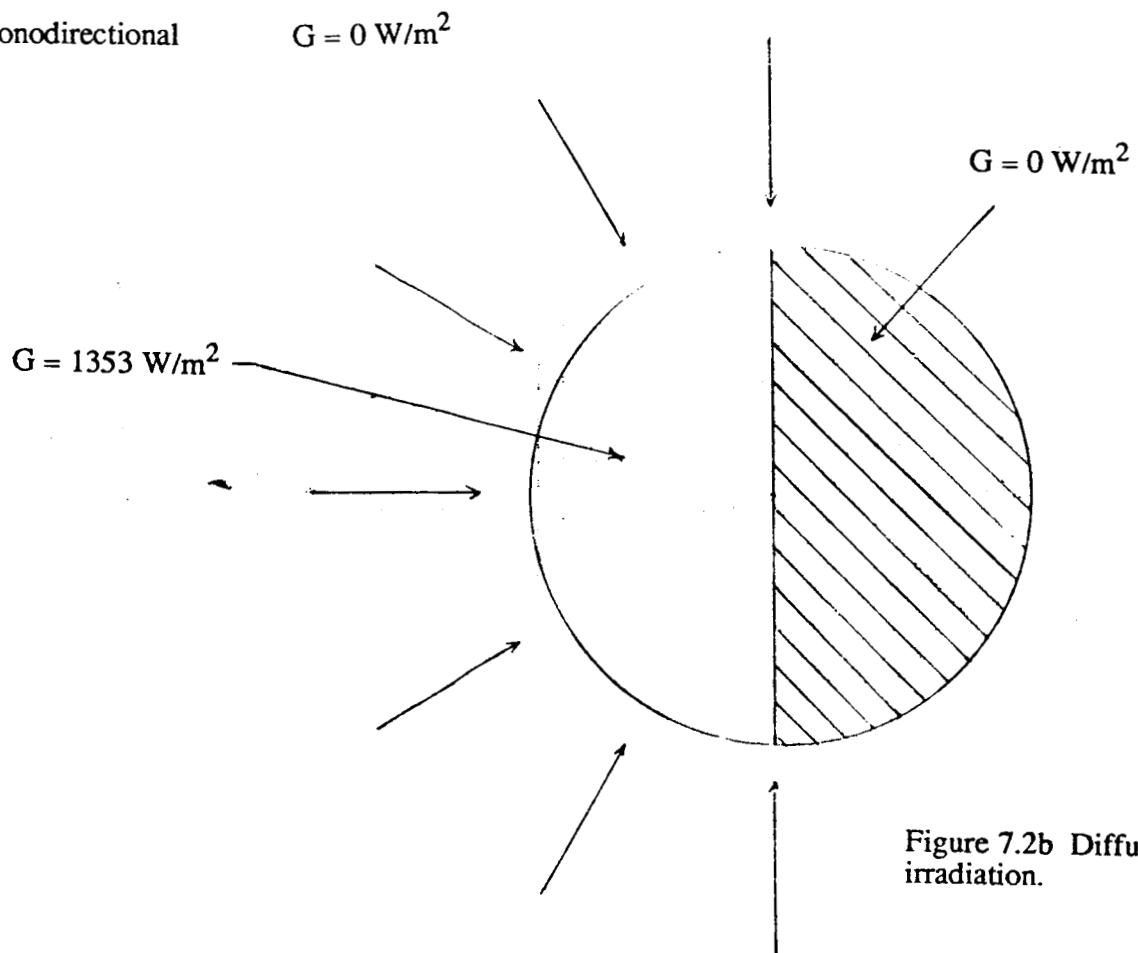


Figure 7.2b Diffuse irradiation.

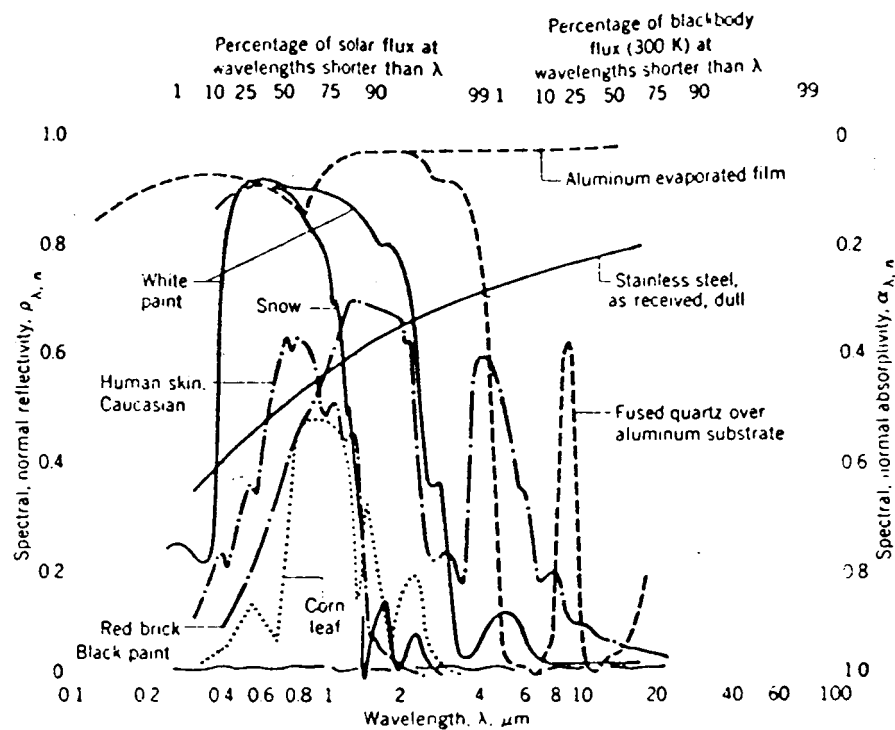


Figure 12.23 Spectral dependence of the spectral, normal absorptivity  $\alpha_{\lambda,n}$  and reflectivity  $\rho_{\lambda,n}$  of selected materials.

Figure 7.3

Equation 7.4

$$\rho_{\lambda} + \alpha_{\lambda} + \tau_{\lambda} = 1$$

where  $\rho_{\lambda}$  = spectral reflectivity

$\alpha_{\lambda}$  = spectral absorptivity

$\tau_{\lambda}$  = spectral transmissivity

Since the tank material will almost invariably be opaque, the transmissivity will be zero. Thus, knowledge of  $\rho_{\lambda}$  implies knowledge of  $\alpha_{\lambda}$ . The absorptivity of stainless steel used in the calculations was made consistently higher than the actual value. Its spectral dependence was approximated as follows:

$$\alpha_{\lambda} = \begin{cases} 0.65 & 0\mu\text{m} < \lambda < 1\mu\text{m} \\ 0.40 & 1\mu\text{m} < \lambda < 4\mu\text{m} \\ 0.30 & 4\mu\text{m} < \lambda < \infty\mu\text{m} \end{cases}$$

The relationship between the total hemispherical absorptivity,  $\alpha$ , in Equation 7.3 and the spectral absorptivity is as follows:

$$\text{Equation 7.5} \quad \alpha \cong \frac{\int \alpha_{\lambda}(\lambda) E_{\lambda,b}(\lambda, T) d\lambda}{\int E_{\lambda,b}(\lambda, T) d\lambda}$$

where  $E_{\lambda,b}(\lambda, T)$  = spectral emissive power of  
a blackbody at temperature T

The value of the integral in the denominator of Equation 7.5 is tabulated for bounds of 0 to  $\lambda$  where  $\lambda$  takes on many incremental values; this is the radiation function  $F_{(0-\lambda)}$ . The radiation function gives the ratio of the total emission from a blackbody at a prescribed temperature for the wavelength interval 0- $\lambda$ . Assuming  $\alpha_{\lambda}(\lambda)$  is constant over certain intervals allows for their removal from the integral in Equation 7.5. Total absorptivity may then be found as follows:

$$\text{Equation 7.6} \quad \alpha = \alpha_{\lambda_1} F(0-\lambda_1) + \alpha_{\lambda_2} (F(0-\lambda_2) - F(0-\lambda_1)) + \dots \\ \alpha_{\lambda_n} (F(0-\infty) - F(0-\lambda_{n-1}))$$

Evaluation of  $\alpha_{\text{solar}}$  in this way yields 0.579 for the solar absorptivity. Multiplying  $\alpha_{\text{solar}}$  and the total solar irradiation  $G_{\text{solar}} (=1353 \text{ W/m}^2)$  yields  $G_{\text{abs}} = 783.4 \text{ W/m}^2$ . Analyzing the irradiation from the Earth in a similar manner reveals  $G_{\text{abs}} = 102.0 \text{ W/m}^2$  for the Earth's radiation. Now, assuming the whole tank is radiating at one temperature, then the tank's emissive power must be half that of the total  $G_{\text{abs}} (=885.4 \text{ W/m}^2)$  since only one half of the tank is being irradiated while both are emitting. Thus it is necessary to find the temperature at which stainless steel has an emissive power of  $885.4 \text{ W/m}^2$ . This may be found from Equation 7.1 by assuming steady state conditions and taking into consideration the temperature dependence of the emissivity  $\epsilon$ ;  $\alpha$  is a weak function of temperature. Figure 7.4 illustrates the effect of increasing tank temperature on emissive power. The necessary operating temperature is slightly less than 500 K.

Since an approximate equilibrium temperature for the propellant tanks is now known, the minimum required tank thickness for assured integrity may now be computed. For stainless steel  $\sigma_{\text{yield}} = 69,000 \text{ psi}$  [6] and from Figure 7.5 the vapor pressure of  $\text{NH}_3$  at 500 K is well above 1000 psi [7].

$$\text{Equation 7.7} \quad t = p_0 R / 2\sigma$$

where  $t$  = tank wall thickness  
 $p_0$  = interior pressure  
 $\sigma$  = design stress of tank material  
 $R$  = tank radius

Using a design stress of 34,500 psi in Equation 7.7, one obtains a tank wall thickness of 7.25 mm for a 1 m diameter tank. This corresponds to a tank mass of 176.5 kg each; at 500 K this would

# Emissive Power vs. Temperature for Lightly Oxidized

Stainless Steel

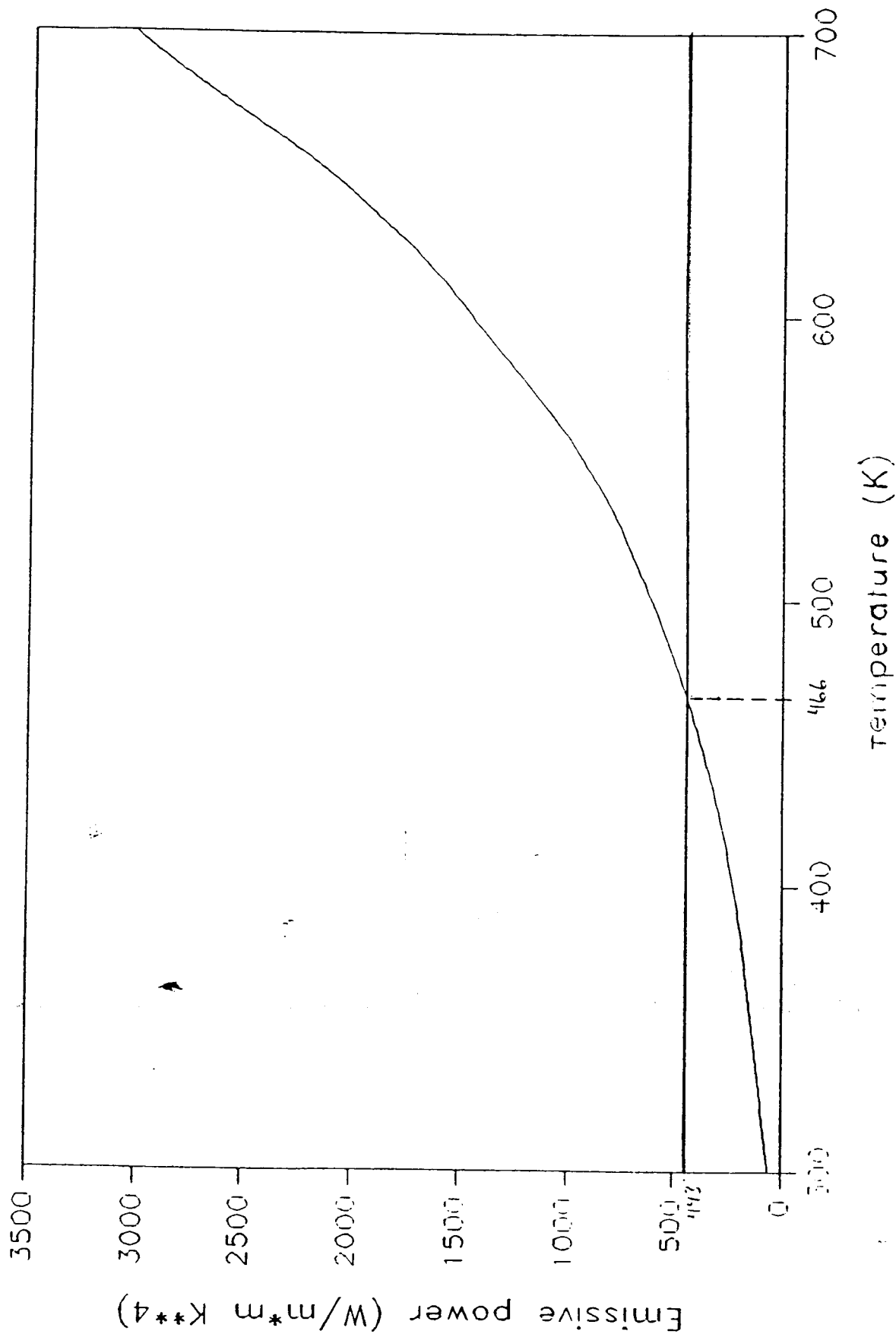


Figure 7.4

require more than 5 tanks (=882 kg) for the MPR-300 transport. This is a rather large mass and the problem deserves further investigation to reduce the tank temperature and pressure, and hence, propellant tank mass.

Upon investigation of the radiative heat transfer characteristics of other materials, a more viable solution was found. The new design requires one additional assumption: that the propellant tanks are always oriented with a highly reflective side toward the sun. The new tank is, as before, constructed from stainless steel; however, one half of each tank is to be coated with an evaporated aluminum film, the other side is to be heavily oxidized. This kind of system is desirable since the aluminum film provides a very high reflectivity while the oxidized stainless steel has a large emissivity. Handling the radiation absorption analysis as before (with the additional assumption), the total irradiation absorbed from the sun and Earth are  $139.7 \text{ W/m}^2$  and  $17.0 \text{ W/m}^2$  respectively for a total of  $156.7 \text{ W/m}^2$  absorbed. Since the emissivity data for aluminum film is available only at 300 K ( $\epsilon = 0.03$  [5]), the energy balance equation (Equation 7.3) will be evaluated at this temperature. From the second term of Equation 7.3, the emissive power of the Al film is  $13.78 \text{ W/m}^2$ . The oxidized side of the tank, however, has an emissive power of  $252.6 \text{ W/m}^2$ . With these figures it is apparent that the tank - if initially above 300 K - would tend to decrease its temperature until an equilibrium condition was achieved at some temperature below 300 K.

Knowing that the tank temperatures can be maintained below 300 K allows for reevaluation of propellant vapor pressures and thus the required tank wall thickness, from Figure 7.5a, the ammonia vapor pressure at 300 K is 166 psi. Using Equation 7.7 for a tank 1 m in diameter, the necessary tank wall thickness is 1.20 mm. This is a marked improvement over 7.25 mm from the prior design; the new tank mass is thus 29.2 kg per unit. The density of  $\text{NH}_3$  at 300 K is  $37.2 \text{ lb}_m/\text{ft}^3$  ( $595 \text{ kg/m}^3$ , see Figure 7.5b). The required propellant mass of 1257 kg will therefore need about 4 tanks for a total propellant tank mass of 116.8 kg: a reduction in total tank mass by a factor of 8 over the first design!

Reiterating some of the conservative approximations: when the tanks are in orbit about the Earth they are being subjected to the most intense irradiation of the entire flight: the idea being

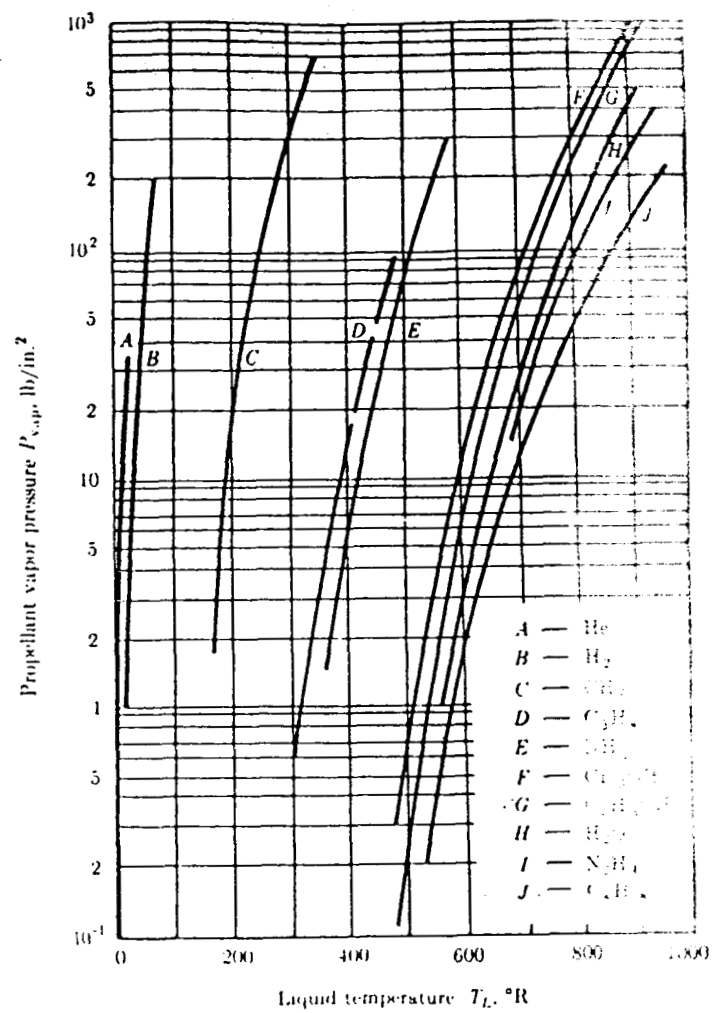


Figure 7.5a [7]

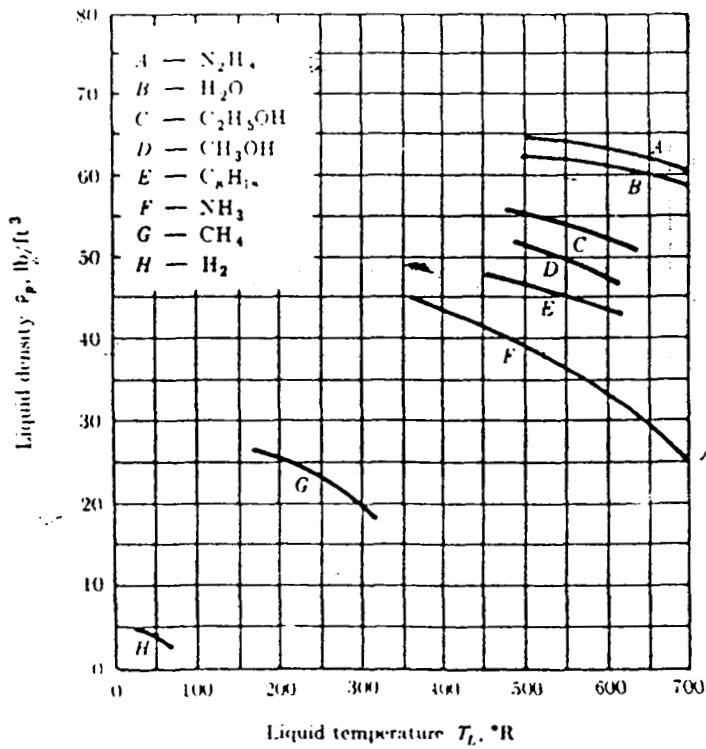


Figure 7.5b [7]

survivability under the most harsh conditions, the actual absorptivities are smaller than those used in the analysis and so the actual total irradiation absorbed  $G_{abs}$  is less, the energy balance showed that the system at 300 K would tend to decrease its temperature and hence actual tank wall thicknesses are larger than necessary to meet pressure requirements at 300 K. Although this design is purely conceptual, the analysis used does present conservative estimates for the pressure tank parameters in a space environ. More importantly, it characterizes the advantages to be had in careful selection of materials.

### **Micrometeoroid damage and failure probabilities**

The vast regions between stellar systems, typically referred to as voids, are in fact not devoid of matter. Although much has yet to be resolved, the existence of particulates and micrometeoroids in space is known and they seriously compromise the survivability of space systems. Careful consideration of micrometeoroid effects is therefore necessary to safeguard against potentially debilitating damage otherwise incurred.

Three principle mechanisms by which meteoroids cause damage are: erosion, spalling, and puncture. The abrasive action of meteoroids on metal surfaces may exceed  $2 \times 10^{-4}$  cm/yr [1]. Meteoroid fluxes may, therefore, lead to significant attrition of a radiator's emissive coating over long periods and result in reducing its radiative power; additional radiator area would then be necessary to accommodate this reduction in radiator effectiveness. In addition, solar protons may erode another  $10^{-3} - 10^{-4}$  cm/yr from exposed surfaces; this effect, however, quickly diminishes as the inverse square of the distance from the sun. Highly reflective surfaces are effected in a similar manner since they are typically either polished metals or an evaporated film. Evaporated films may be on the order of several atomic diameters thick and thus its degradation below the acceptable limits of performance may occur in much less than a year's time. Thickening of the film layer may sufficiently extend performance life.

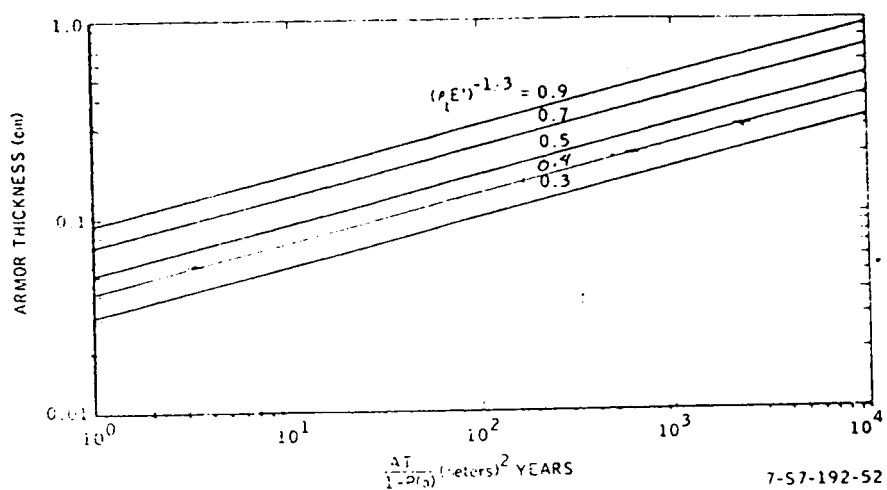
Spalling, the process by which small metal chips are expelled from an interior surface due to an external impact, may cause problems in the power conversion system if the metal chips reach



the turbine or bearings. Due to the large size of the heat pipe radiator in the MPR-300 system, the majority of the spalling is expected to occur there. The radiator modules, however, do not carry primary coolant but their own working fluid. Thus, spalling would have its most serious effect on the heat pipe manifold and other piping in the power conversion cycle, all of which carry primary coolant.

The puncture of system piping by micrometeoroids also poses a serious problem. The resulting holes in piping and tanks allows the slow release of coolant and propellant respectively. This implies the necessity of a reserve coolant supply as well as extra propellant and redundant heat pipes. Since only sparse knowledge of the meteoroid environment in space exists, it is difficult to establish protection criteria. However, the meteoroid penetration theory developed by Summers and Charters [8] provides a means of calculating the probability of no system failure due to meteoroid damage ( $P_0$ ) as a function of exposed area, exposure time, and target material properties ( see Figure 7.6 and Table 7.2). With these data it is possible to view the required propellant tank wall thickness from another design perspective as well as provide an estimate of the additional heat pipe modules necessary to accommodate failures.

Considering a propellant temperature of 300 K, the required tank wall thickness was found to be 1.2 mm stainless steel. From Table 7.2 and Figure 7.6, this yields a  $P_0$  of 0.968 for all 4 propellant tanks in a nine month flight. For reliability, an additional tank of  $\text{NH}_3$  was included in the design. Computation of the appropriate number of redundant heat pipes was carried out similarly. The heat pipe walls are composed of stainless steel 1.5 mm thick and the sensitive area is  $29.5 \text{ m}^2$ ; a nine month flight yields  $P_0 = 0.864$ . Subsequently, about 13.6% more heat pipes were added to the design to replace the number of expected module failures.



Meteoroid Armor Thickness Requirements

Figure 7.6 [8]

METEOROID ARMOR MATERIAL PROPERTIES

Material	Density (g/cc)	$E \times 10^{-6}$ (kg/cm <sup>2</sup> )	$\rho_t E'^{-1/3}$	$\rho_t^{2/3} E'^{-1/3}$
Beryllium	1.85	2.9	0.57	1.1
Aluminum	2.7	0.63	0.83	2.2
Iron	7.9	2.1	0.39	3.1
Molybdenum	10.2	3.3	0.31	3.2
Stainless Steel	7.9	1.9	0.4	3.2
Copper	8.9	1.1	0.46	4.1

1 kg/cm<sup>2</sup> = 14.4 psi

Table 7.2 [8]

## References

1. W. R. Corliss, "Propulsion Systems for Space Flight", pp 19-62, McGraw-Hill Book Company, Inc., New York, Toronto, London, 1960.
2. H. F. Crouchy, "Nuclear Space Propulsion", chapter 12, Astronuclear Press, Granada Hills, California, 1965.
3. R. Jahn, "Physics of Electric Propulsion", chapter 8, pp 240-246, McGraw-Hill Book Company, Inc., 1968.
4. E. K. Inall, "High Power High Energy Pulse Production and Application", chapter 1, ANU Press Canberra, 1978.
5. F. P. Incropera, D. P. DeWitt, "Fundamentals of Heat and Mass Transfer", chapter 12, John Wiley & Sons, New York, Chichester, Brisbane, Toronto, Singapore, 1985.
6. J. H. Rust, "Nuclear Power Plant Engineering", chapter 8, S. W. Holland Company, Atlanta, Georgia, 1979.
7. R. W. Bussard, R. D. DeLauer, "Fundamentals of Nuclear Flight", pp 360-367, McGraw-Hill Book Company, New York, St. Louis, San Francisco, London, Sydney, 1967.
8. H. M. Dieckamp, "Nuclear Space Power Systems", pp 112-123, Atomics International (a division of North American Aviation, Inc.), Canoga Park, California, 1967.

#### IV. Conclusions and Recommendations

Several aspects of the design project should be further examined and a more detailed analysis performed. Additional neutronics work could include two dimensional and multigroup computations and examining the possibility of reducing the core size by increasing the enrichment of the fuel. Thermal hydraulics calculations could be modified by introducing radial conduction into the model which would lead to two dimensional computations. Changes in the Brayton cycle would consist of including an intercooler and regenerator to improve cycle efficiency. Modifications in the heat rejection system design would be performing a finite difference analysis on the fins and devising a mechanism so that the radiator system could be collapsed for entry into the Martian atmosphere. Additional propulsion analysis should include a more detailed calculation of trajectory, a more specific entry and landing plan such as aerobraking, and reduction of trip time. Overall areas which could be more closely examined are pressure losses and materials. A more realistic piping diagram should be developed and piping losses calculated for the system. Also, the pressure drop across the grid plate of the reactor needs to be reanalyzed. A materials search needs to be performed to identify materials compatible with carbon dioxide at high temperatures and identify a material for the turbine/compressor shaft.

## V. Summary

### Overall Hardware Layout

The layout of the pebble bed nuclear space reactor and associated equipment can be seen in Figure V.1 which is a spacecraft conceptual configuration. A more detailed configuration of the power system can be observed in Figure V.2. Note the jet engine type design where the turbine, reactor, and compressor are welded together. The reactor configuration can be seen in Figure V.3. Also note that the heat pipe radiators complete the power configuration loop but are not shown.

### System Schematic:

The thermodynamic cycle can be seen in Figure V.4. This is a Brayton cycle consisting of CO<sub>2</sub> gas coolant and the reactor, turbine, heat pipe radiator, and compressor loop. The four state points on the figure will be referred to later.

### Final Design Description:

The final design description can be observed in Tables V.1 through V.7. The only parameters not listed in the previous tables are the total mass of the CO<sub>2</sub> coolant which is 17.2 kg and the mass flow rate of 2.43 kg/s.

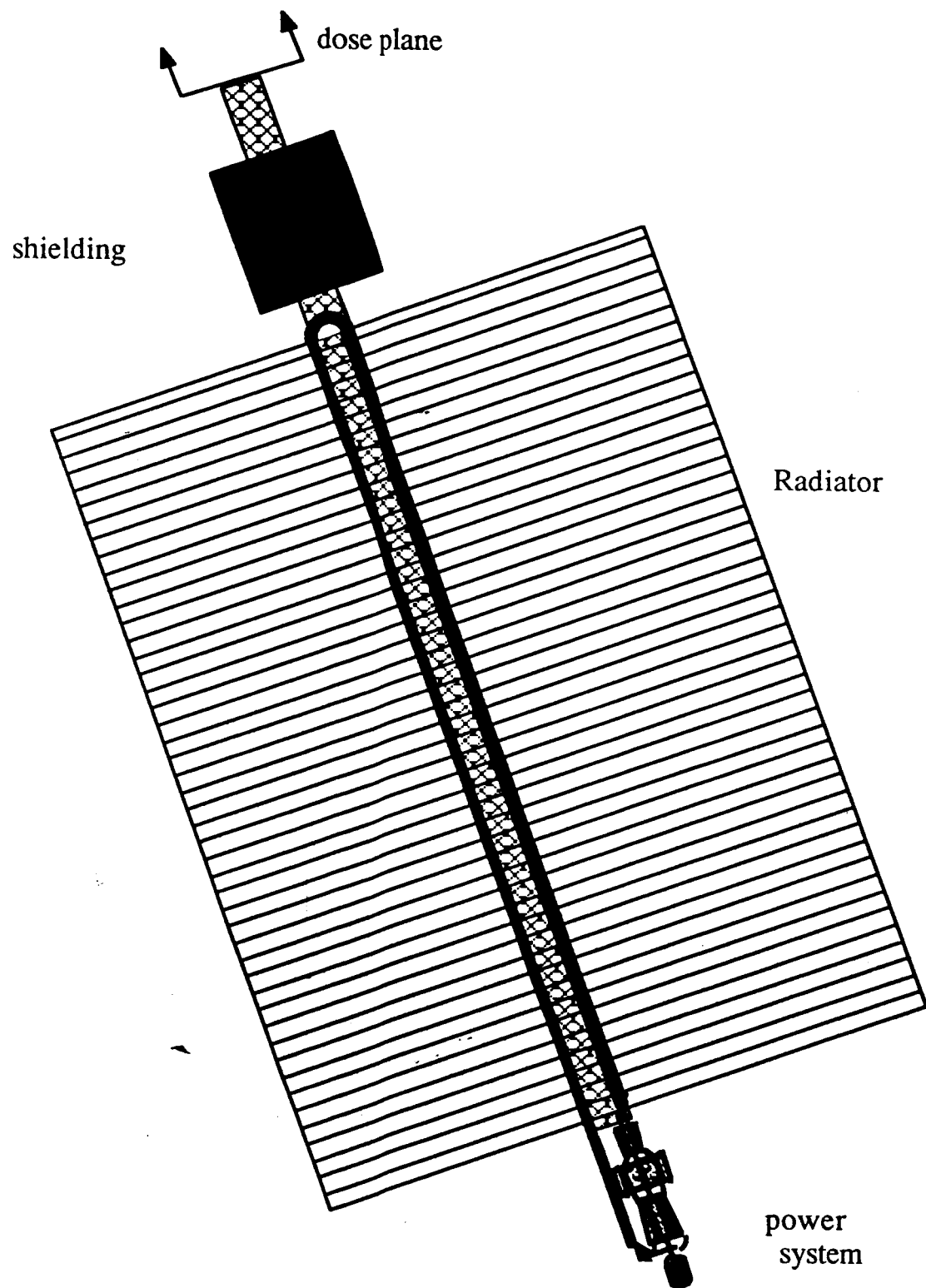


Figure V.3 Space Craft Configuration

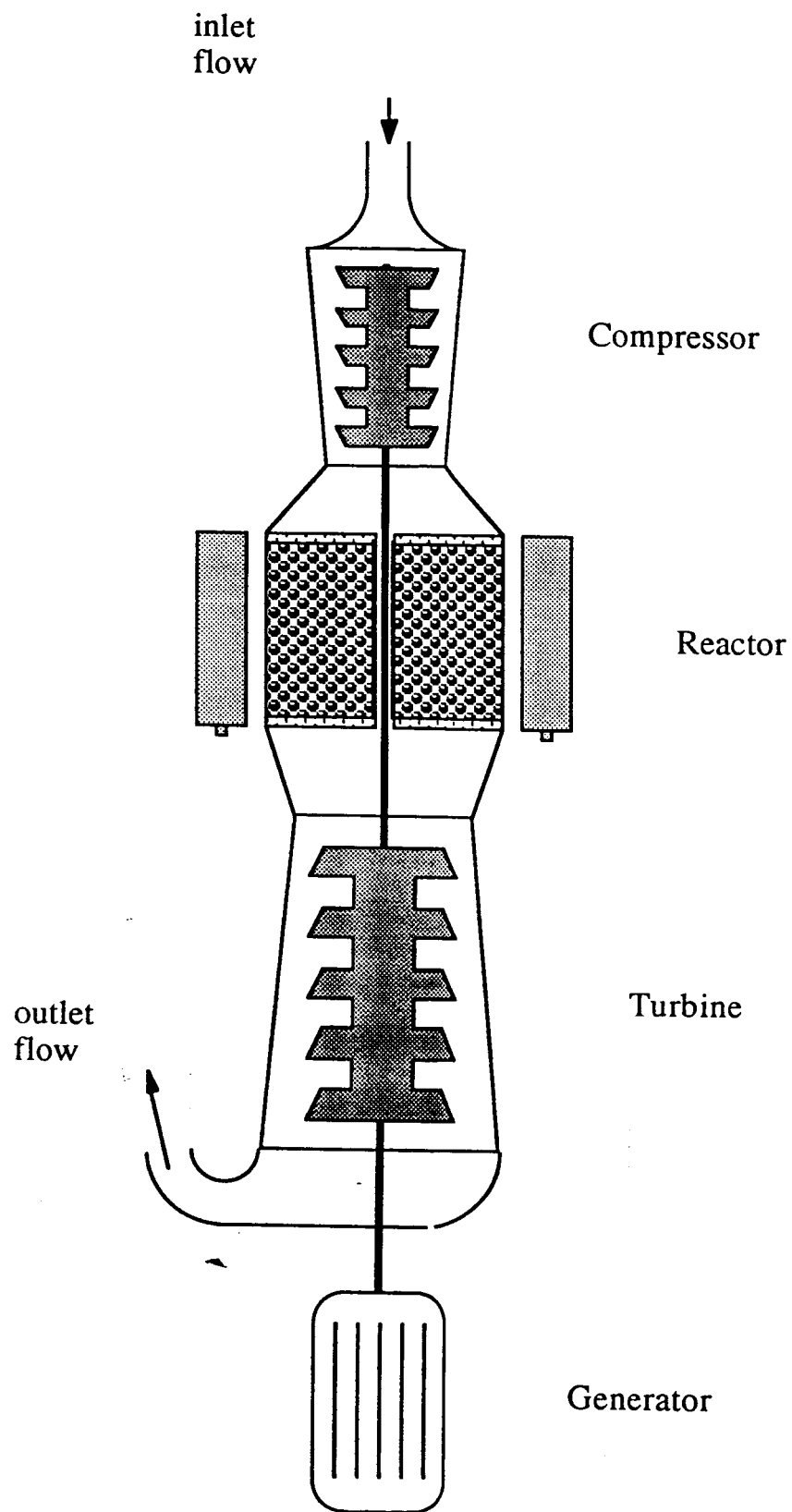


Figure V.2 Power System

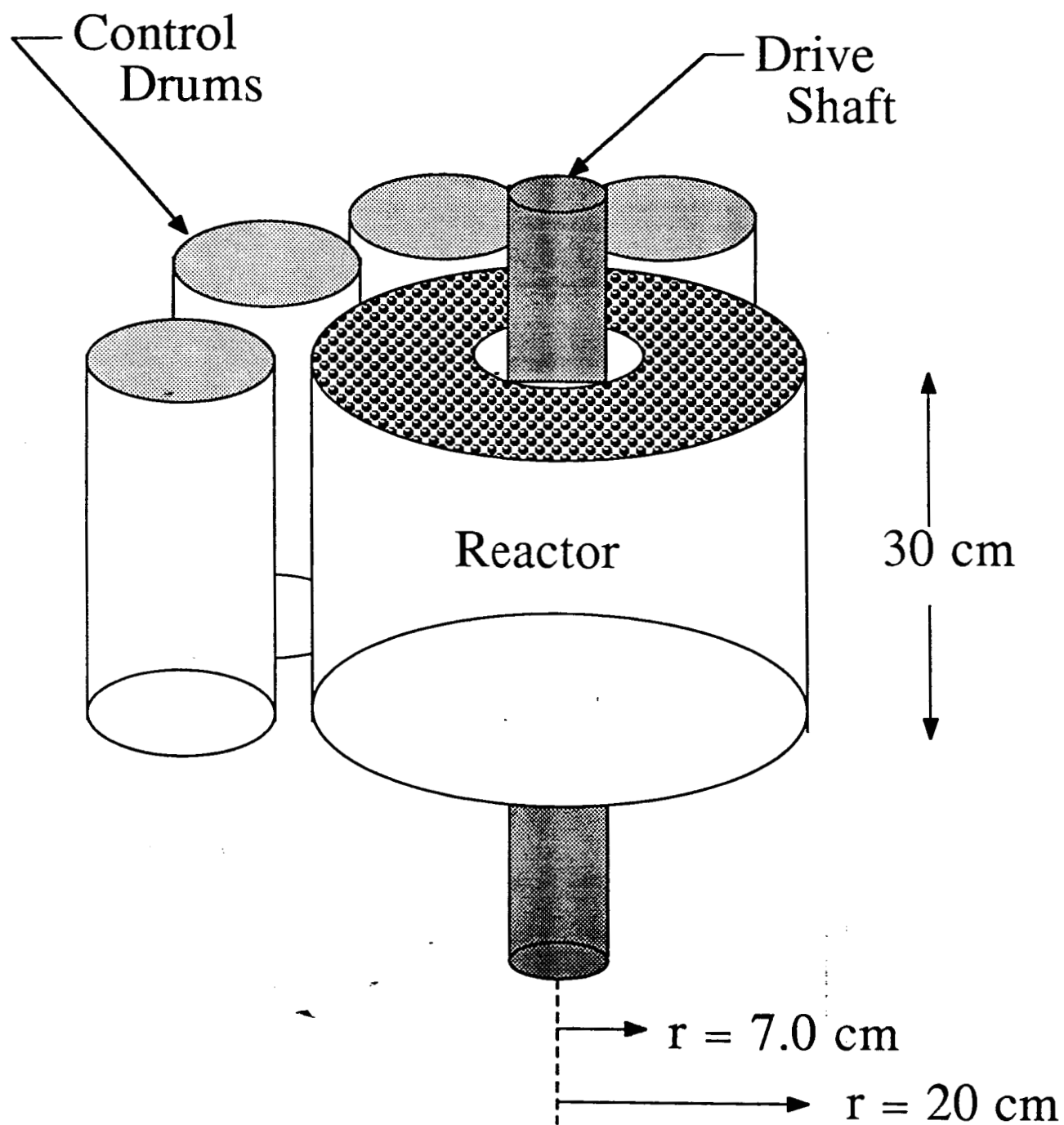


Figure V.3 Conceptual Reactor Design



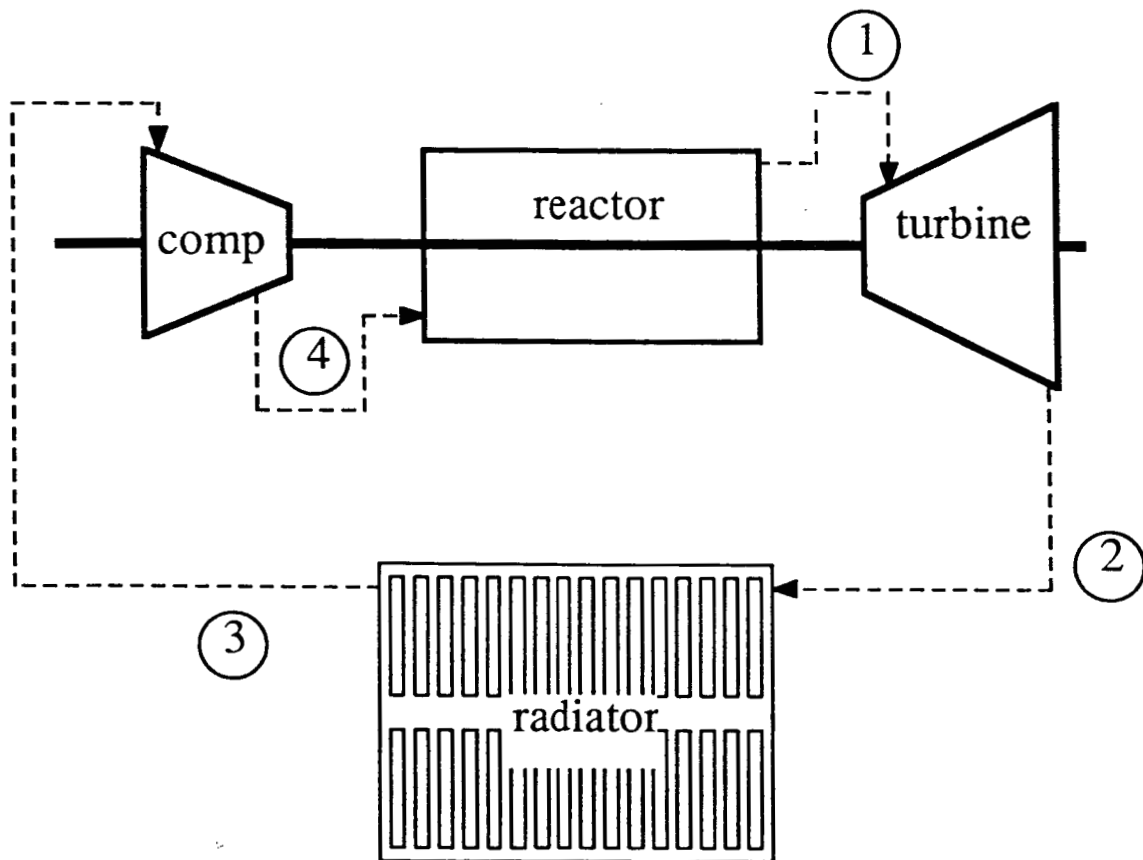


Figure V.4 System Brayton Cycle

Table V.1: Reactor parameters

Thermal Output	1 MWt
Electrical Output	300 kWe
Power Density	3.1 W/cm <sup>3</sup>
Active Core Volume	32,330 cm <sup>3</sup>
CO2 fraction	0.37
UN fraction	0.315
PyC fraction	0.315
Peak to Average Flux Ratio	1.12
Power Production	
Thermal Flux	91%
Fast Flux	9%
Control Drum Diameter	10 cm
Reactor core annular radius	12.8 cm
Central Shaft Diameter	6 cm
Reactor Vessel Wall Diameter	
inner	4 cm
outer	3 cm
Peak Fluxes	
thermal	4.7E13
fast	7.3E14
Reactor Height	30 cm
Height to diameter ratio	0.81

Table V.2 Fuel Parameters

UN Fuel Mass	145.7 kg
Fuel Enrichment	8.0%
Mass of U-235	10.8 kg
critical mass at startup	7.6 kg
fuel burnup	3.2 kg
Specific Power	6.86 kW/kg UN
# of Fuel Pebbles	180,000
Pebble Diameter	0.602 cm

Table V.3 System Thermodynamic States\*

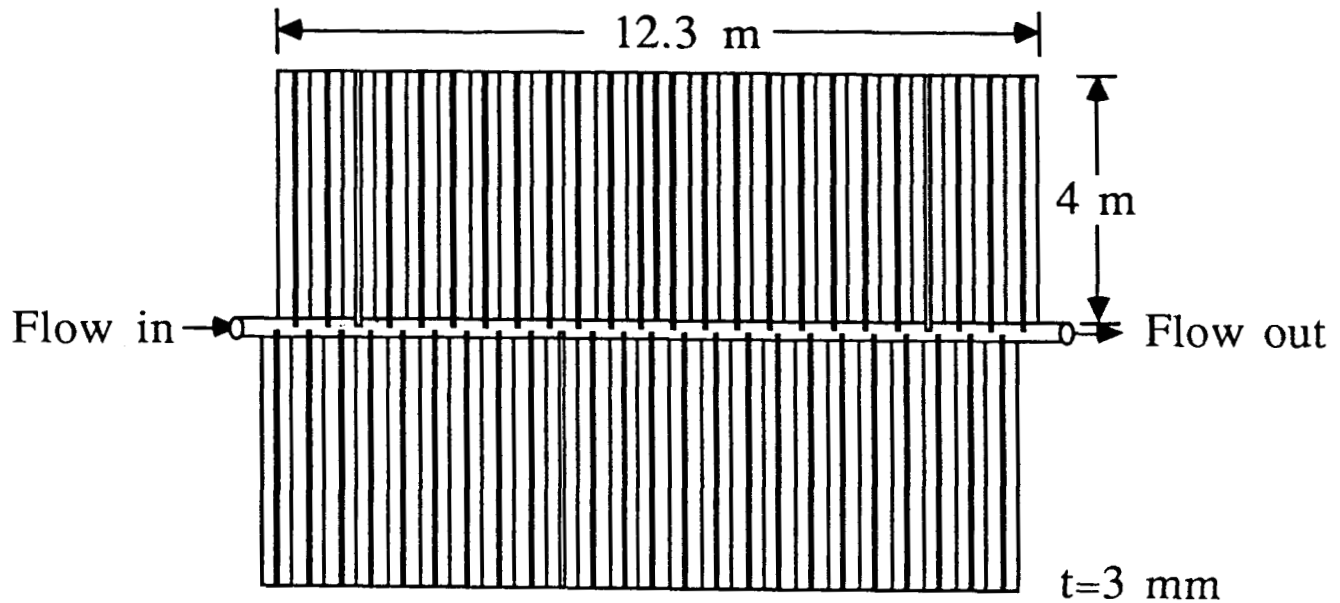
State	Temperature (K)	Pressure (MPa)
1	850	6.58
2	677	3.76
3	428	3.63
4	500	6.90

Table V.4 System Energy Balance

$Q_{in}$	=	1038 kW
$Q_{out}$	=	738 kW
$W_T$	=	513 kW
$W_C$	=	213 kW
$h$	=	28.9 %

\* Refer to Figure V.3 for the state points.

Table V.5 - Heat Rejection System



• Total Area	98 m <sup>2</sup>
• Mass	1870 kg
• Heat Pipes	246
• Container Material	SS 304
• Coating	Al <sub>2</sub> O <sub>3</sub>
• Working Fluid	Sodium
• Fin Material	Aluminum

Table V.6 Propulsion

• MPD Thrusters	13 units
• Total Craft Mass	8000 kg
• Total Thrust	3.6 N
• Propellant	NH <sub>3</sub>
• Mass	1571 kg
• Flight Time	242 days
• MPD power	296 kW
• Specific Impulse	6140 seconds

Table V.7 Masses

• Propulsion	2350 kg
• Reactor system and shield	1029 kg
• Rotating machinery	565 kg
• Ducting and structure	700 kg
• Heat pipes	2130 kg
• Miscellaneous	480 kg
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• Total Mass	7254 kg